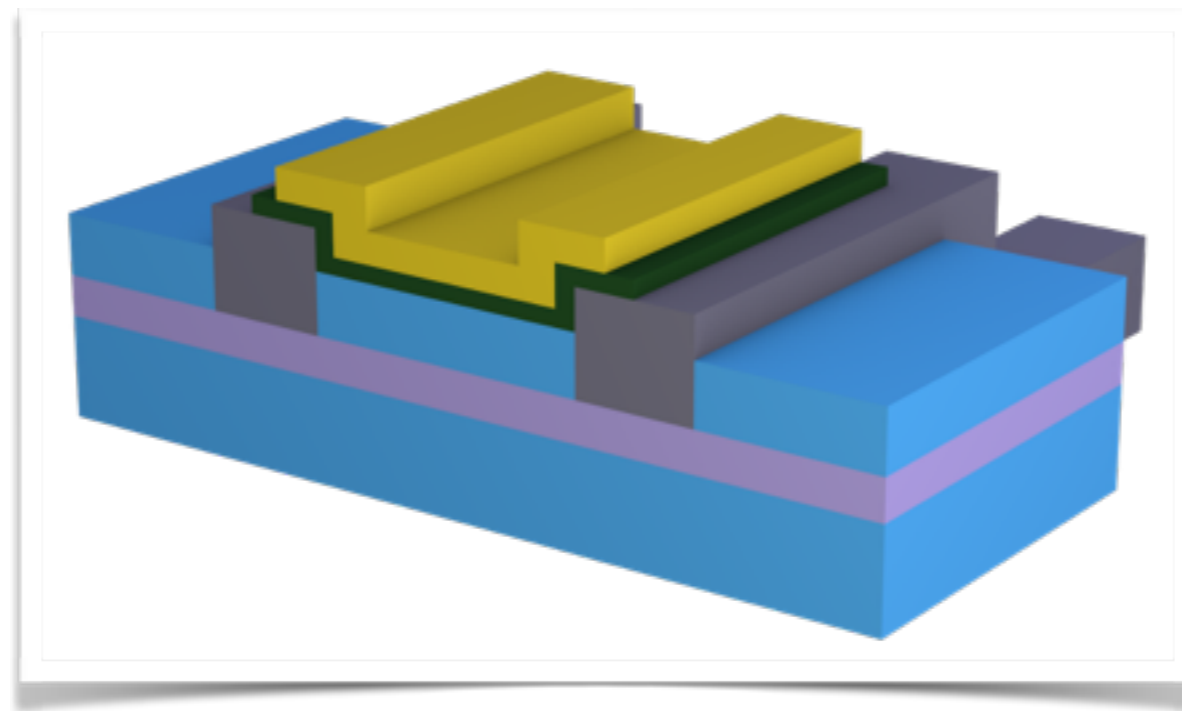


Erwann Bocquillon

Gapless Andreev bound states in HgTe-based topological Josephson junctions

TopoLyon 2016
03/10/2016

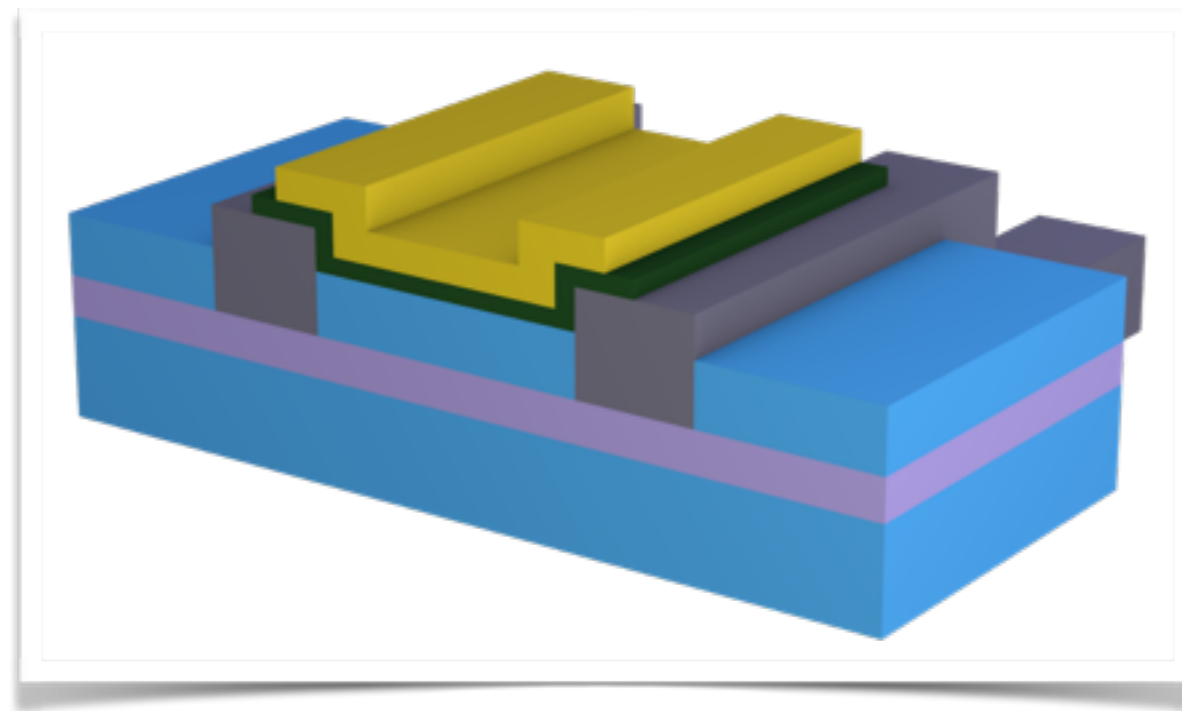


Physikalisches Institut (EP3)
Universität Würzburg, Am Hubland, D-97074 Würzburg
<http://www.physik.uni-wuerzburg.de/EP3/>

Erwann Bocquillon

Gapless Andreev bound states in HgTe-based topological Josephson junctions

TopoLyon 2016
03/10/2016



Physikalisches Institut (EP3)
Universität Würzburg, Am Hubland, D-97074 Würzburg
<http://www.physik.uni-wuerzburg.de/EP3/>



Uni. Würzburg

- ▷ PhD students : J. Wiedenmann,
P. Leubner
- ▷ Staff : C. Brüne
H. Buhmann
L.W. Molenkamp
- ▷ Invited : T.M. Klapwijk
- ▷ Theory : F. Domínguez, E.M. Hankiewicz



European Research Council
Established by the European Commission

Unterstützt von / Supported by



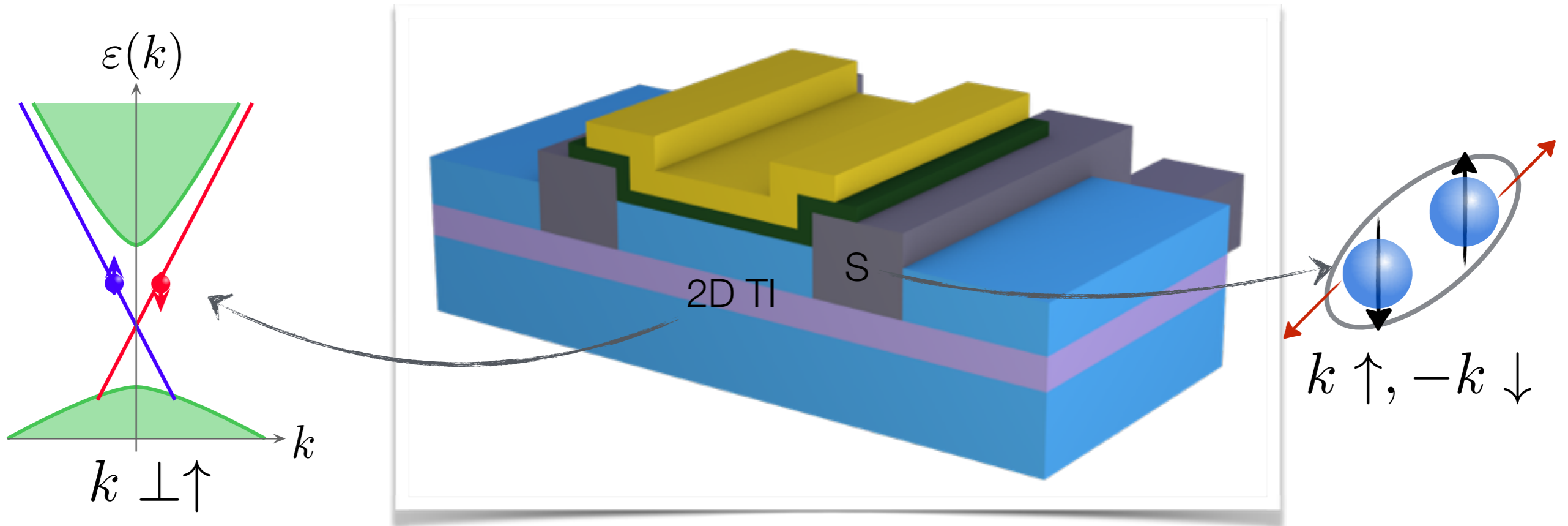
Alexander von Humboldt
Stiftung / Foundation



RIKEN, Tōkyō

- ▷ R.S. Deacon, K. Ishibashi, S. Tarucha

Induced superconductivity in a 2D TI



Cooper pair of helical Dirac fermions
 \Rightarrow helical pairing
 \Rightarrow p -type correlations

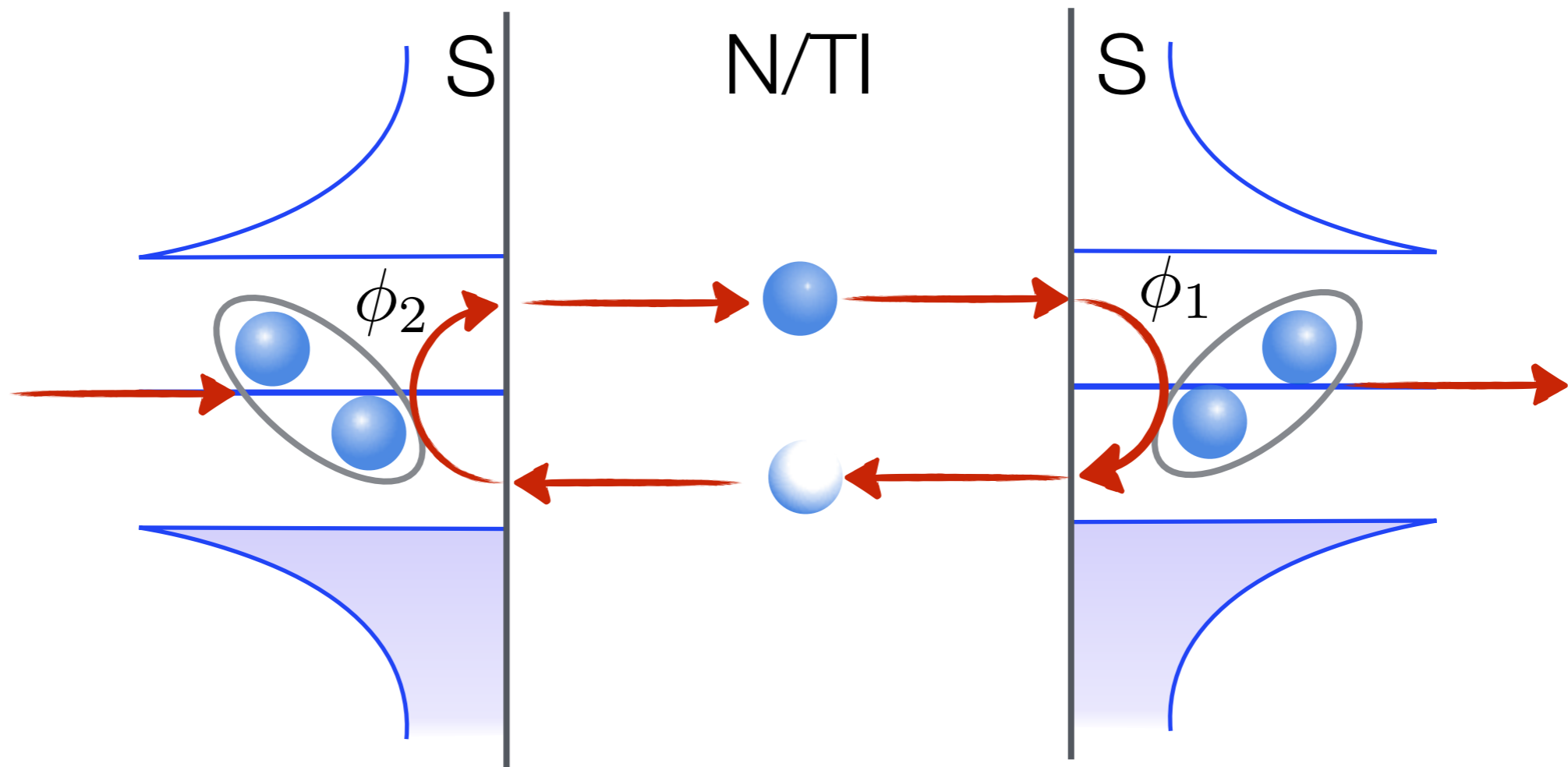
gapless Andreev bound states
 \Rightarrow Majoranas

Fu et al., PRB **79**, 161408 (2009)

spin-orbit coupling
 \Rightarrow φ_0 -junctions

Dolcini et al., PRB **92**, 035428 (2015)

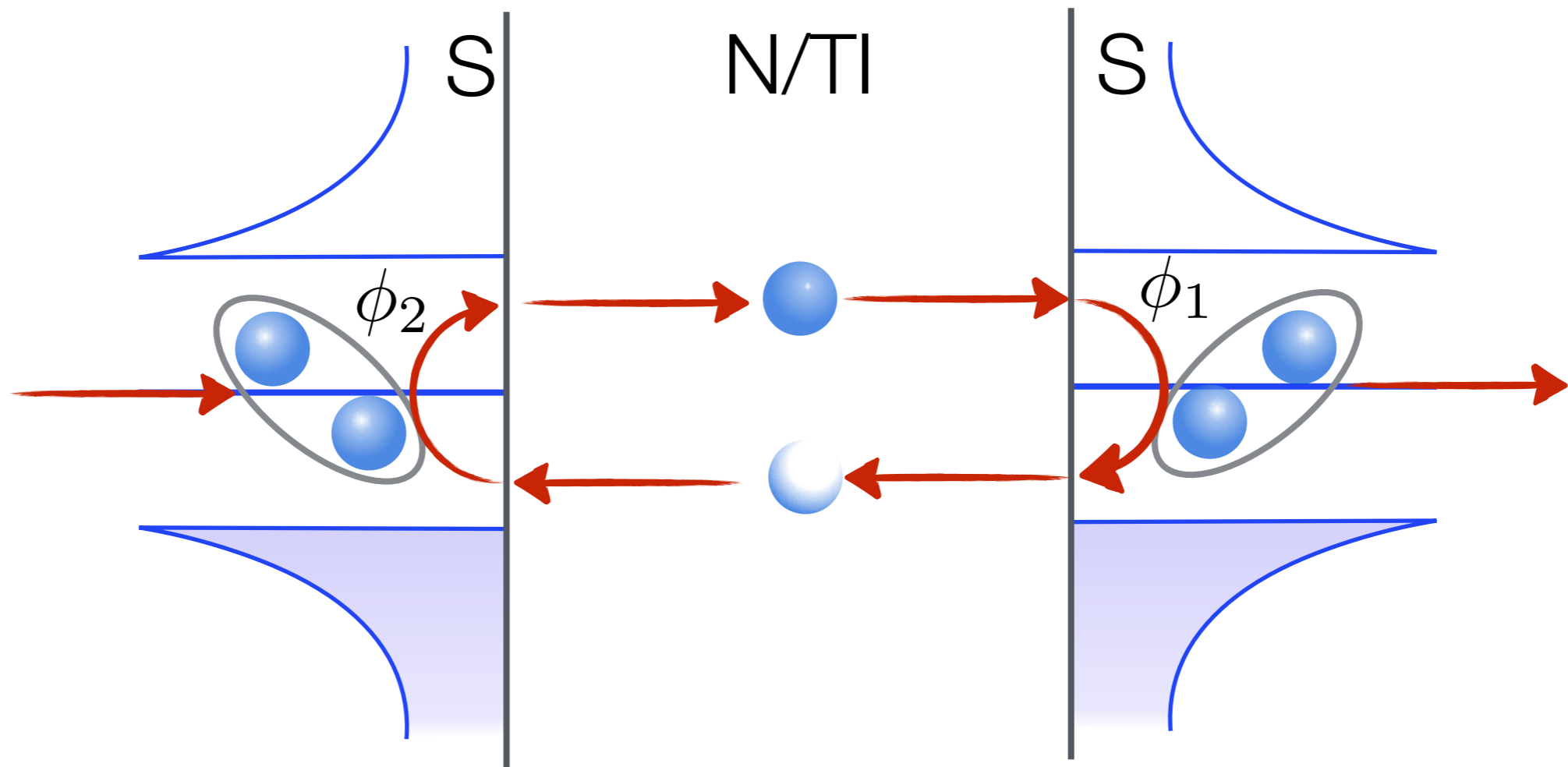
Andreev reflections



Andreev reflection

- ▷ Cooper pair created in S
- ▷ hole reflected in N
- ▷ phase coherence

Andreev reflections



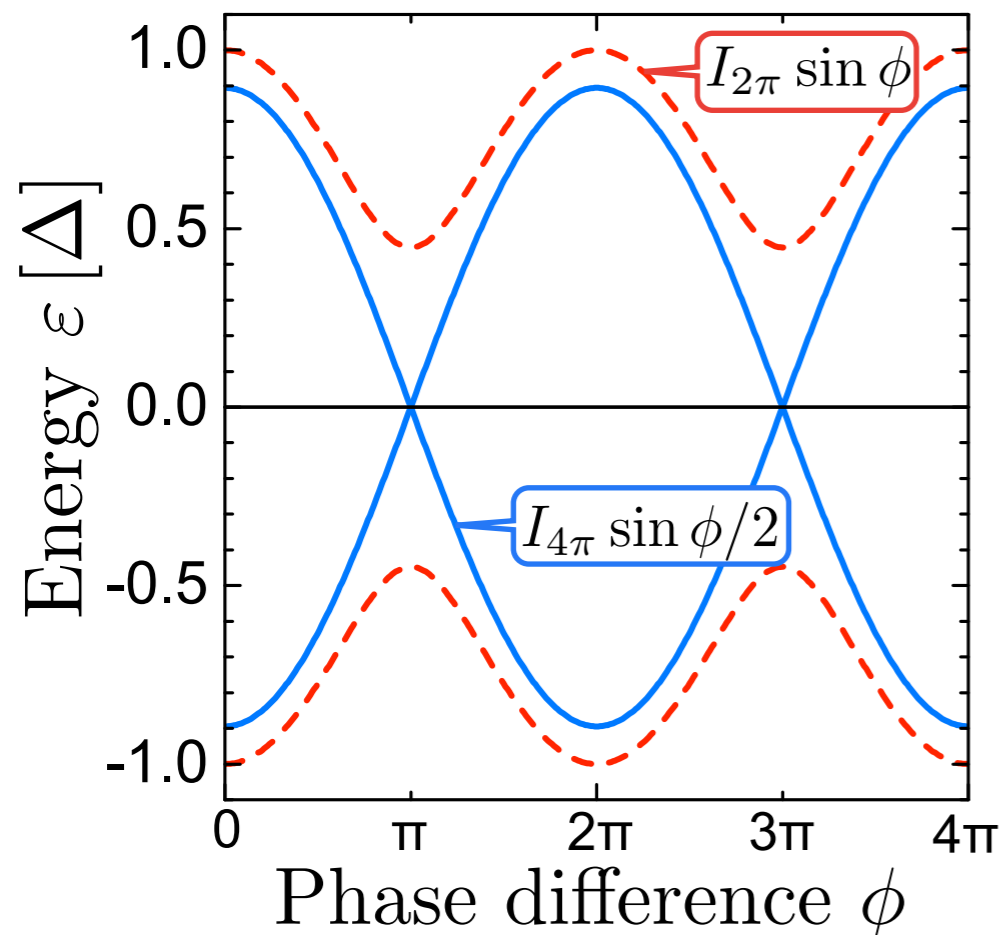
Andreev reflection

- ▷ Cooper pair created in S
- ▷ hole reflected in N
- ▷ phase coherence

Andreev bound states

- ▷ resonant modes of the "cavity"
- ▷ energy levels $\varepsilon_n(\phi)$
- $\phi = \phi_1 - \phi_2$

Gapless Andreev bound states



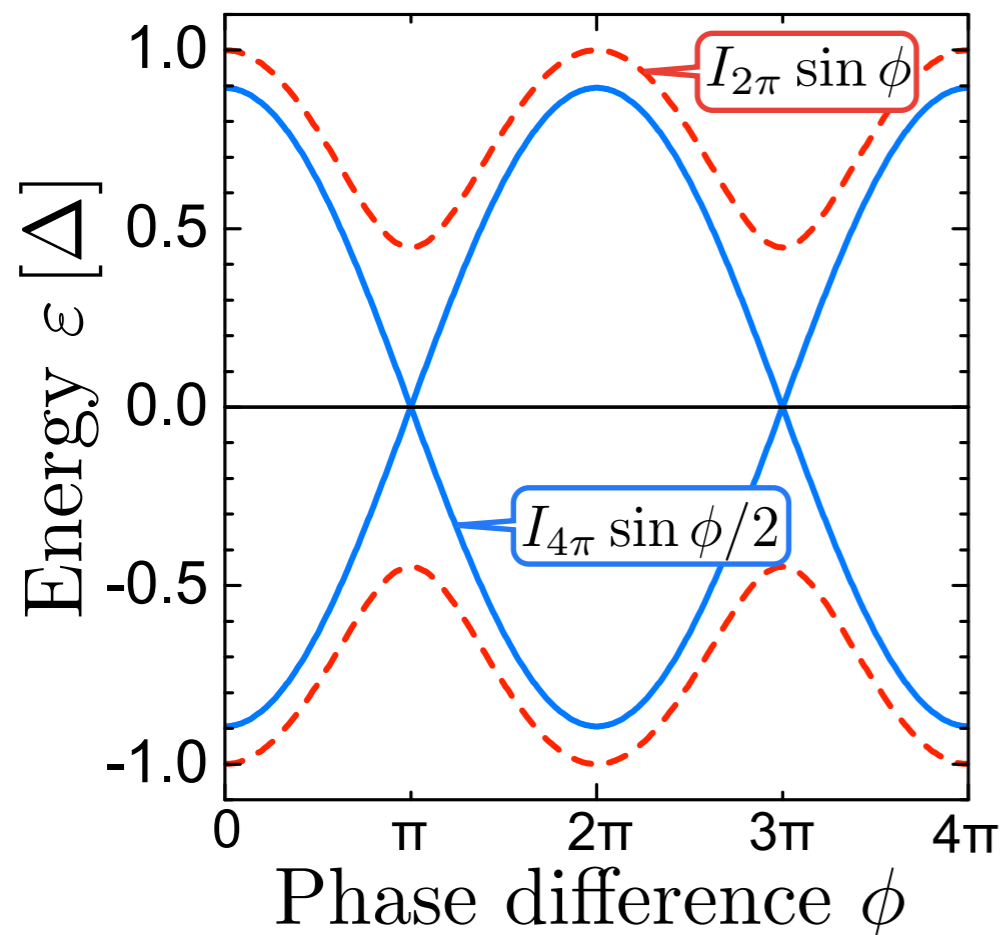
Andreev bound states

- ▶ 2π gapped states (bulk)
 $I_{2\pi} \sin \phi$ (+harmonics)
- ▶ 4π gapless topological state (edge)
 $I_{4\pi} \sin \phi/2$
- ▶ fractional Josephson effect

Kwon *et al.*, JLTP **30**, 613 (2004)

Fu *et al.*, PRB **79**, 161408 (2009)

Gapless Andreev bound states



Andreev bound states

- ▷ 2π gapped states (bulk)
 $I_{2\pi} \sin \phi$ (+harmonics)
- ▷ 4π gapless topological state (edge)
 $I_{4\pi} \sin \phi / 2$
- ▷ fractional Josephson effect

Kwon *et al.*, JLTP **30**, 613 (2004)

Fu *et al.*, PRB **79**, 161408 (2009)

Difficult detection

- ▷ 2π bulk states \Rightarrow $2\pi/4\pi$ mixture
- ▷ finite lifetime \Rightarrow 2π -periodicity restored
- ▷ interactions \Rightarrow 8π -periodicity
- ▷ Landau-Zener transitions \Rightarrow 4π -periodicity

Pikulin *et al.*, PRB **86**, 140504 (2012)

Badiane *et al.*, CRP **14**, 840 (2013)

Zhang *et al.*, PRL **113**, 036401 (2014)

Peng *et al.*, ArXiv 1609.01896 (2016)

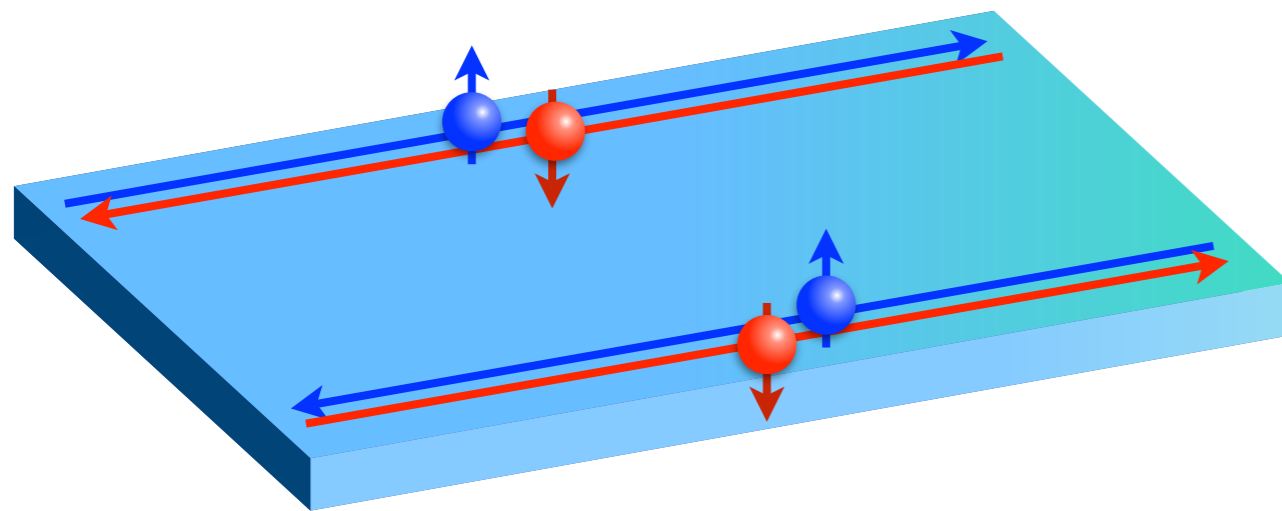
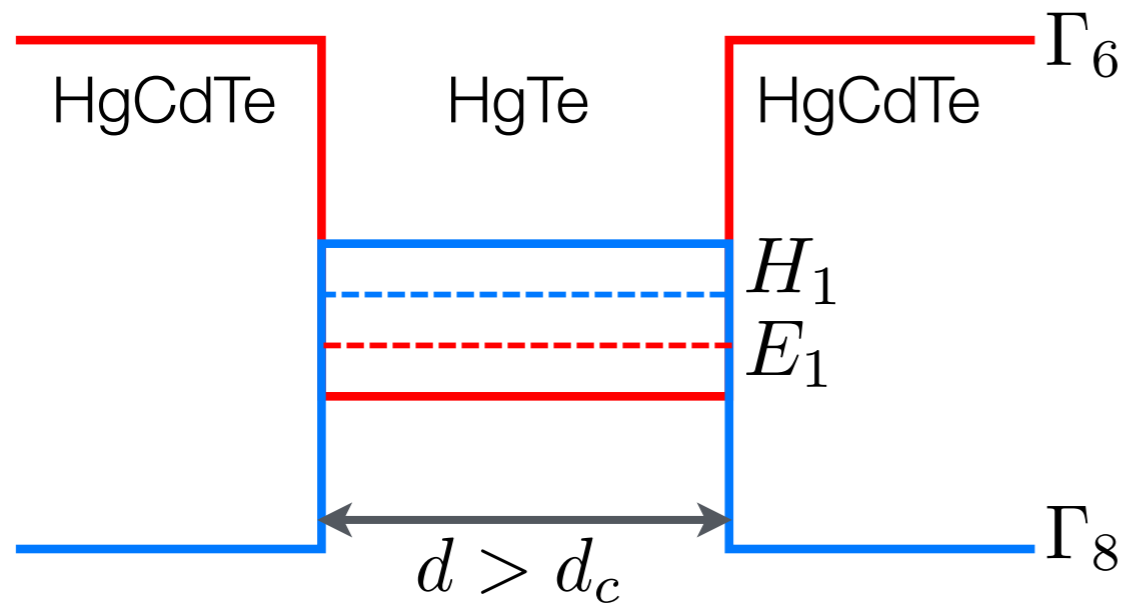
Hui *et al.*, ArXiv 1609.02909 (2016)

- ▶ Josephson junctions in a 2D TI
- ▶ Josephson emission
- ▶ Response to AC excitation (Shapiro steps)

Bocquillon *et al.*, Nat. Nano, DOI: 10.1038/NNANO.2016.159

Deacon *et al.*, ArXiv 1603.09611 (2016)

QSH effect in HgTe

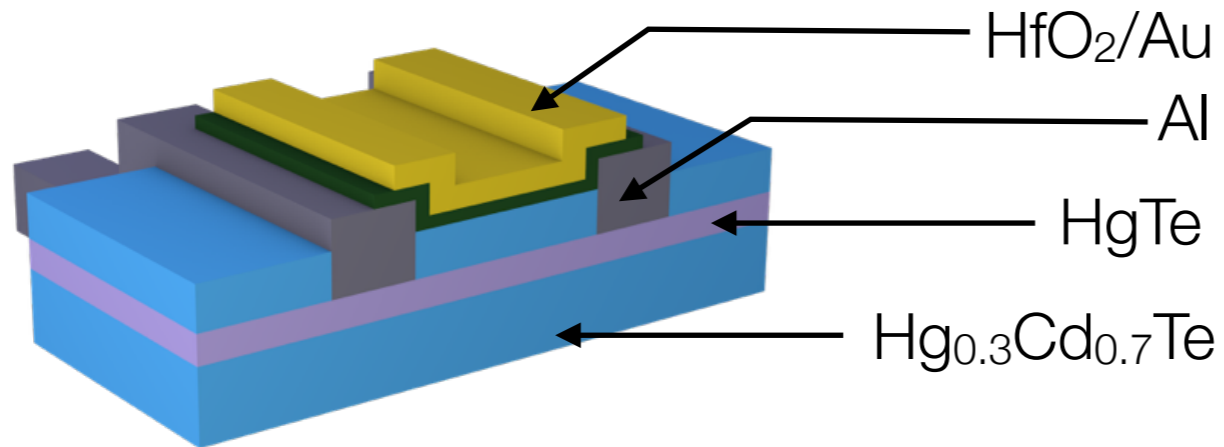


- ▷ Band inversion in bulk HgTe
⇒ inverted QW
⇒ Quantum Spin Hall effect

- ▷ MBE growth, $\mu \approx 3 \cdot 10^5 \text{ cm}^2\text{V}^{-1}\text{s}^{-1}$
- ▷ QSH if $d > d_c \approx 6.3 \text{ nm}$ at $B = 0$!
- ▷ trivial if $d < d_c$

Bernevig *et al.*, Science **314**, 1757 (2006)
König *et al.*, Science **318**, 766 (2007)

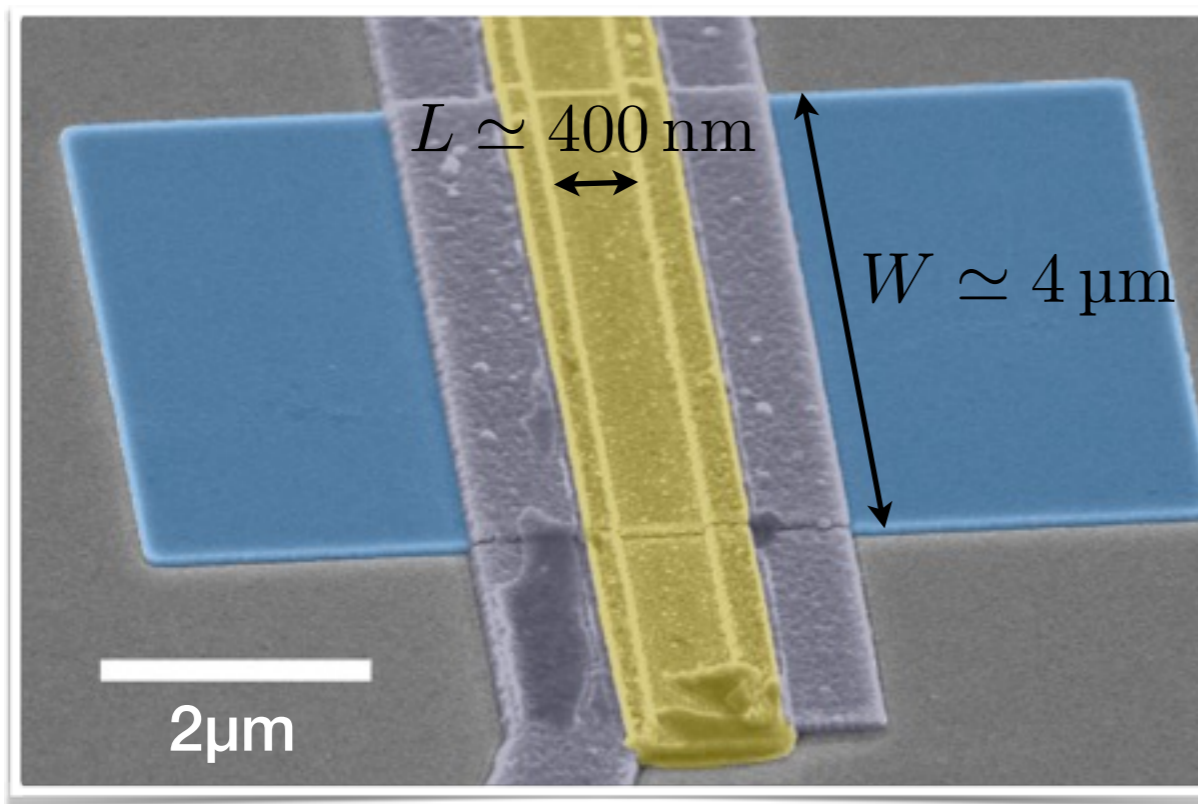
Quantum spin Hall junctions



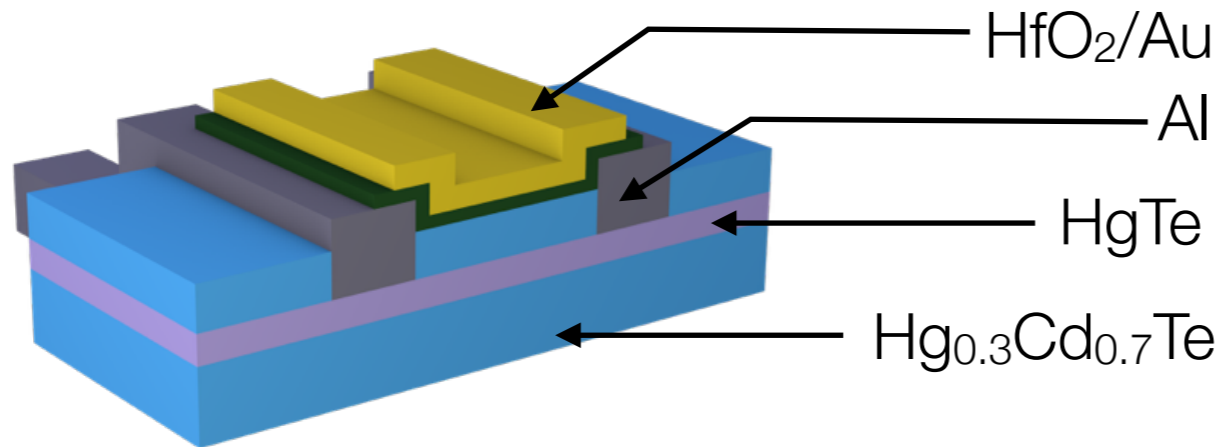
Josephson junctions

- ▷ $\mu \approx 3 \cdot 10^5 \text{ cm}^2 \text{V}^{-1} \text{s}^{-1}$
- ▷ Al contacts (in situ)
- ▷ HfO₂/Au gate
- ▷ no overlap of edge states
- ▷ ballistic / intermediate

$$L \ll l \quad L \lesssim \xi$$



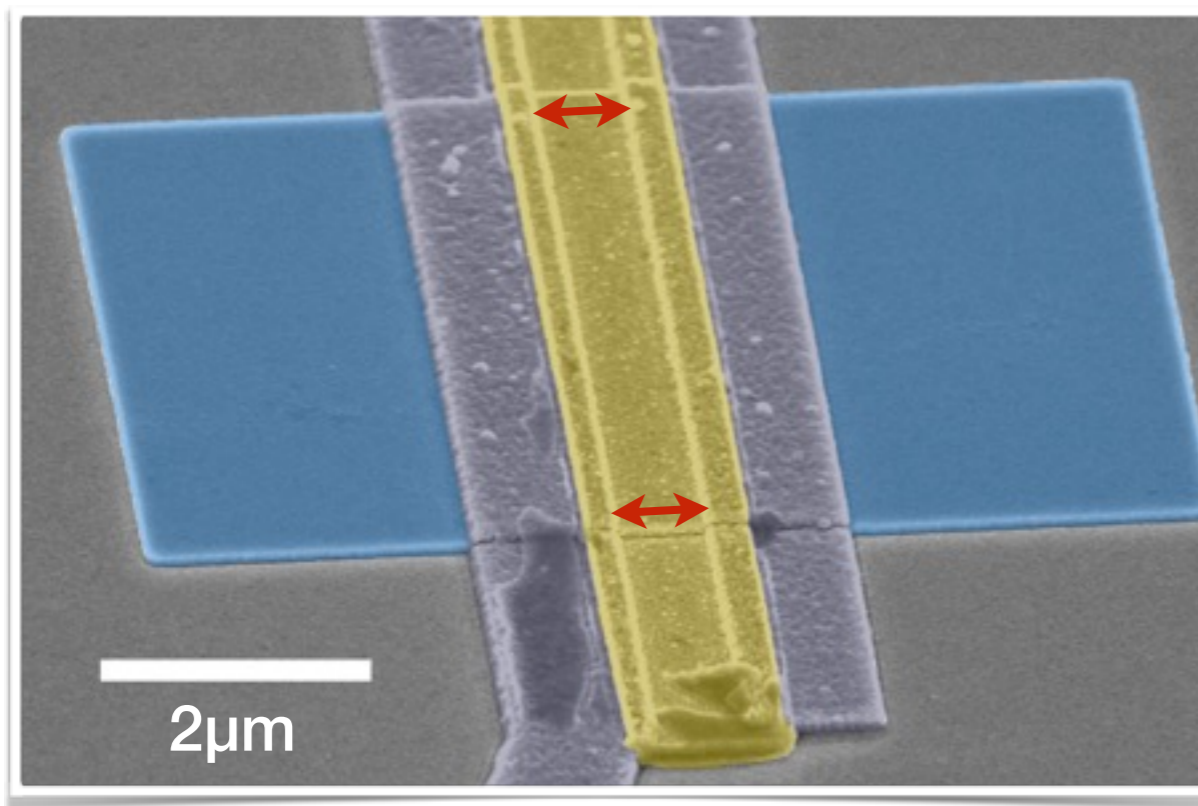
Quantum spin Hall junctions

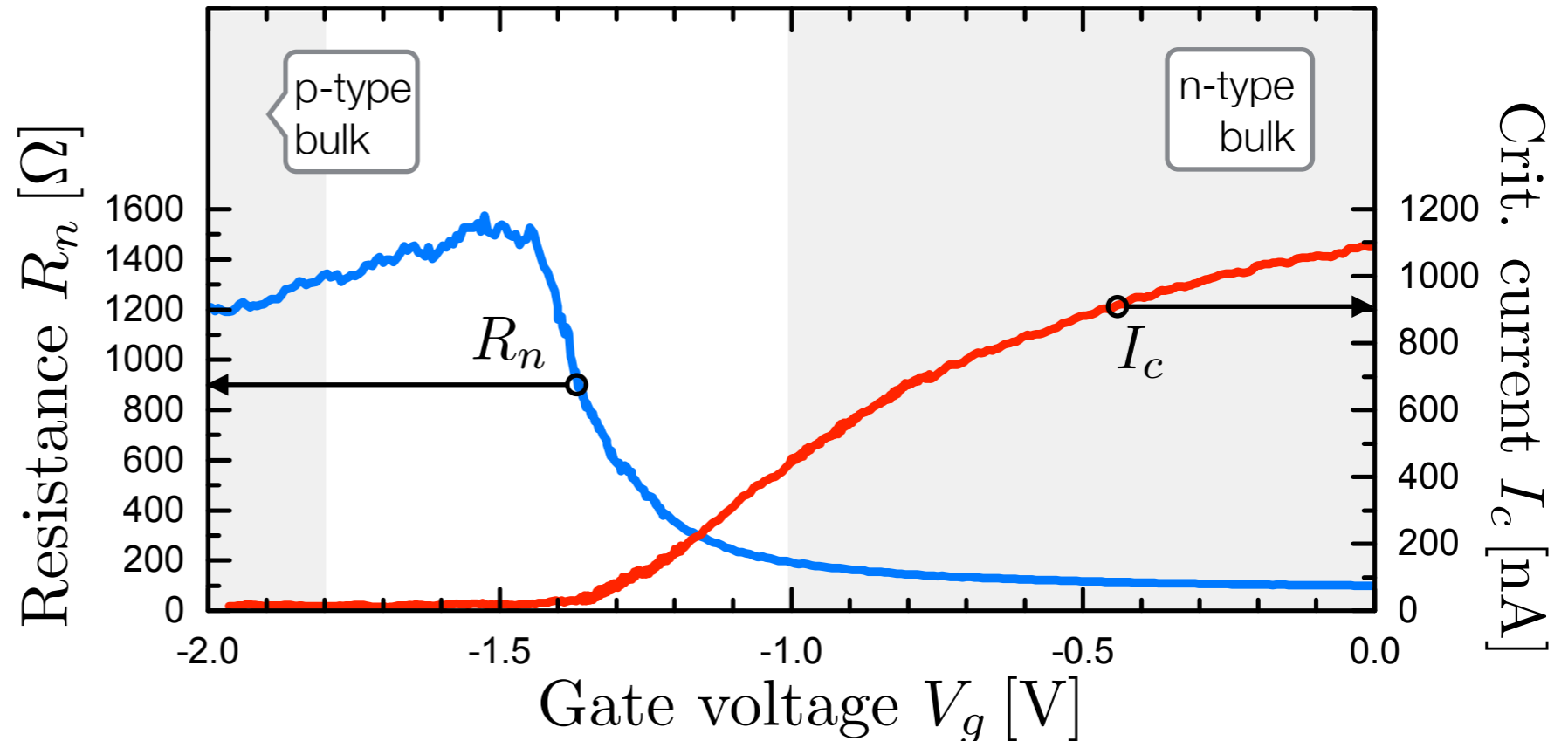
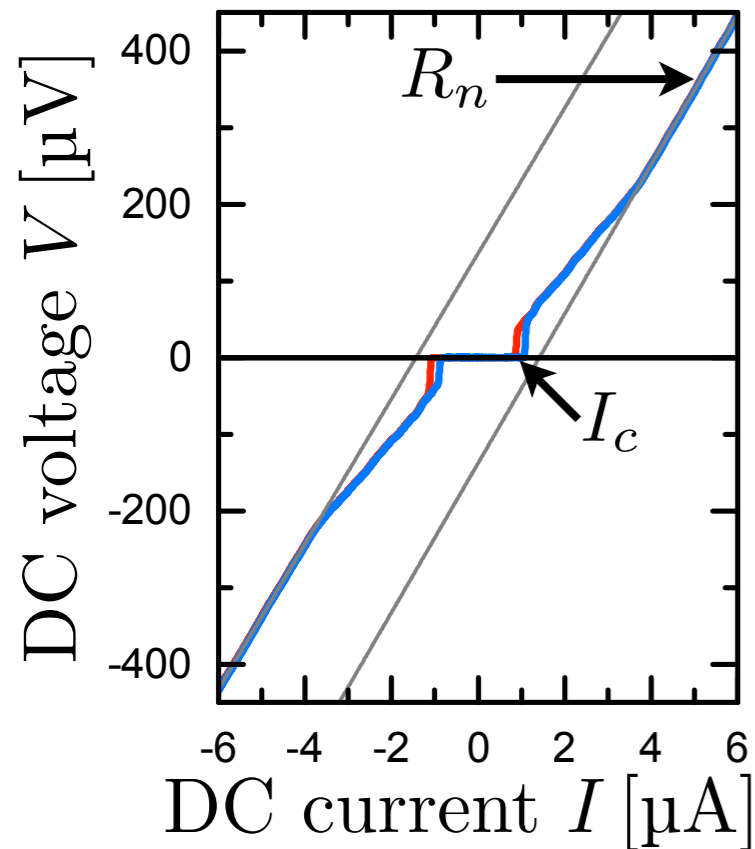


Josephson junctions

- ▷ $\mu \approx 3 \cdot 10^5 \text{ cm}^2 \text{V}^{-1} \text{s}^{-1}$
- ▷ Al contacts (in situ)
- ▷ HfO₂/Au gate
- ▷ no overlap of edge states
- ▷ ballistic / intermediate

$$L \ll l \quad L \lesssim \xi$$





I-V curve

- ▷ weak hysteresis visible
- ▷ excess current
⇒ Andreev reflections

Gate dependence

- ▷ 3 regimes : p , n , and QSH
- ▷ asymmetry between n and p

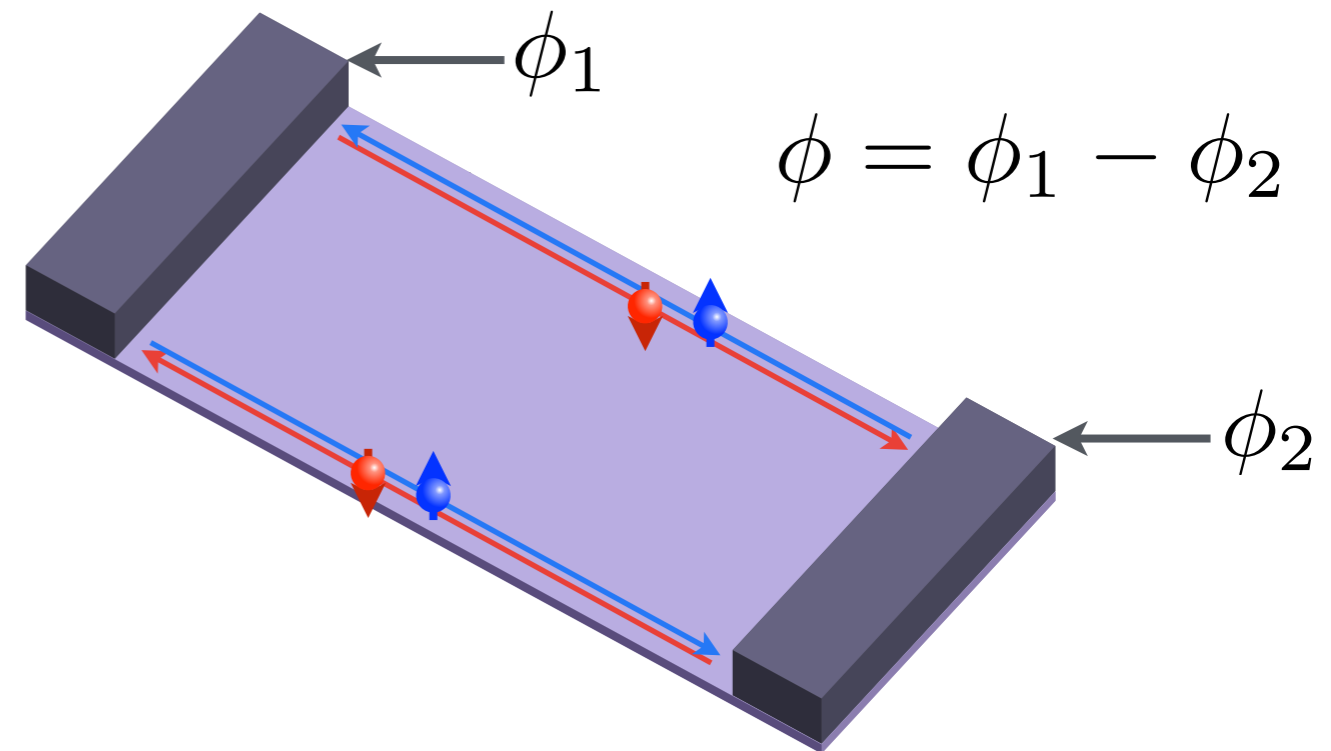
Blonder *et al.*, PRB **25**, 4515 (1982)

Josephson equations

$$\frac{d\phi}{dt} = \frac{2eV}{\hbar}$$

$$I_S(\phi) = I_c \sin \phi$$

$$\Rightarrow \text{Josephson frequency } f_J = \frac{2eV}{\hbar}$$

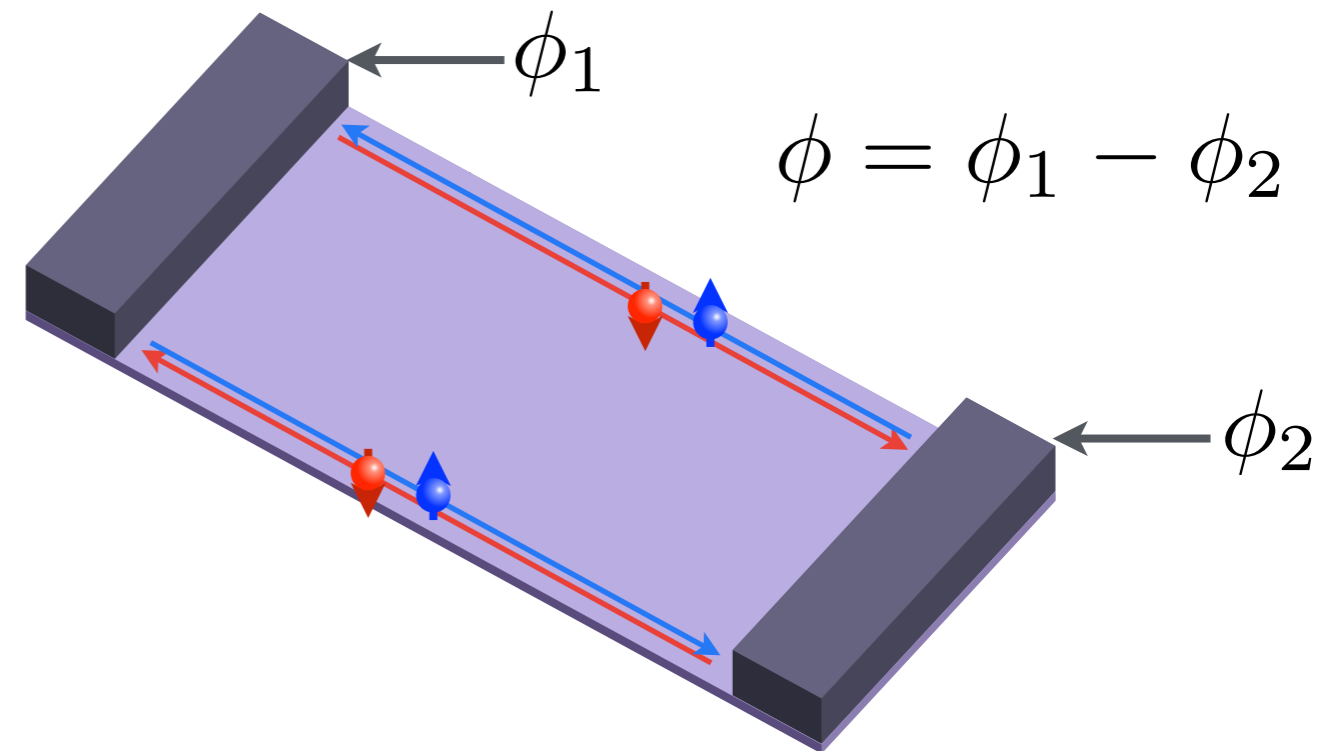


Josephson equations

$$\frac{d\phi}{dt} = \frac{2eV}{\hbar}$$

$$I_S(\phi) = I_c \sin \phi$$

$$\Rightarrow \text{Josephson frequency } f_J = \frac{2eV}{\hbar}$$



Fractional Josephson effect

$$\sin \phi \rightarrow \sin \phi/2$$

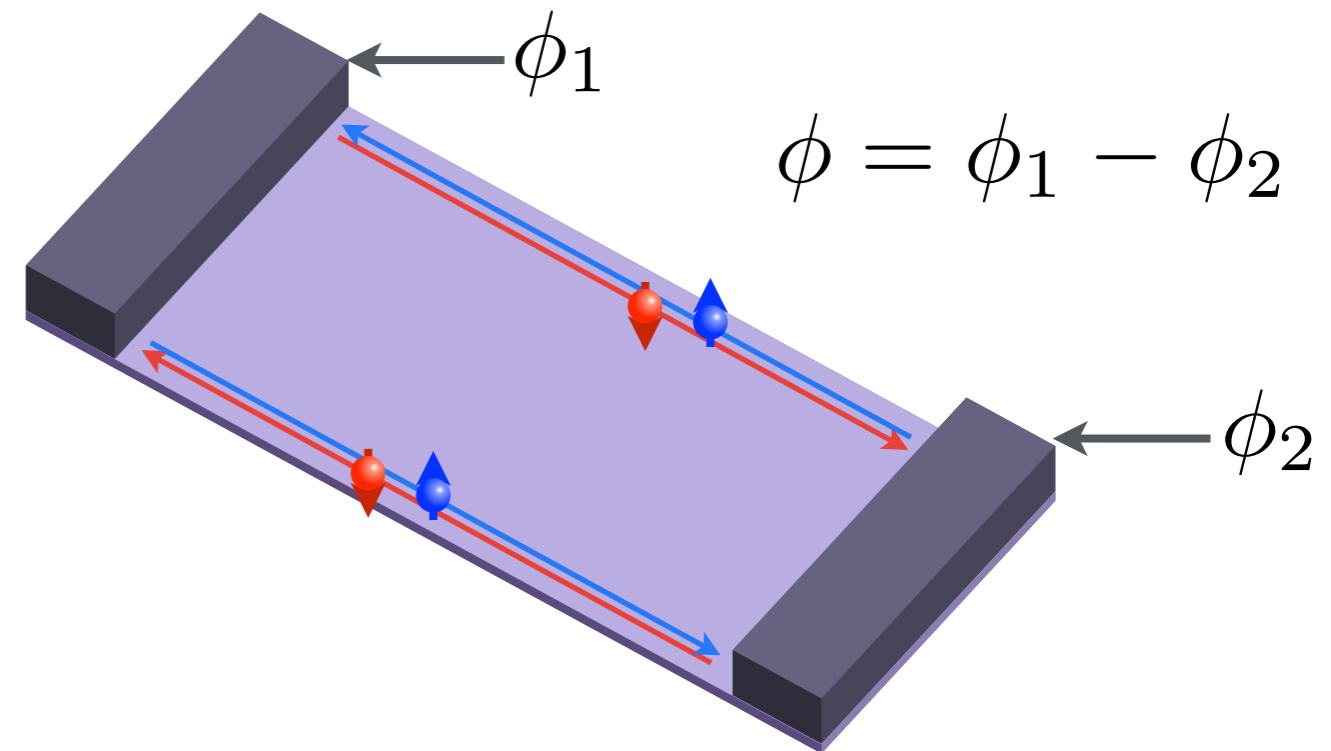
$$f_J \rightarrow f_J/2$$

Josephson equations

$$\frac{d\phi}{dt} = \frac{2eV}{\hbar}$$

$$I_S(\phi) = I_c \sin \phi$$

$$\Rightarrow \text{Josephson frequency } f_J = \frac{2eV}{\hbar}$$



Fractional Josephson effect

$$\sin \phi \rightarrow \sin \phi/2$$

$$f_J \rightarrow f_J/2$$

Detection

- ▶ ‘listening’ to Josephson emission
- ▶ beatings with ac excitation (Shapiro steps)

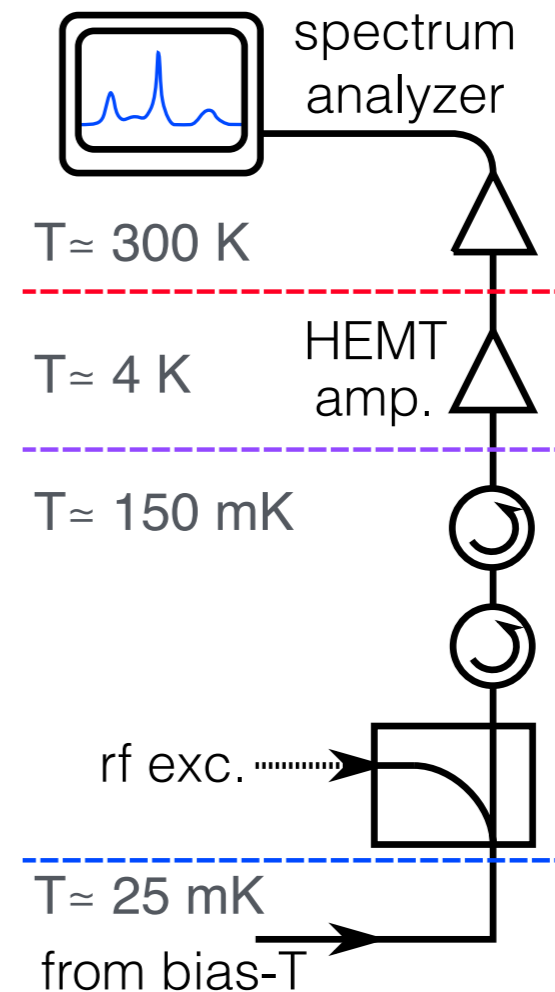
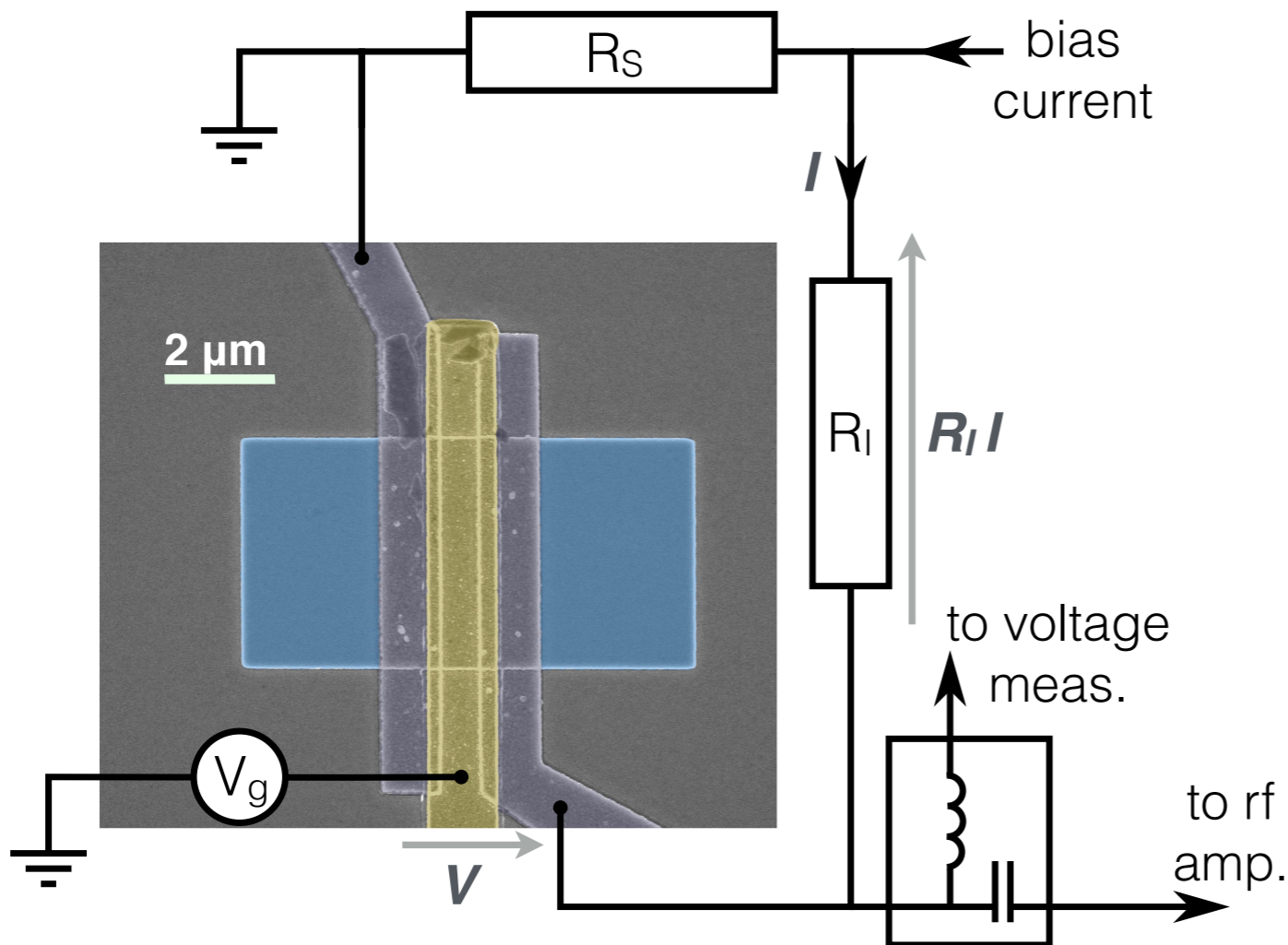
Detection setup

▷ voltage bias

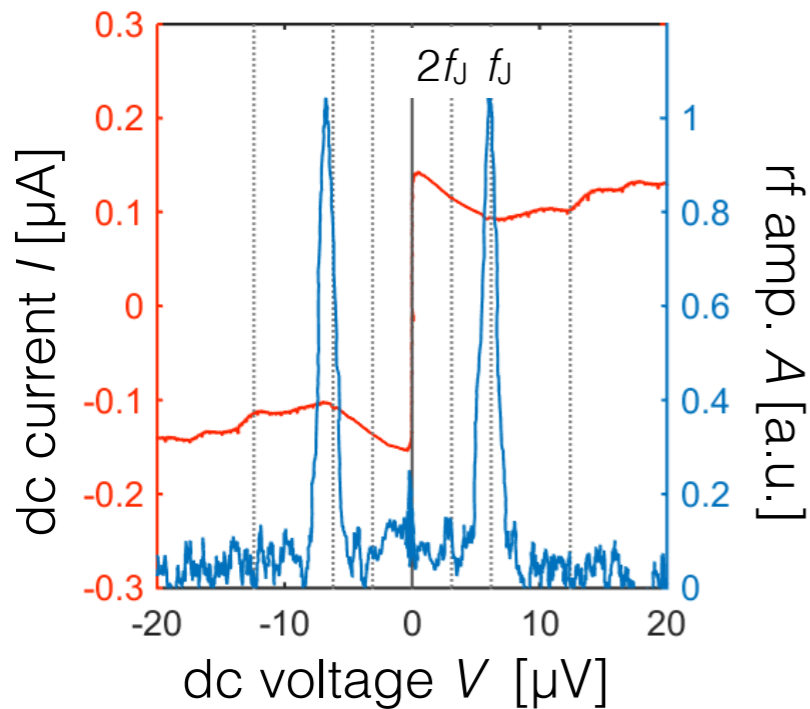
shunt resistance R_s
current resistance R

▷ rf amplification setup

1 cryo amp. (+ 2 amps at RT)
0.1 fW (-130 dBm) in 8 MHz

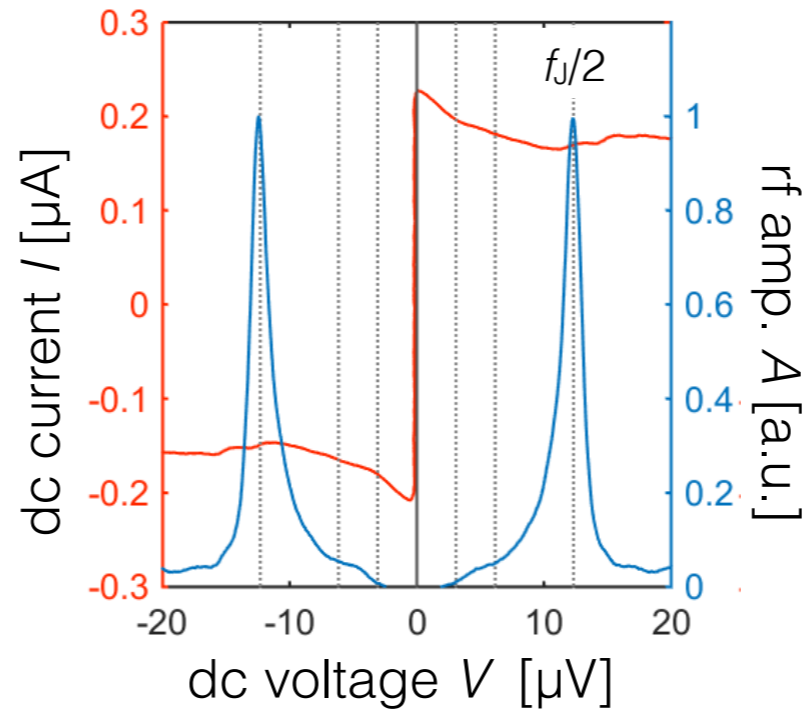


Trivial QW

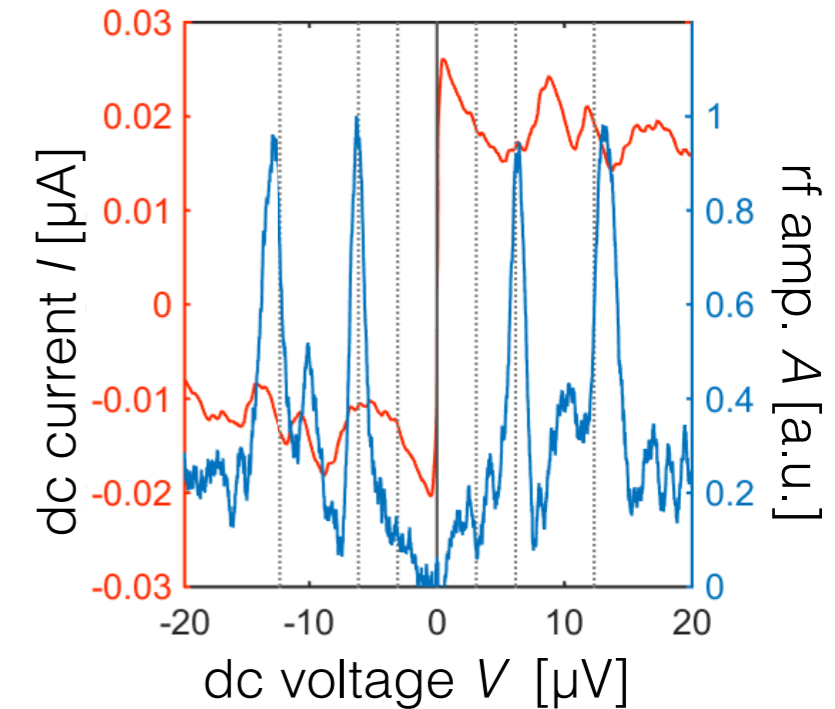


Topological QW

n- and QSH regime



p-regime



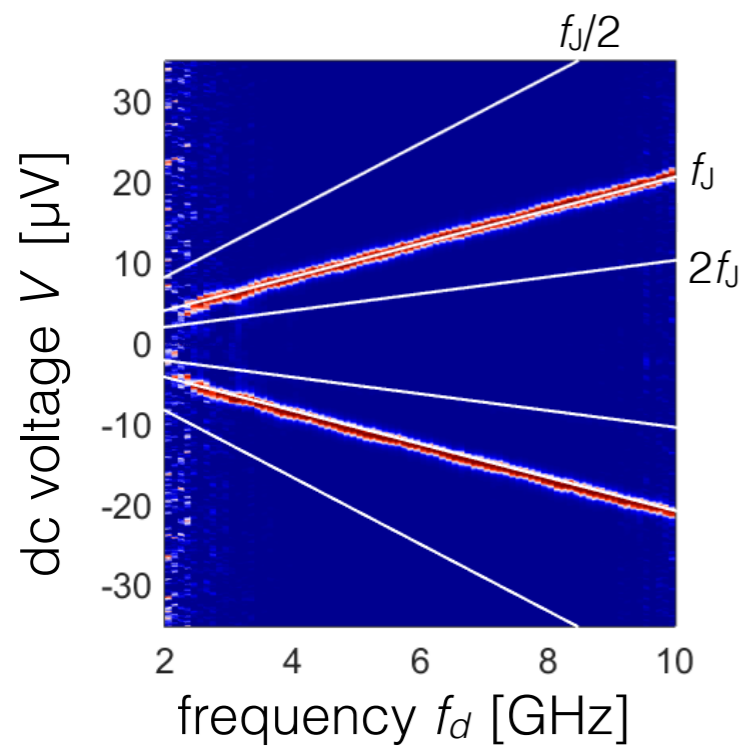
- ▶ voltage V swept
- ▶ integrated power at $f_d=3$ GHz (in 8 MHz bandwidth)

- ▶ trivial QW : signal at $f_d=f_J$
- ▶ topological QW : at $f_d=f_J$ and $f_J/2$

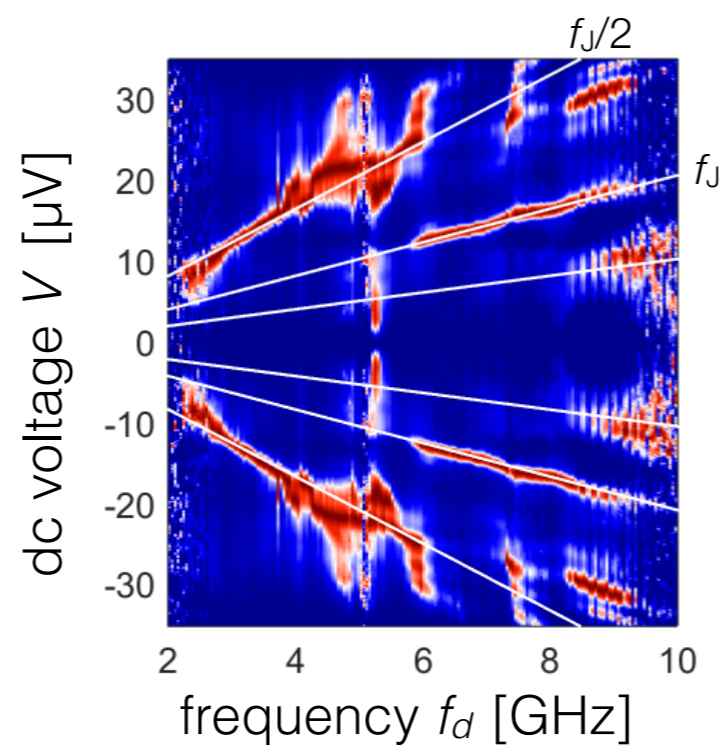
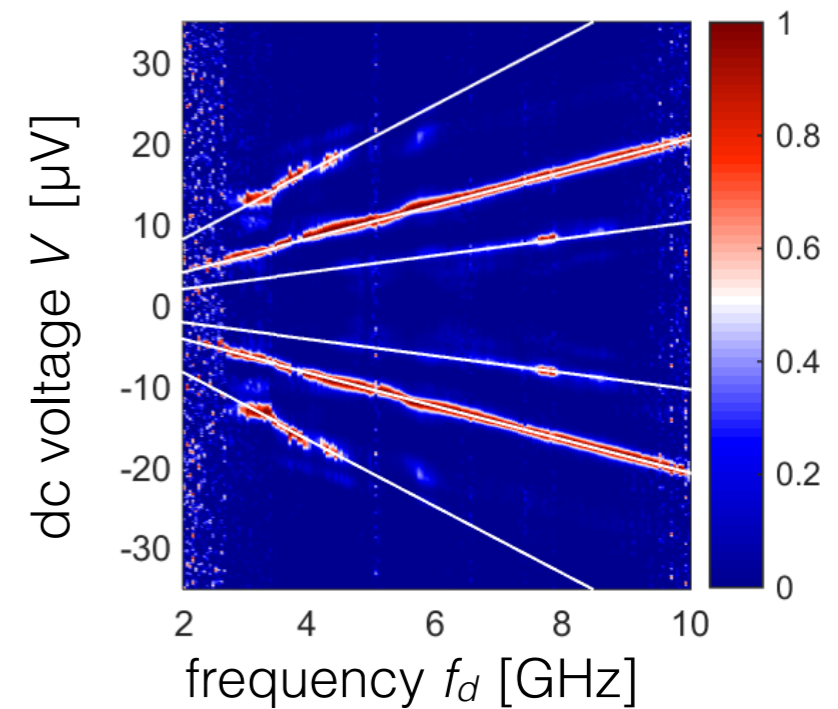
Deacon *et al.*, submitted, ArXiv 1603.09611 (2016)

Frequency dependence

Trivial QW

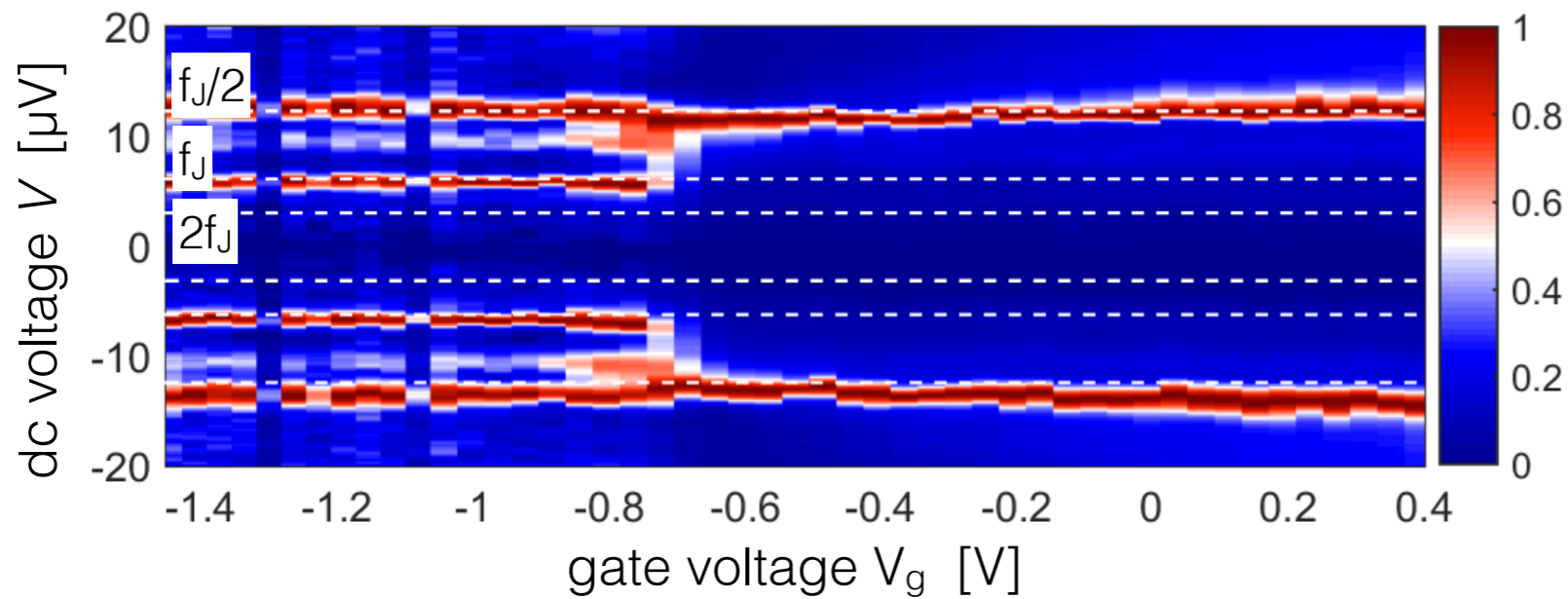


Topological QW

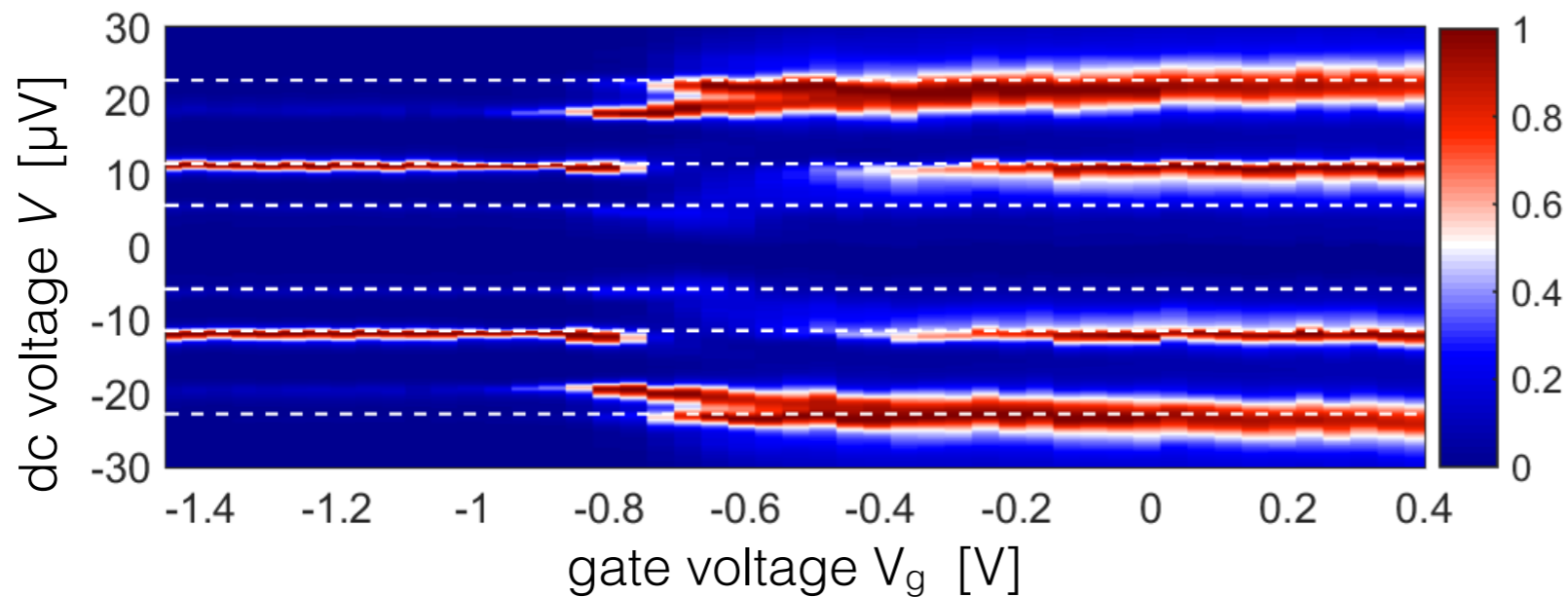
n- and QSH regime*p*-regime

- ▶ Stronger $f_J/2$ signal at low frequencies
- ▶ Relative intensities of $f_J/2$ and f_J depending on V_g

Gate voltage dependence

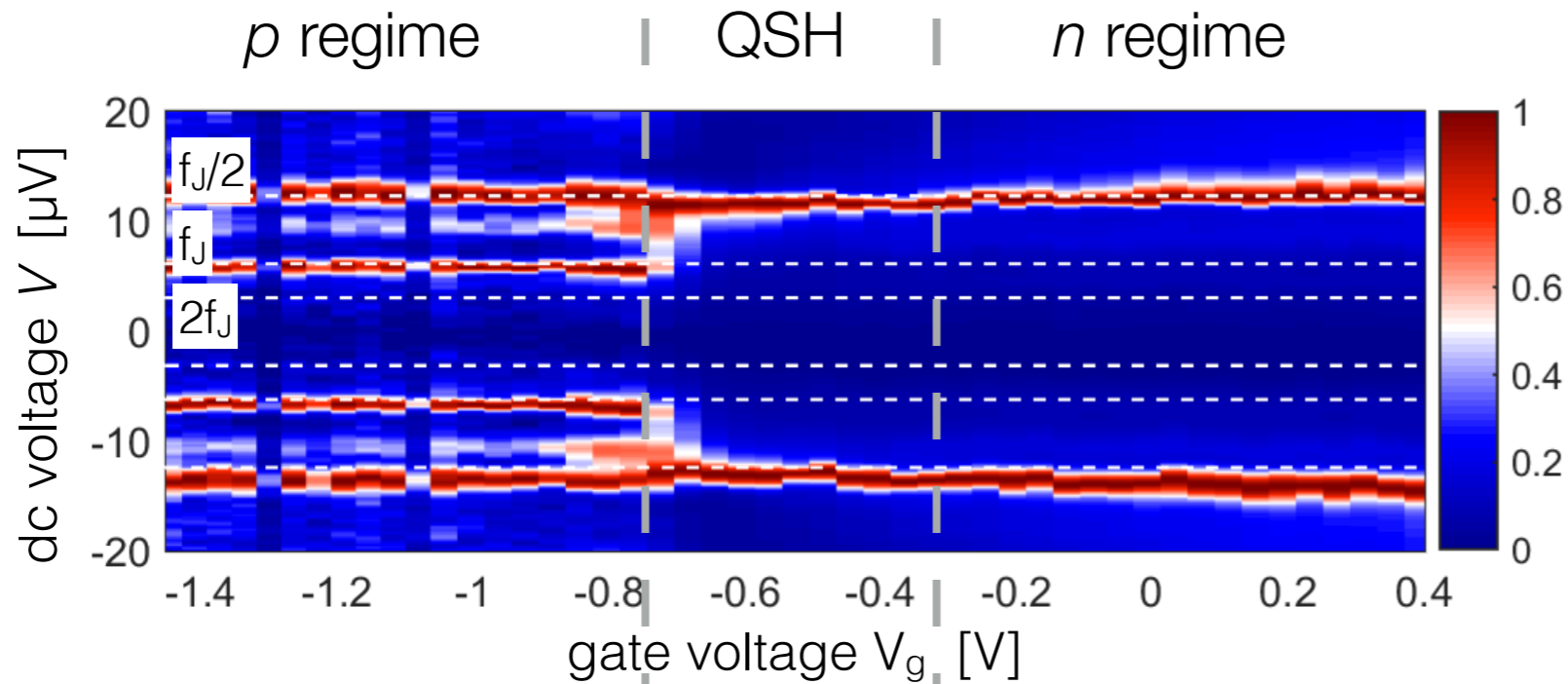


Low frequency
 $f_d = 3 \text{ GHz}$

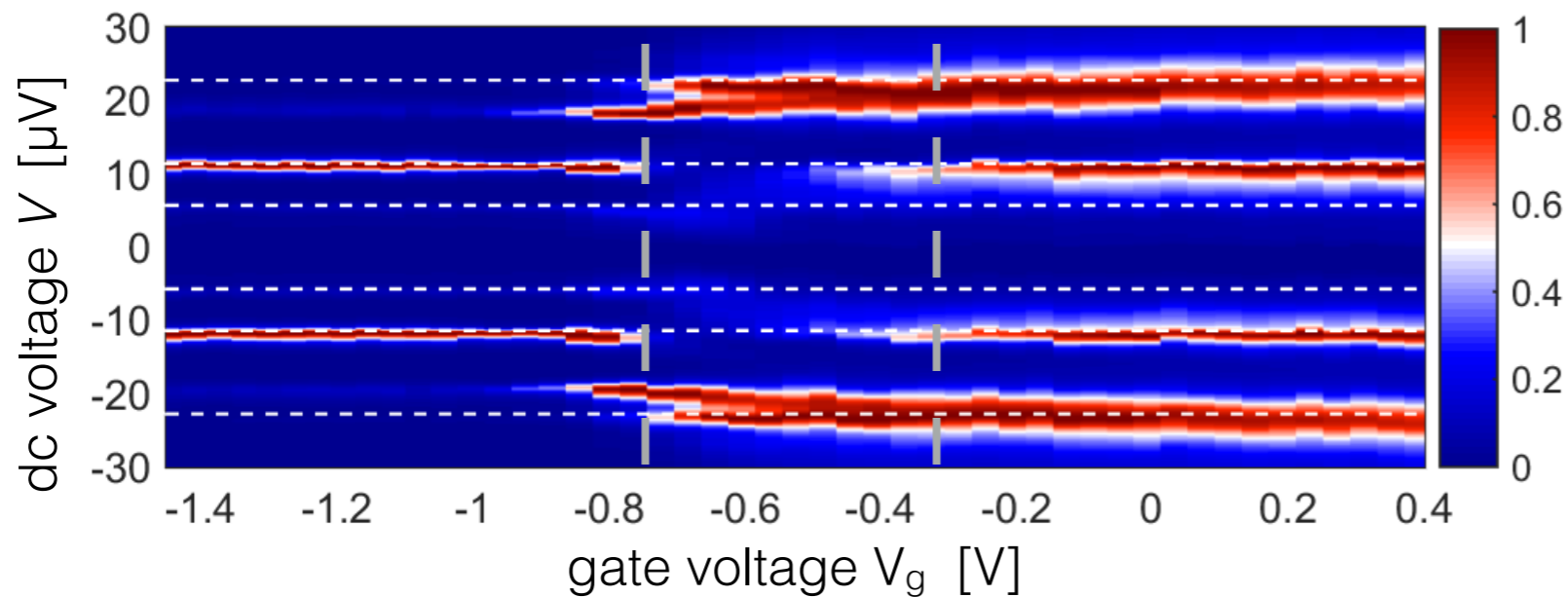


High frequency
 $f_d = 5.5 \text{ GHz}$

Gate voltage dependence



Low frequency
 $f_d = 3$ GHz



High frequency
 $f_d = 5.5$ GHz

Phase-locked motion

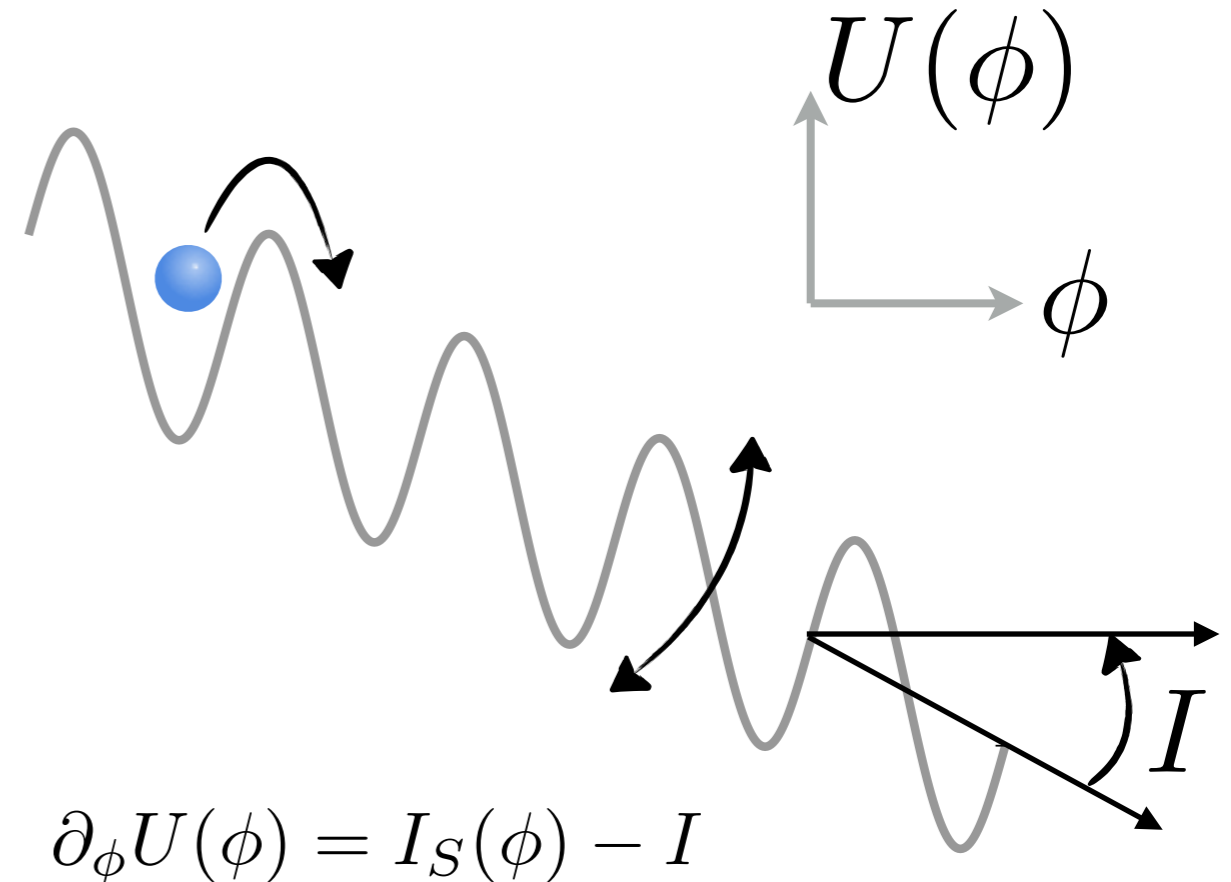
- ▷ phase dynamics (RSJ model)

$$\frac{d\phi}{dt} = \frac{2eV}{\hbar}$$

$$I = I_S(\phi) + \frac{\hbar}{2eR} \dot{\phi}$$

- ▷ motion locked to rf excitation

$$\frac{\Delta\phi}{\Delta t} = \frac{2\pi n}{1/f} \Rightarrow V_n = n \frac{hf}{2e}$$



Shapiro, PRL **11**, 80 (1963)

Russer, J. App. Phys. **43**, 2008 (1972)

Phase-locked motion

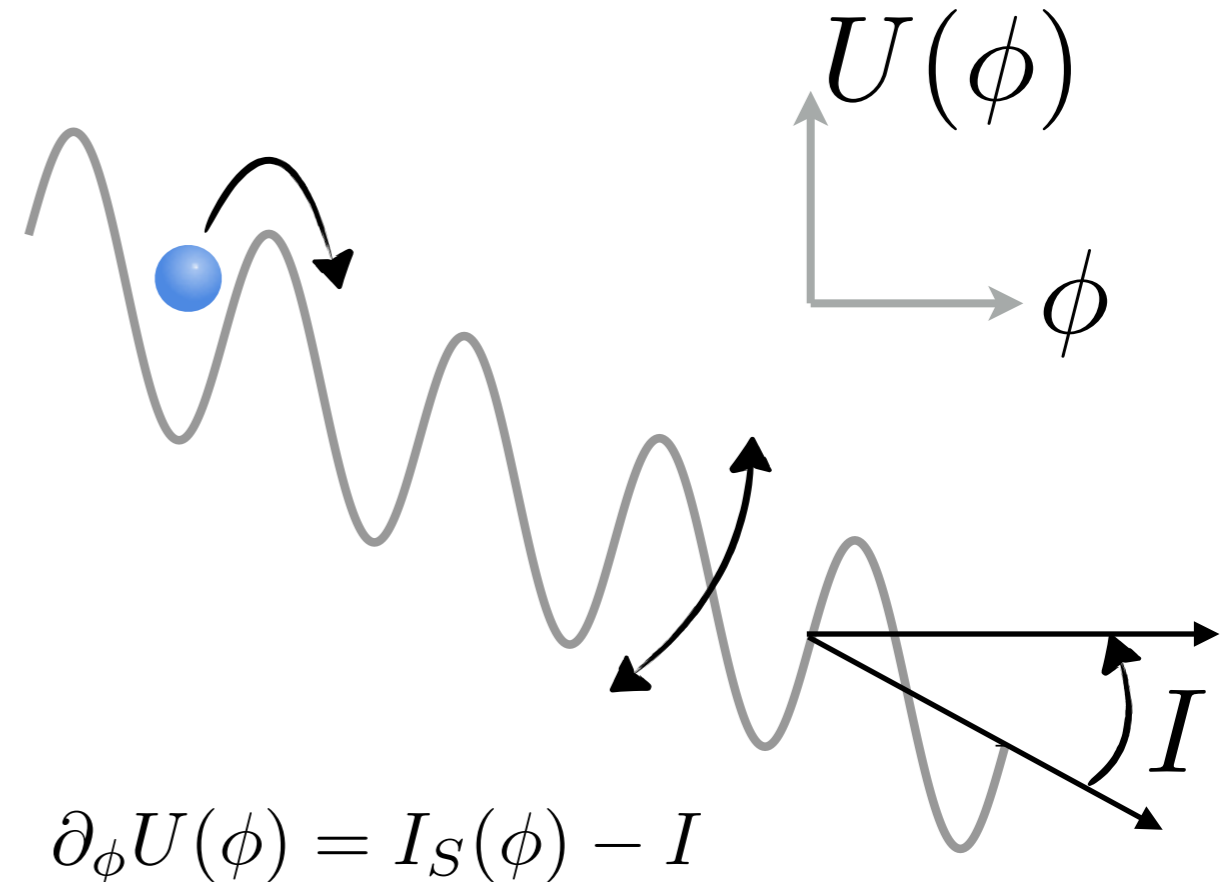
- ▷ phase dynamics (RSJ model)

$$\frac{d\phi}{dt} = \frac{2eV}{\hbar}$$

$$I = I_S(\phi) + \frac{\hbar}{2eR} \dot{\phi}$$

- ▷ motion locked to rf excitation

$$\frac{\Delta\phi}{\Delta t} = \frac{2\pi n}{1/f} \Rightarrow V_n = n \frac{hf}{2e}$$



Shapiro, PRL **11**, 80 (1963)

Russer, J. App. Phys. **43**, 2008 (1972)

Phase-locked motion

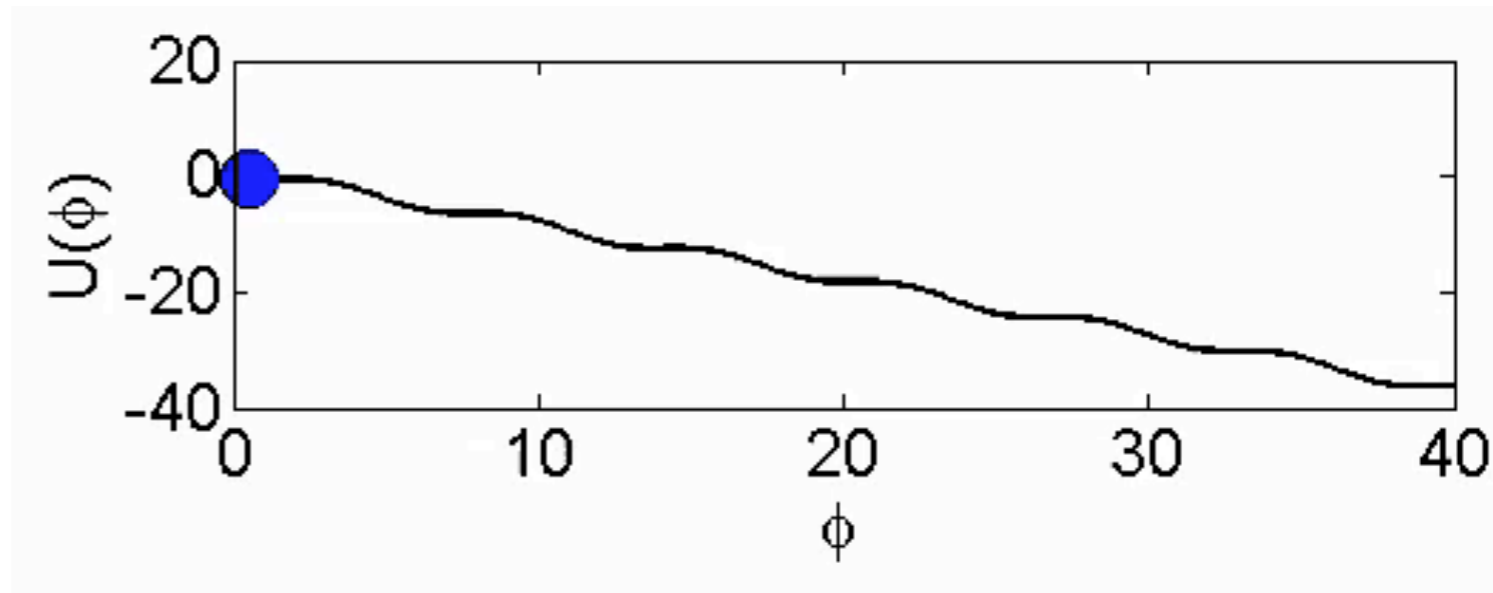
- ▷ phase dynamics (RSJ model)

$$\frac{d\phi}{dt} = \frac{2eV}{\hbar}$$

$$I = I_S(\phi) + \frac{\hbar}{2eR} \dot{\phi}$$

- ▷ motion locked to rf excitation

$$\frac{\Delta\phi}{\Delta t} = \frac{2\pi n}{1/f} \Rightarrow V_n = n \frac{hf}{2e}$$



Shapiro, PRL **11**, 80 (1963)

Russer, J. App. Phys. **43**, 2008 (1972)

Phase-locked motion

- ▷ phase dynamics (RSJ model)

$$\frac{d\phi}{dt} = \frac{2eV}{\hbar}$$

$$I = I_S(\phi) + \frac{\hbar}{2eR} \dot{\phi}$$

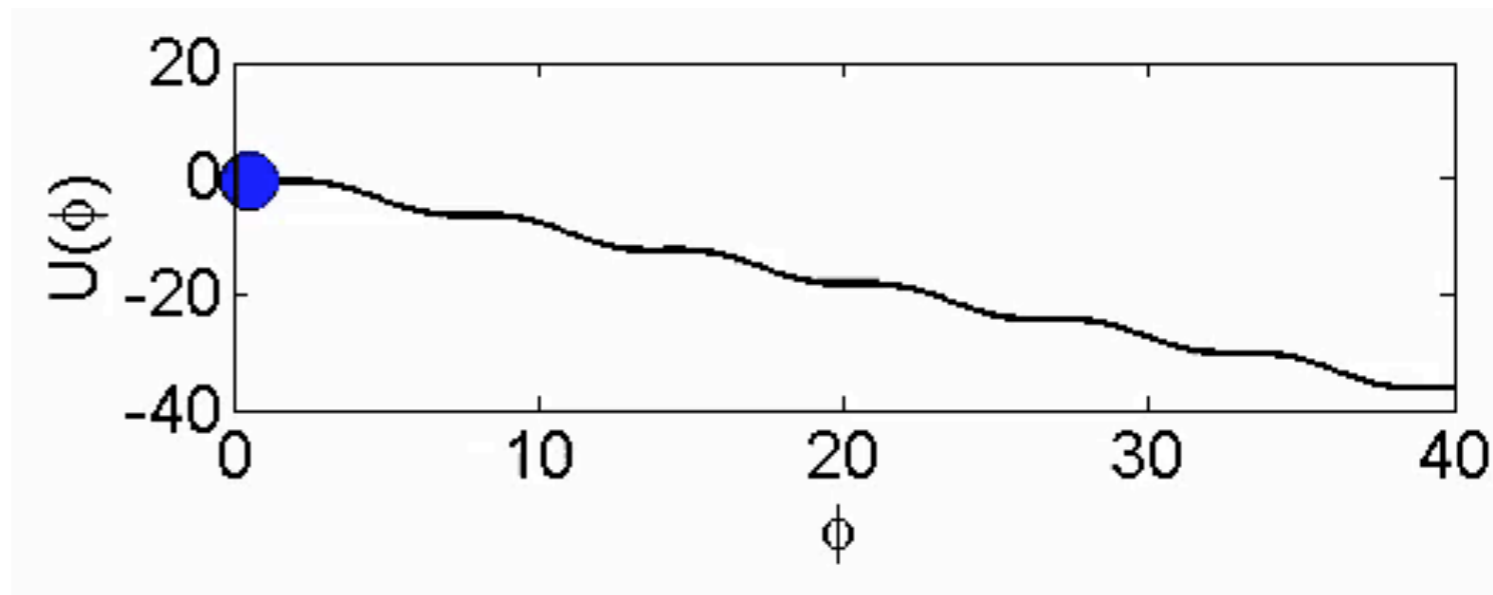
- ▷ motion locked to rf excitation

$$\frac{\Delta\phi}{\Delta t} = \frac{2\pi n}{1/f} \Rightarrow V_n = n \frac{hf}{2e}$$

4π-periodic supercurrent

- ▷ doubled steps
 $\sin \phi \rightarrow \sin \phi/2$
 $V_n \rightarrow V_{2n}$

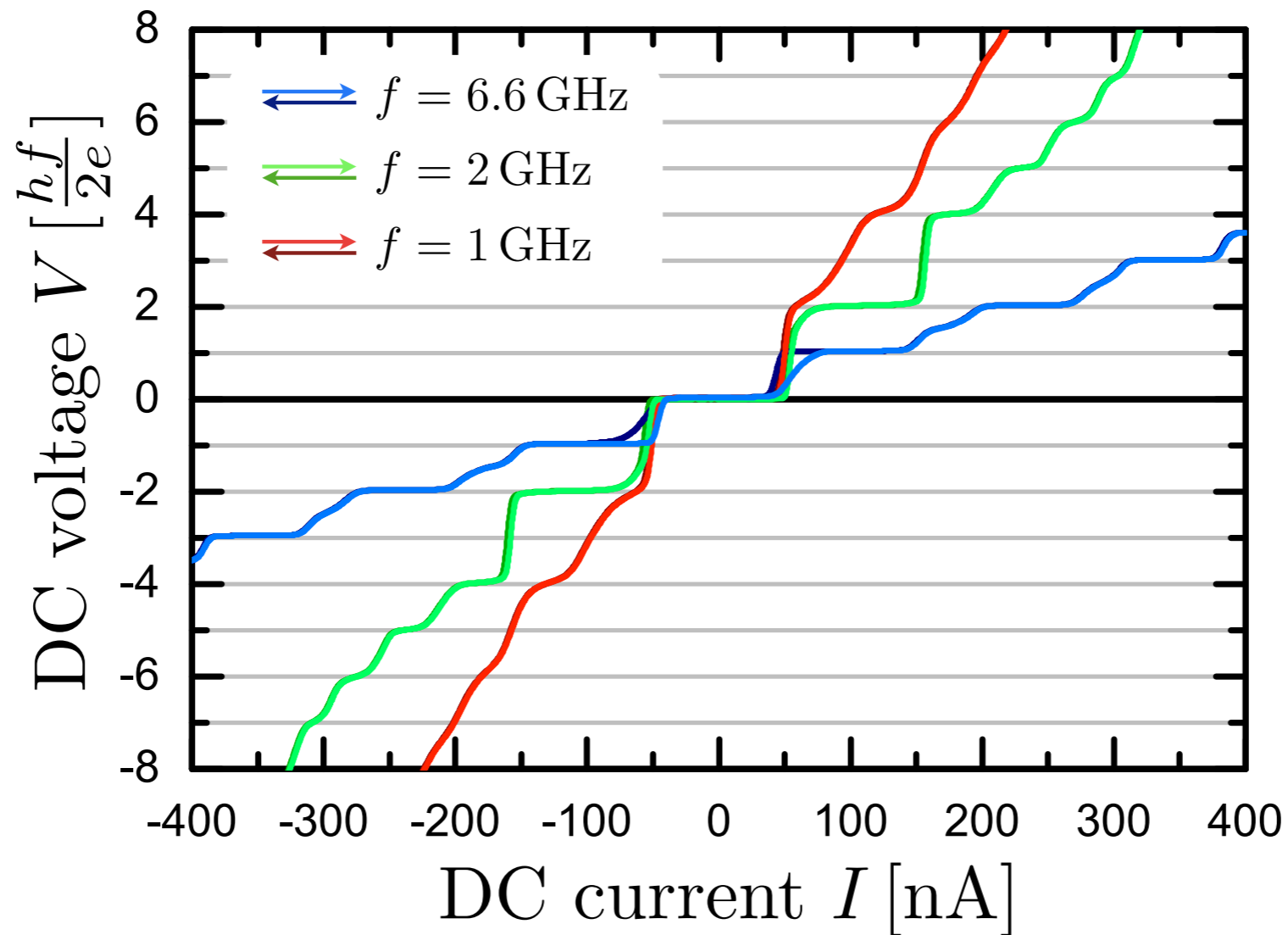
- ▷ mixture 2π/4π ?



Shapiro, PRL **11**, 80 (1963)

Russer, J. App. Phys. **43**, 2008 (1972)

Shapiro response : frequency



Shapiro response

- ▶ >12 steps visible
- ▶ weak hysteresis on 1st step

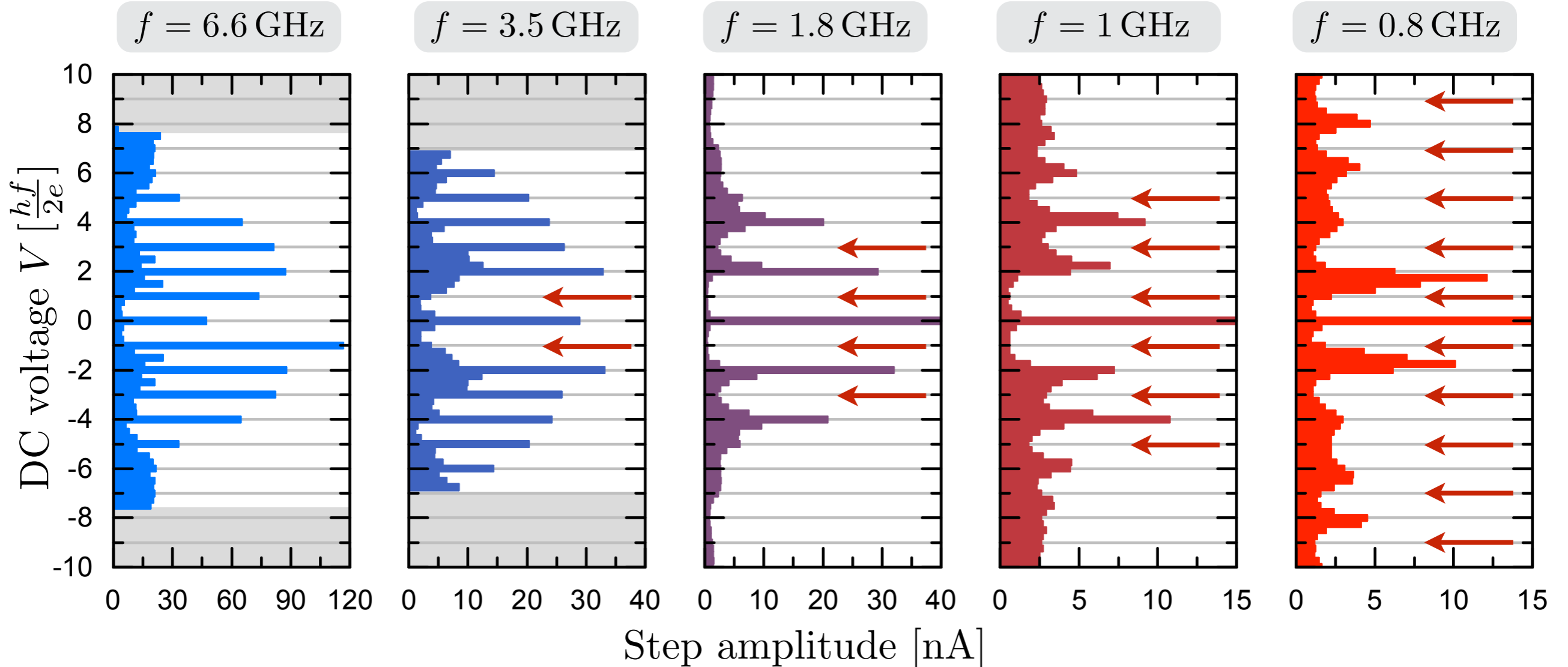
- ▶ multiple odd steps missing
- ▶ at low frequency $f \lesssim 4$ GHz

Rokhinson *et al.*, Nat. Phys. **86**, 146503 (2012)

Wiedenmann *et al.*, Nat. Comms **7**, 10303 (2016)

Bocquillon *et al.*, Nat. Nano, DOI: 10.1038/NNANO.2016.159

Shapiro response : frequency



Shapiro response

- ▷ >12 steps visible
- ▷ weak hysteresis on 1st step

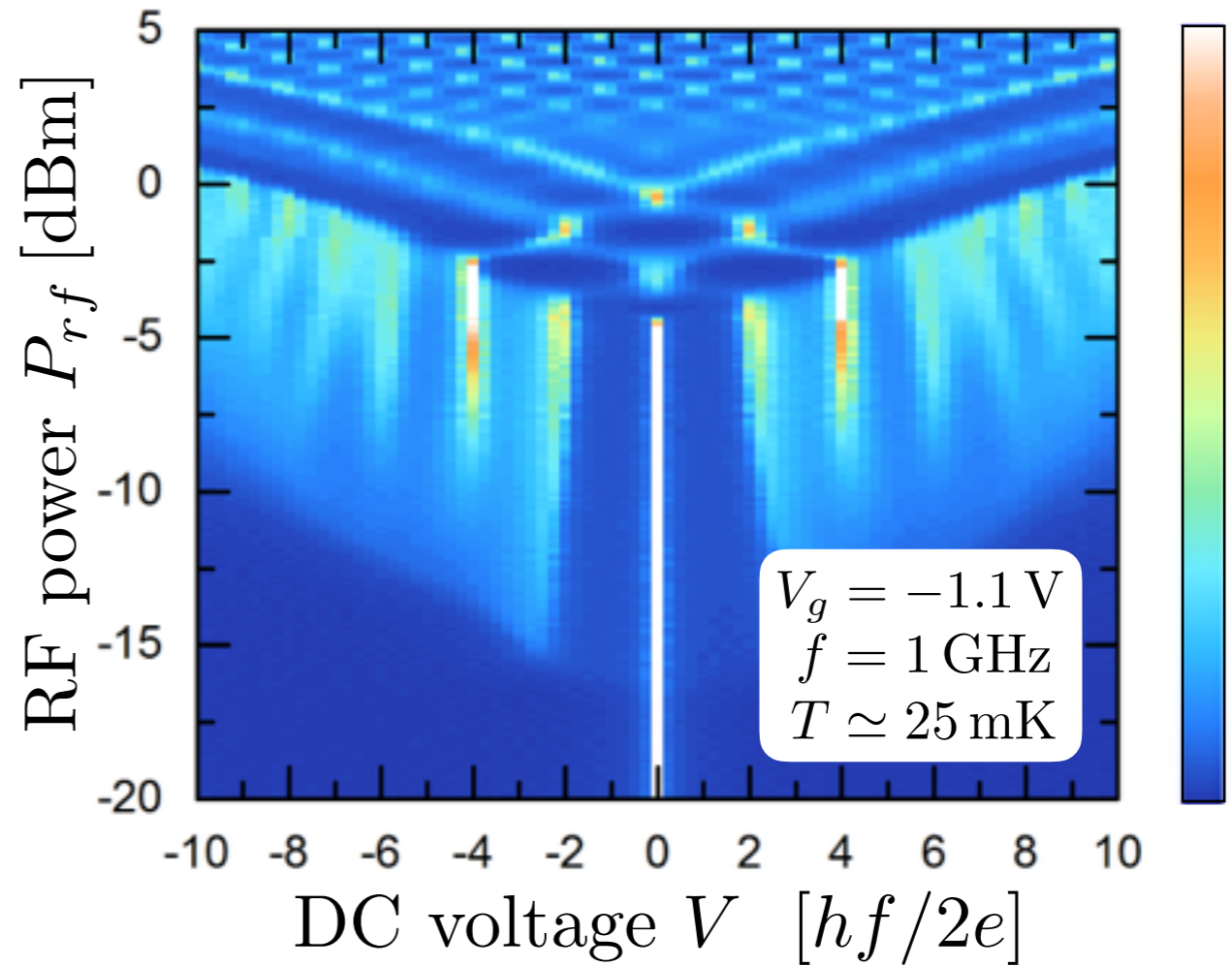
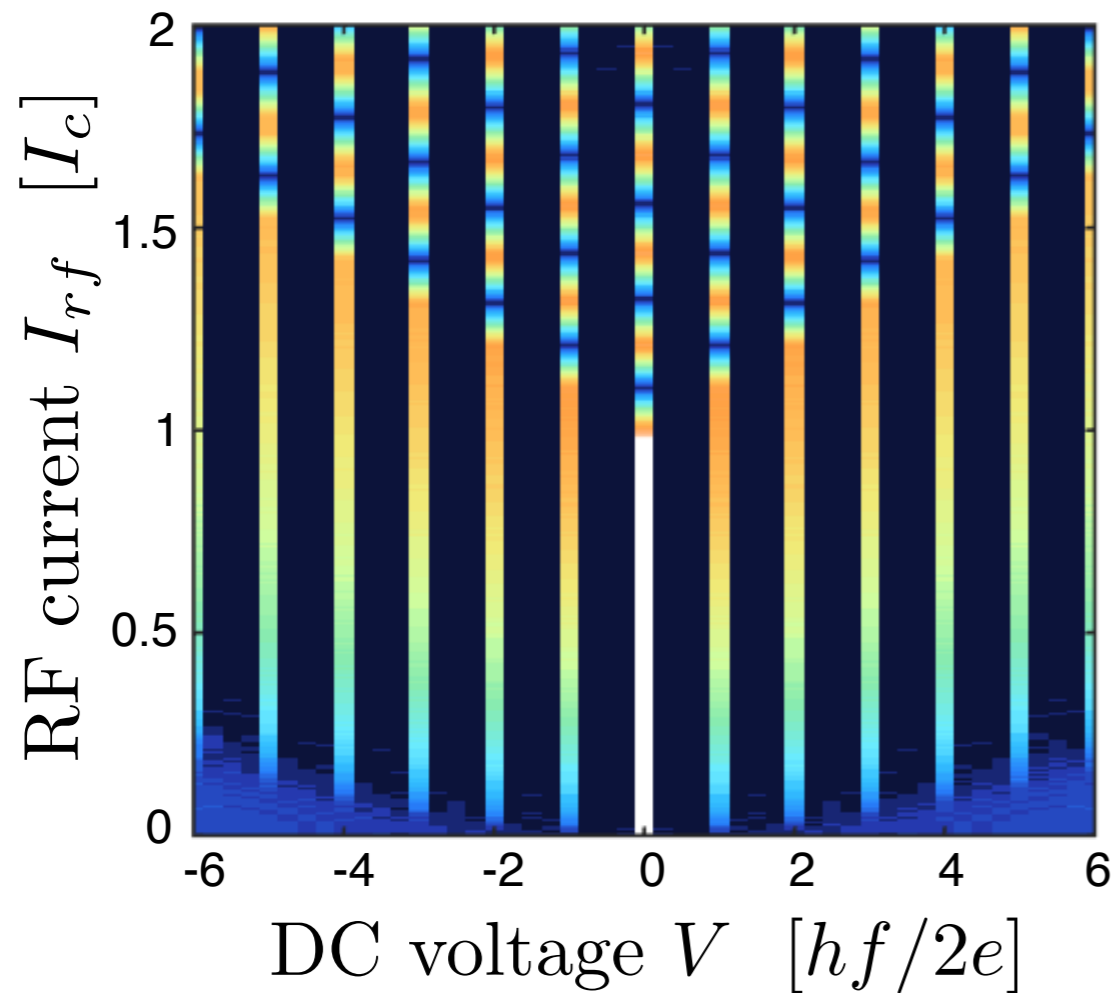
- ▷ multiple odd steps missing
- ▷ at low frequency $f \lesssim 4$ GHz

Rokhinson *et al.*, Nat. Phys. **86**, 146503 (2012)

Wiedenmann *et al.*, Nat. Comms **7**, 10303 (2016)

Bocquillon *et al.*, Nat. Nano, DOI: 10.1038/NNANO.2016.159

Shapiro response : power



Simulated response

- ▷ at low power : steps forming
- ▷ at high power : oscillatory pattern

Our device

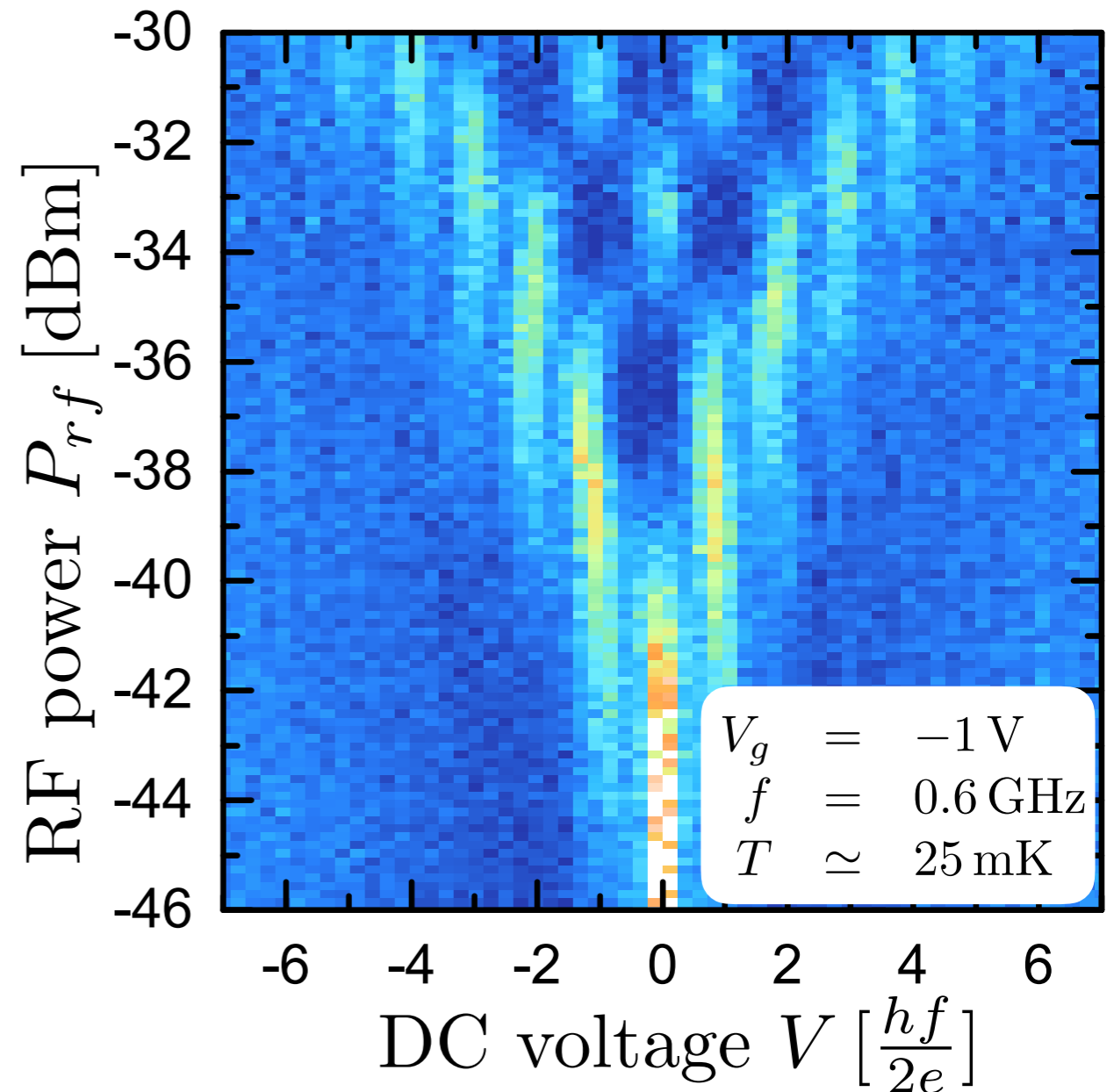
- ▷ missing $n=1,3,5$
- ▷ « dark fringes »

Non-topological QW

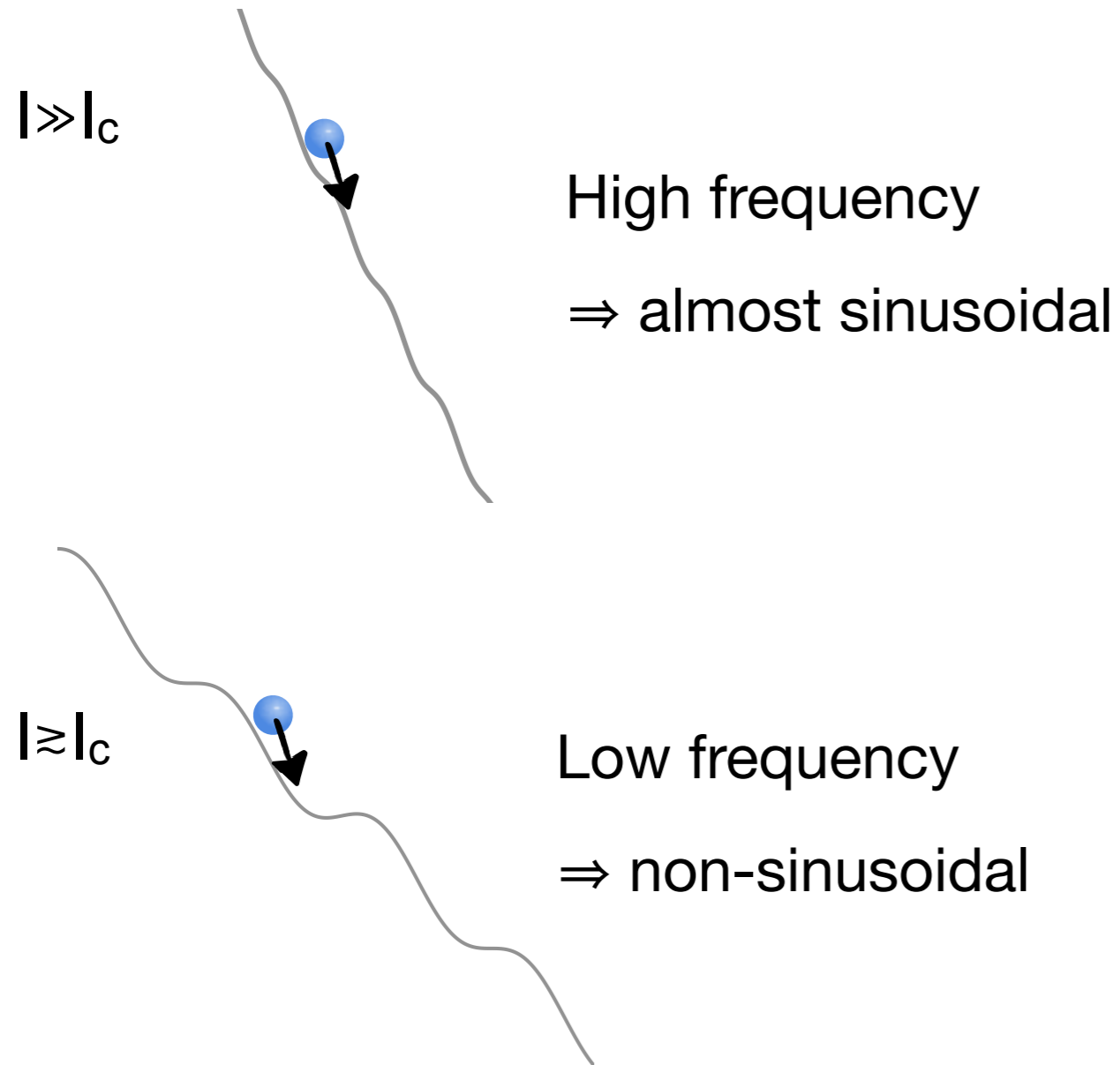
- ▶ narrow well (5 nm)
- ▶ no band inversion
- ▶ similar mobility $1.5 \cdot 10^5 \text{ cm}^2\text{V}^{-1}\text{s}^{-1}$

Shapiro steps

- ▶ no missing steps
- ▶ n -, p - regimes and gap verified
- ▶ down to $f = 0.6 \text{ GHz}$

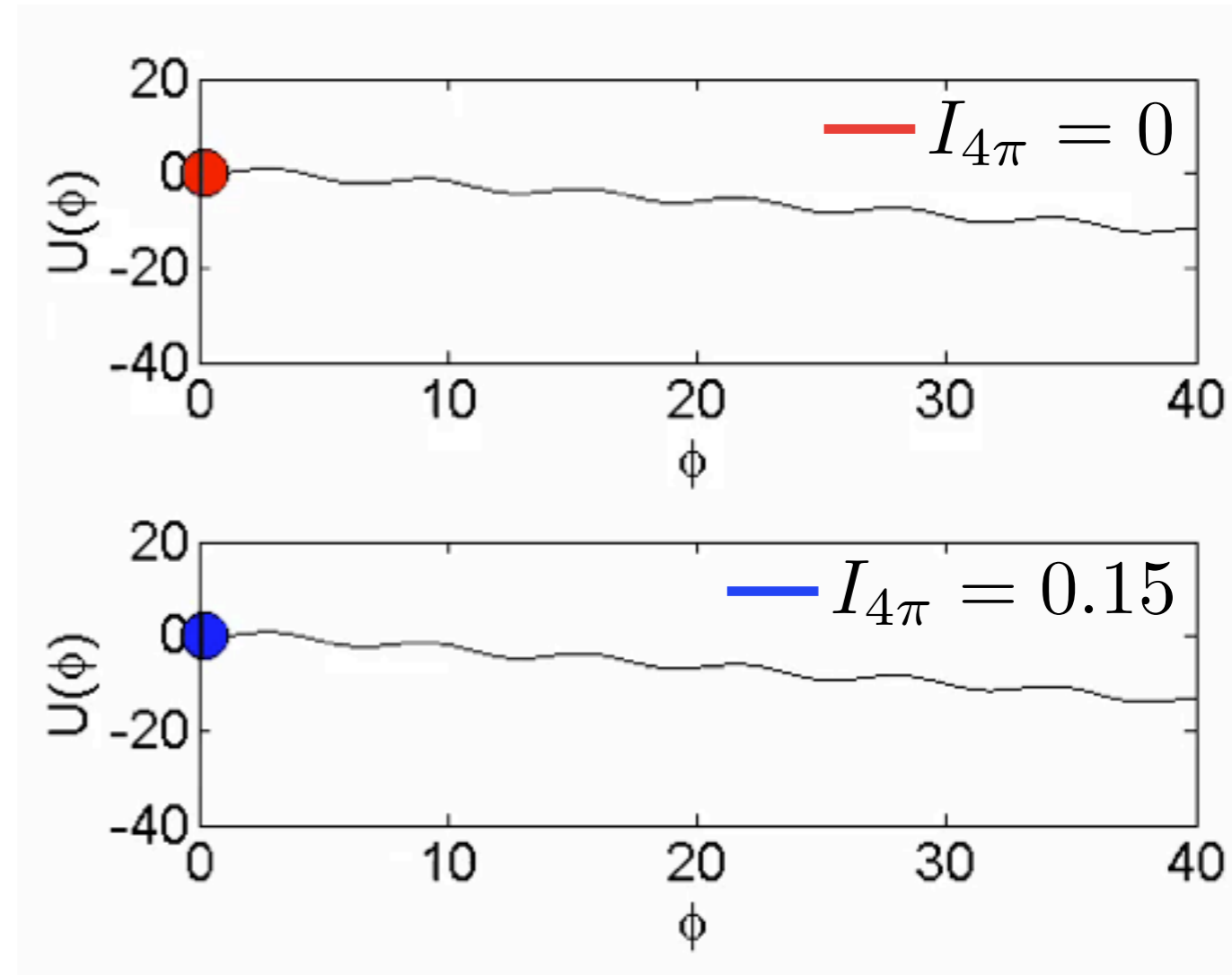
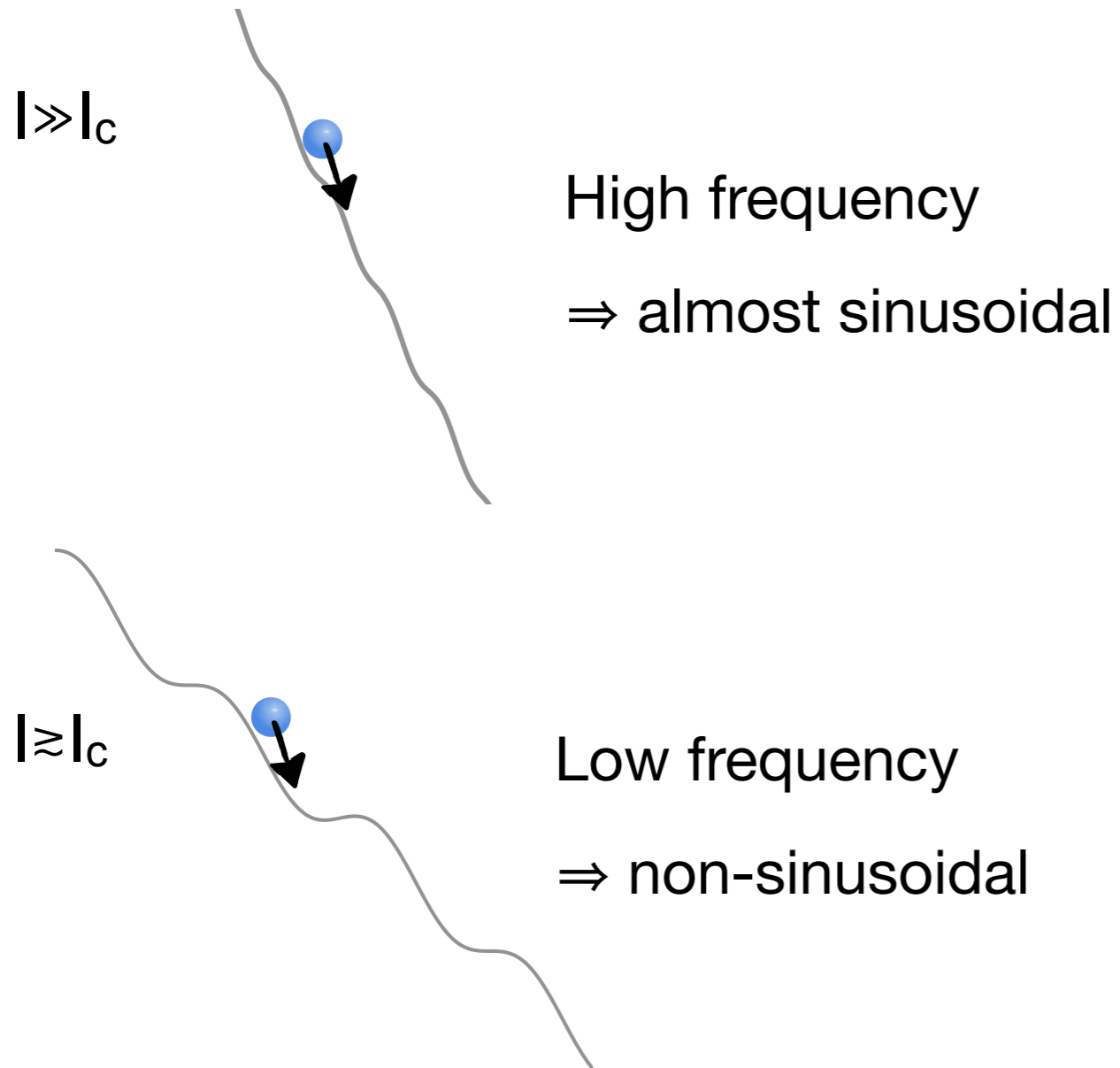


RSJ frequency dependence

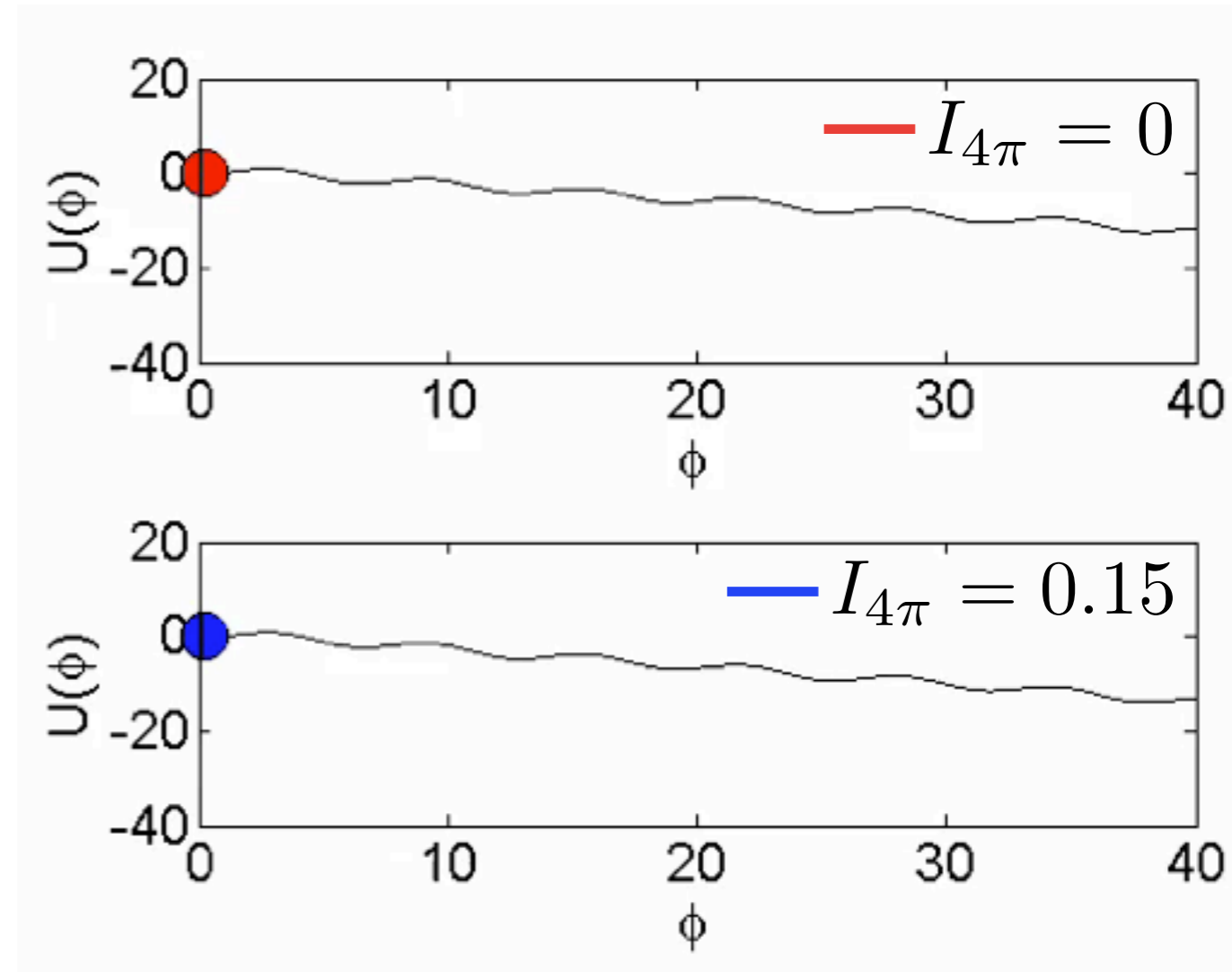
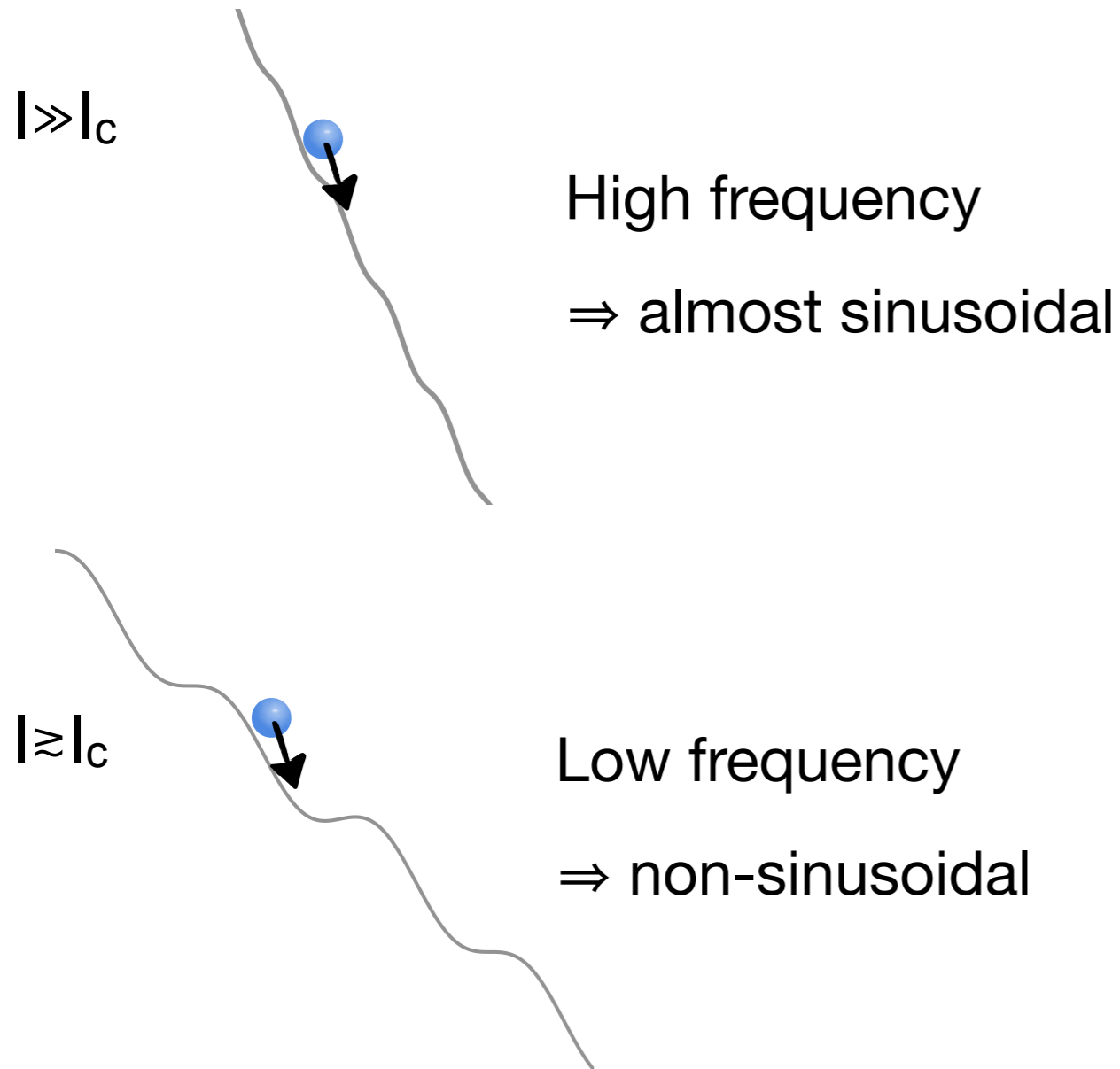


Domínguez *et al.*, PRB **86**, 146503 (2012)

RSJ frequency dependence

Domínguez *et al.*, PRB **86**, 146503 (2012)

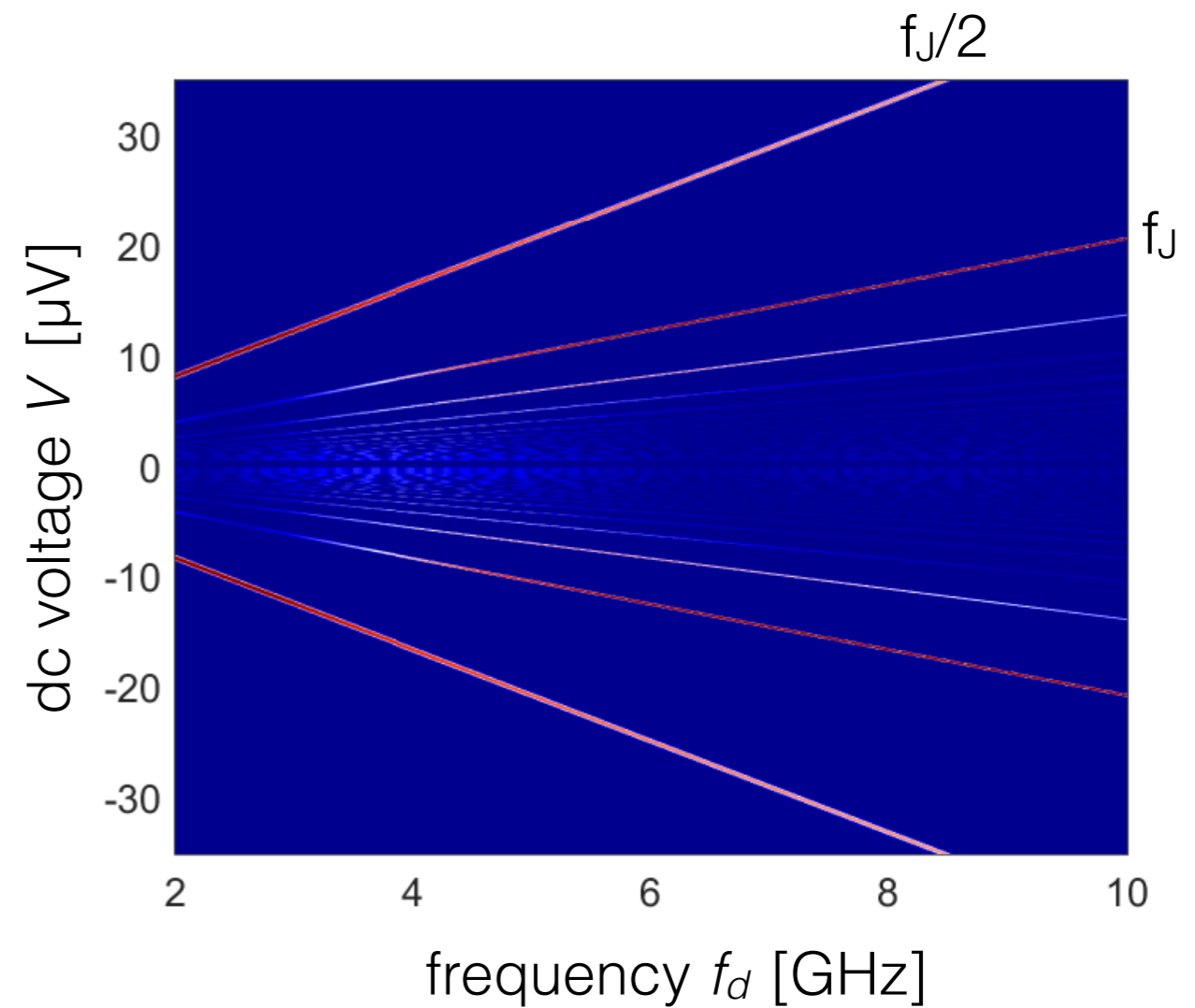
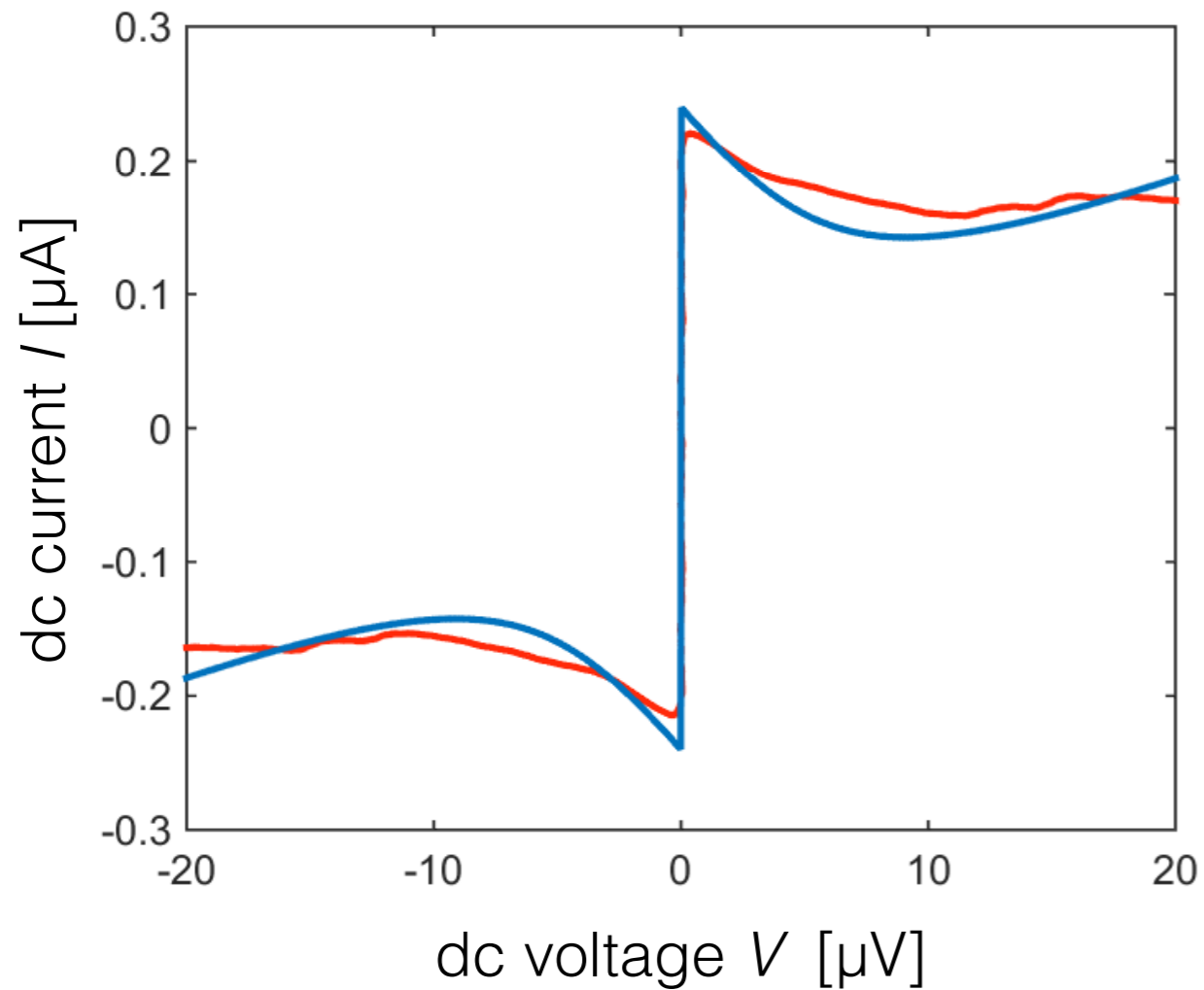
RSJ frequency dependence



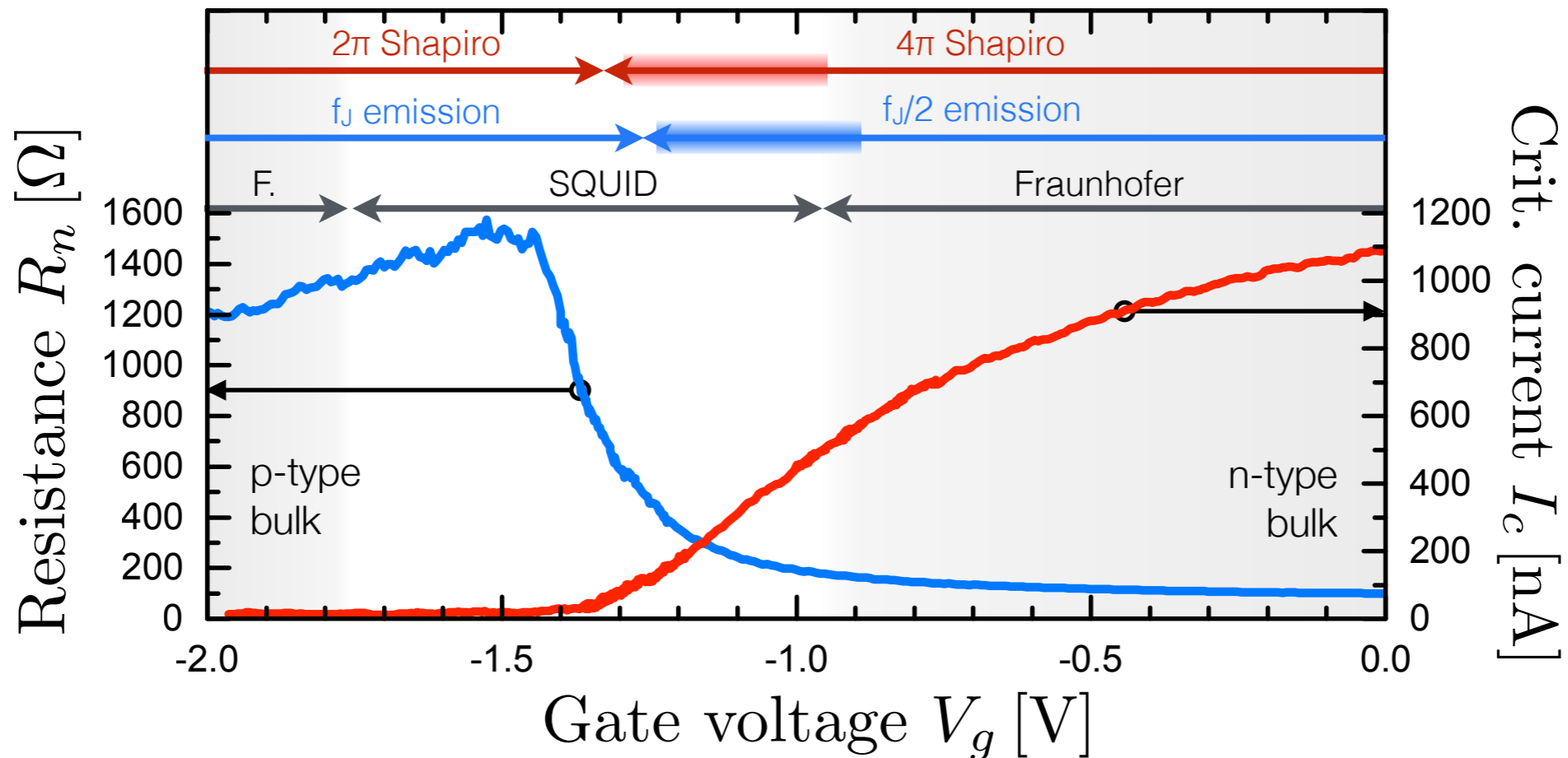
⇒ crossover $f_{4\pi}$ yields : 1-3 modes

⇒ no Landau-Zener transitions ?

Domínguez *et al.*, PRB **86**, 146503 (2012)



Theory : F. Domínguez & E. M. Hankiewicz



Fractional Josephson effect

- ▶ even sequence of Shapiro steps
- ▶ emission at $f_J/2$
- ▶ Landau-Zener transitions unlikely
- ▶ edge currents (SQUID pattern)
- ▶ contribution: 1-3 modes
- ▶ coexistence with conduction band ?
- ▶ discrepancy with R_n ?

Dai et al., PRB **77**, 125319 (2008)
Hart et al., Nat. Phys. **10**, 638 (2014)

Spectroscopy of Andreev bound states

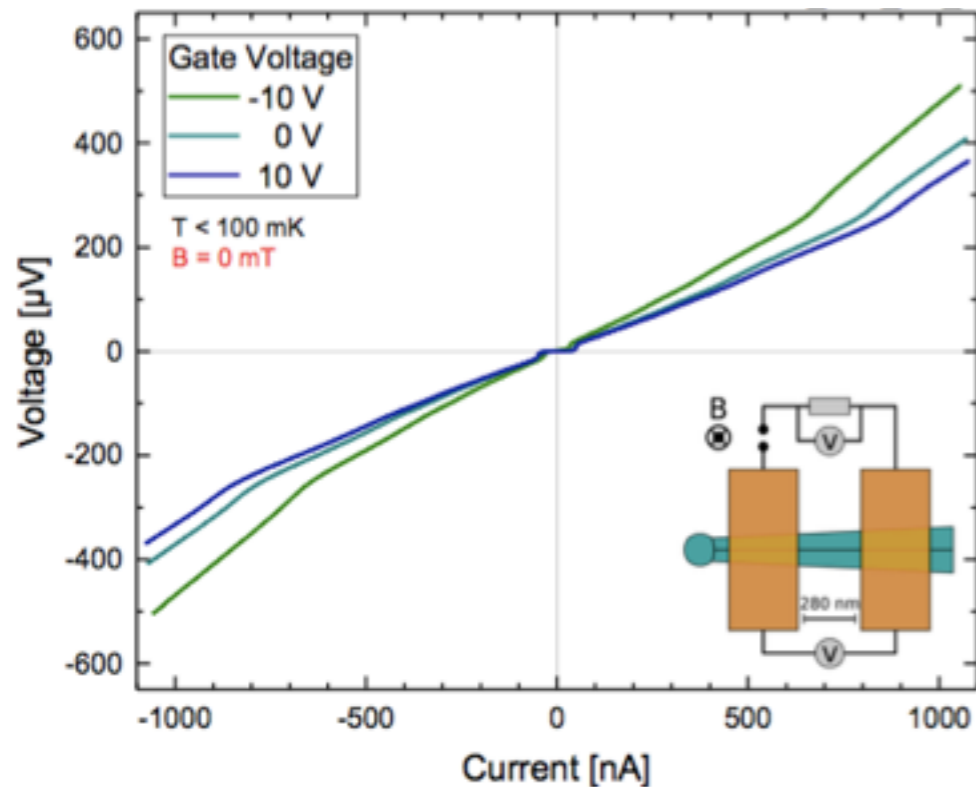
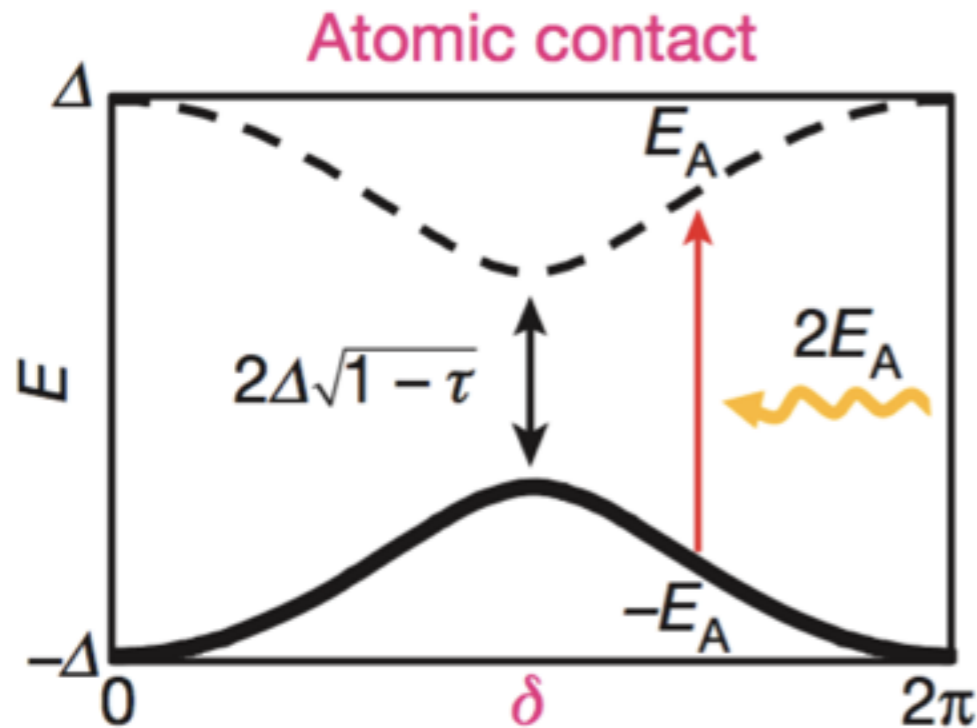
- ▷ tunneling DOS
- ▷ absorption spectroscopy
- ▷ SN junctions

⇒ towards Majorana qu-bits

Pillet *et al.*, Nature Phys. **6**, 965 (2010)
 Bretheau *et al.*, Nature **499**, 312 (2013)
 Astafiev *et al.*, Science **327** 840 (2010)
 Peng *et al.*, arXiv 1604.04287 (2016)

Other HgTe systems

- ▷ HgTe nanowires
- ▷ QSH, QH, QAH, Weyl



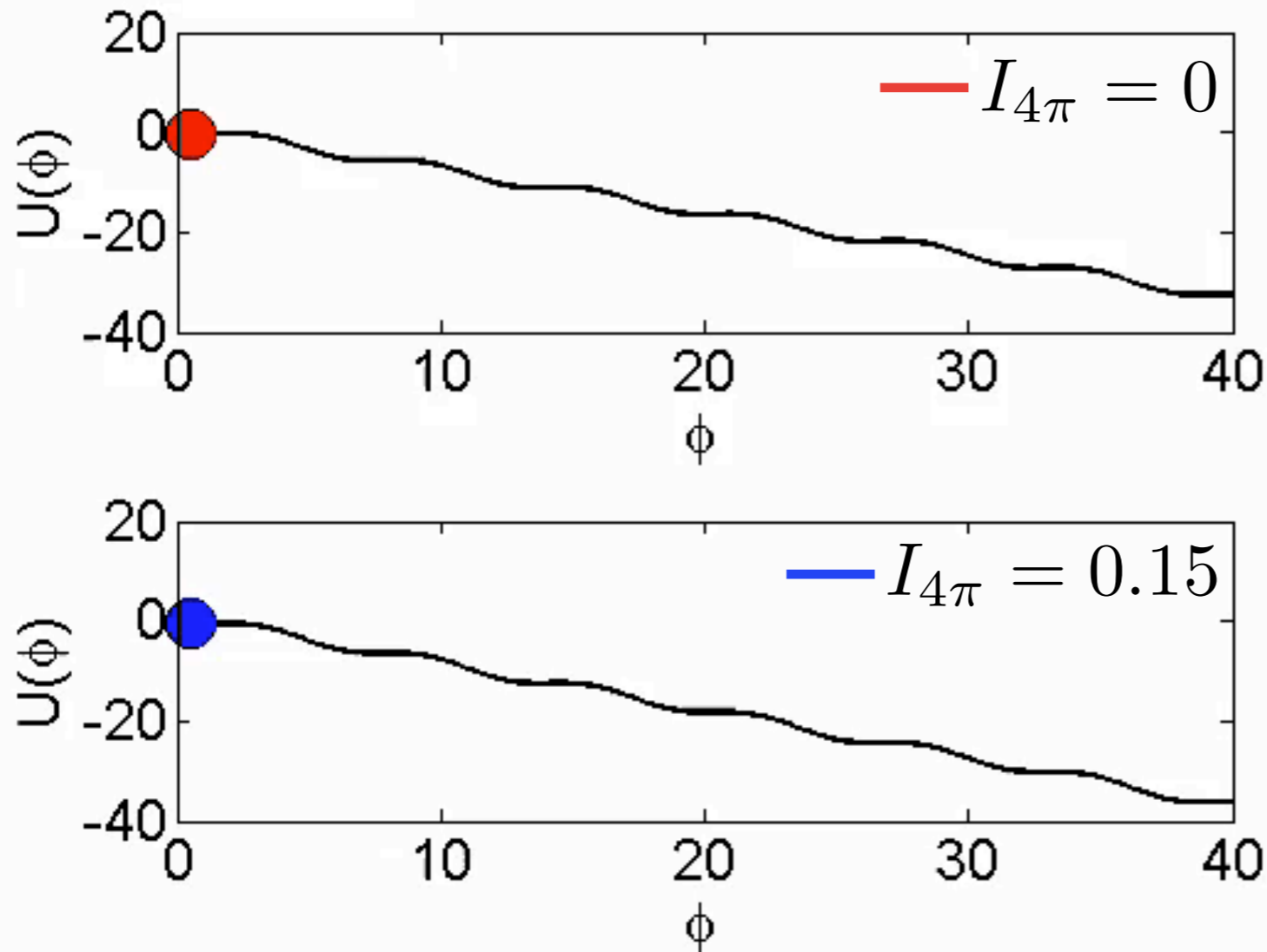
Thank you for your attention !

Open positions!

Post-docs

PhD students





RSJ simulations

- ▷ anharmonic motion
- ▷ doubled step due to 4π
- ▷ transition at $f \sim f_{4\pi} = \frac{2eR_n I_{4\pi}}{h}$

4π -supercurrent

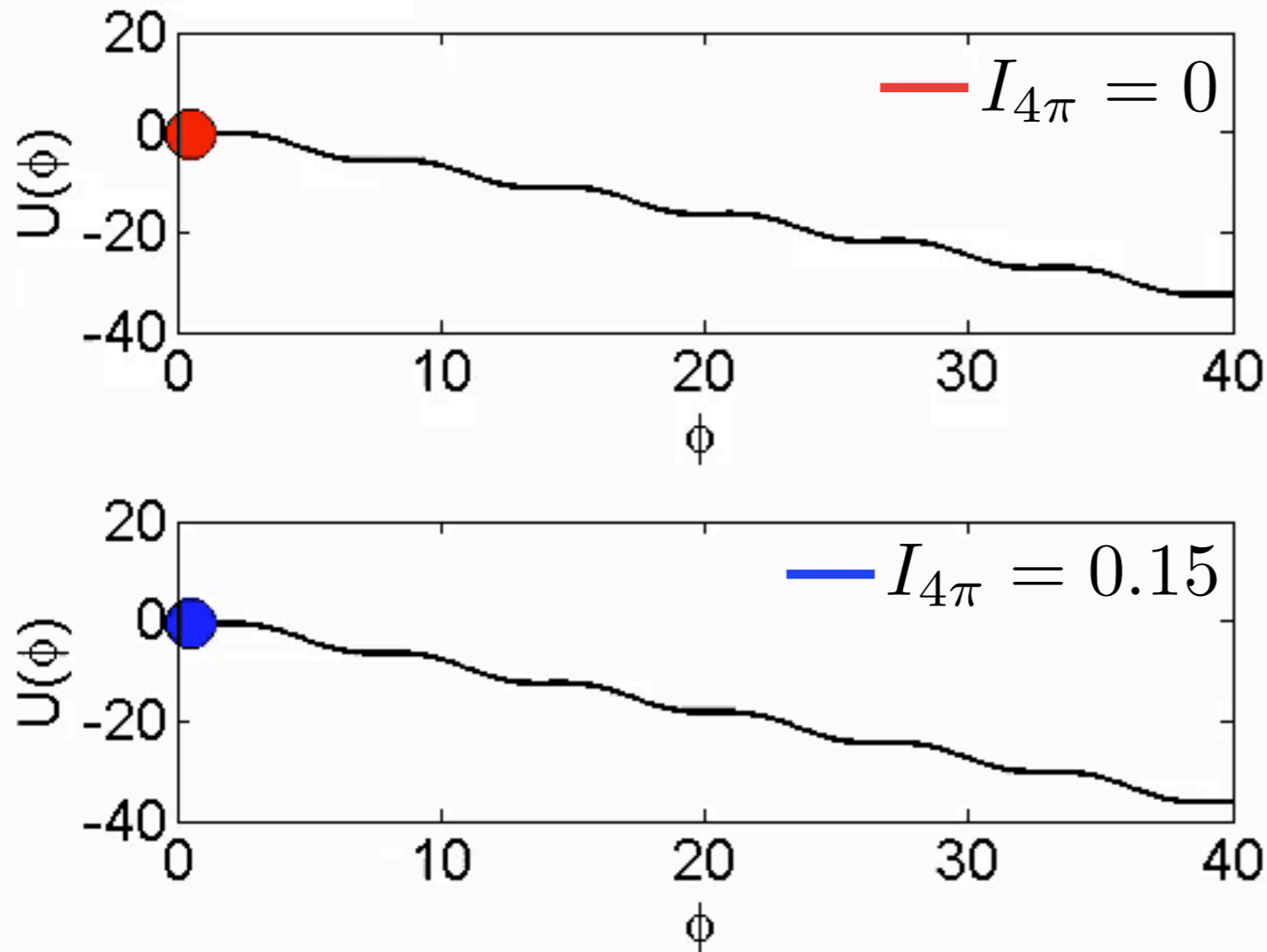
- ▷ ≈ 20 nA at $V_g = -1.1$ V
- ▷ 1-2 modes

Landau-Zener transitions

- ▷ stronger at high frequency

Dominguez *et al.*, PRB **86**, 146503 (2012)

Wiedenmann *et al.*, Nat. Comms **7**, 10303 (2016)



RSJ simulations

- ▷ anharmonic motion
- ▷ doubled step due to 4π
- ▷ transition at $f \sim f_{4\pi} = \frac{2eR_n I_{4\pi}}{h}$

4π -supercurrent

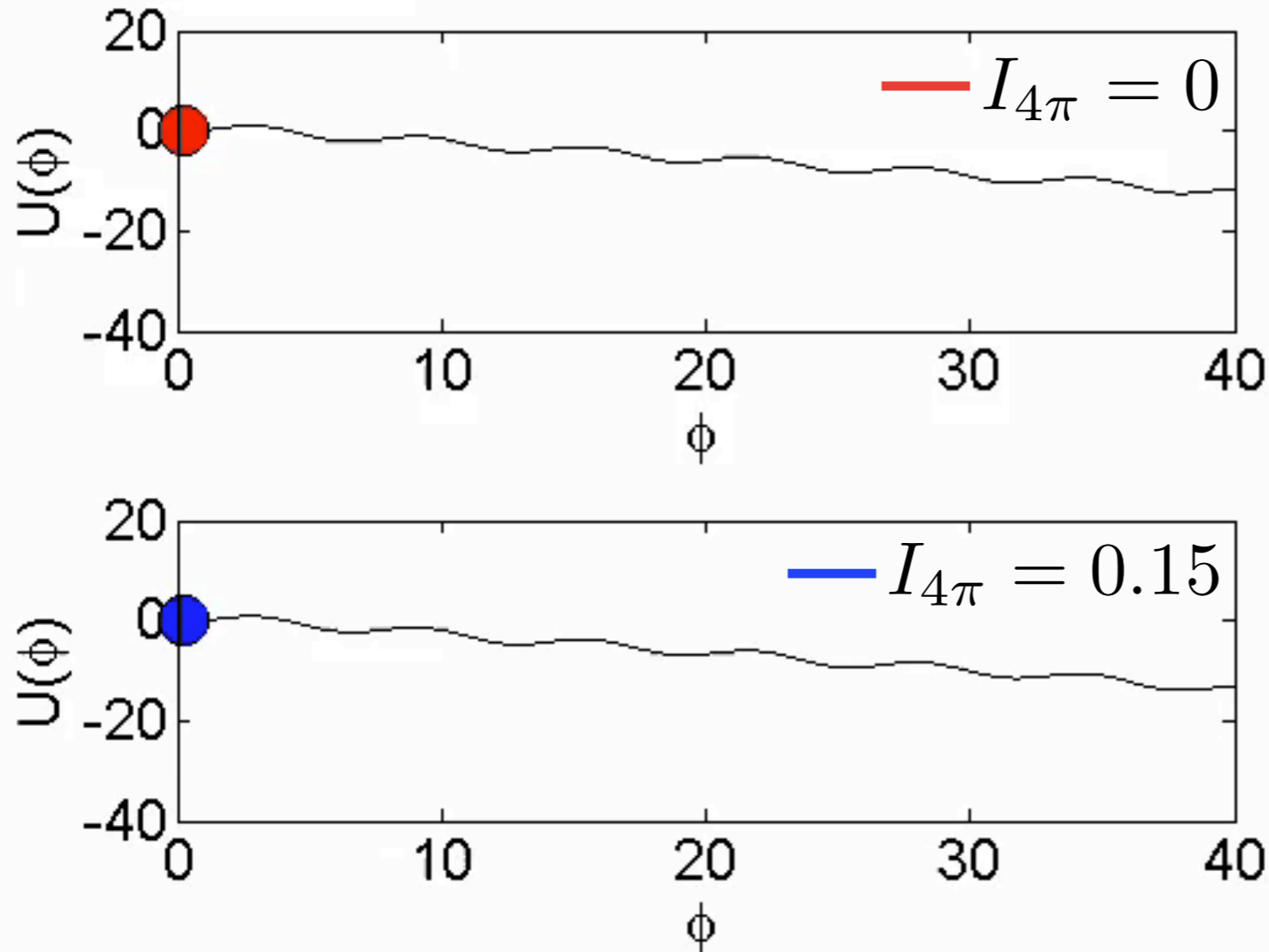
- ▷ ≈ 20 nA at $V_g = -1.1$ V
- ▷ 1-2 modes

Landau-Zener transitions

- ▷ stronger at high frequency

Dominguez *et al.*, PRB **86**, 146503 (2012)

Wiedenmann *et al.*, Nat. Comms **7**, 10303 (2016)



RSJ simulations

- ▷ anharmonic motion
- ▷ doubled step due to 4π
- ▷ transition at $f \sim f_{4\pi} = \frac{2eR_n I_{4\pi}}{h}$

4π -supercurrent

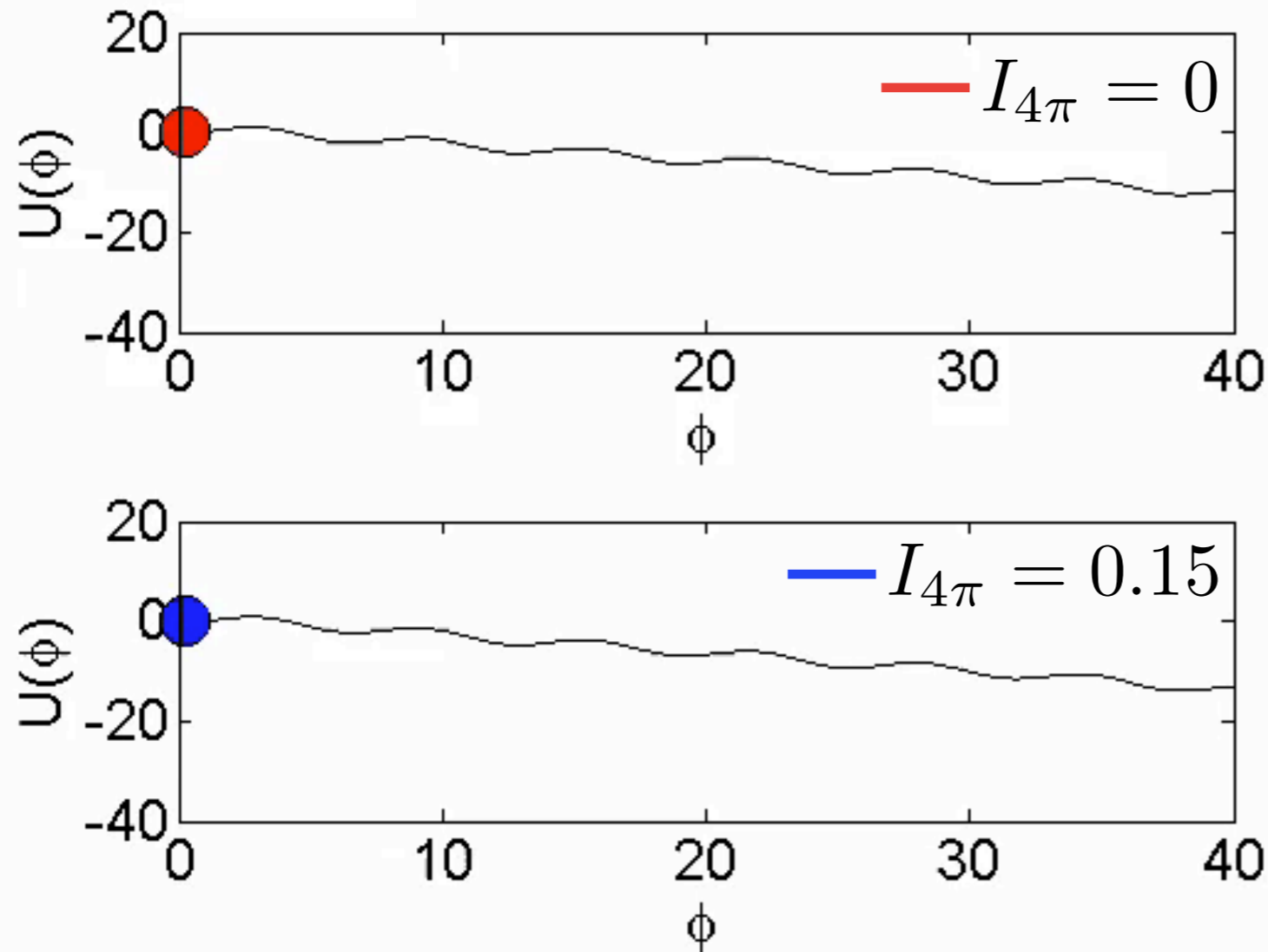
- ▷ ≈ 20 nA at $V_g = -1.1$ V
- ▷ 1-2 modes

Landau-Zener transitions

- ▷ stronger at high frequency

Dominguez *et al.*, PRB **86**, 146503 (2012)

Wiedenmann *et al.*, Nat. Comms **7**, 10303 (2016)



RSJ simulations

- ▷ anharmonic motion
- ▷ doubled step due to 4π
- ▷ transition at $f \sim f_{4\pi} = \frac{2eR_n I_{4\pi}}{h}$

4π -supercurrent

- ▷ ≈ 20 nA at $V_g = -1.1$ V
- ▷ 1-2 modes

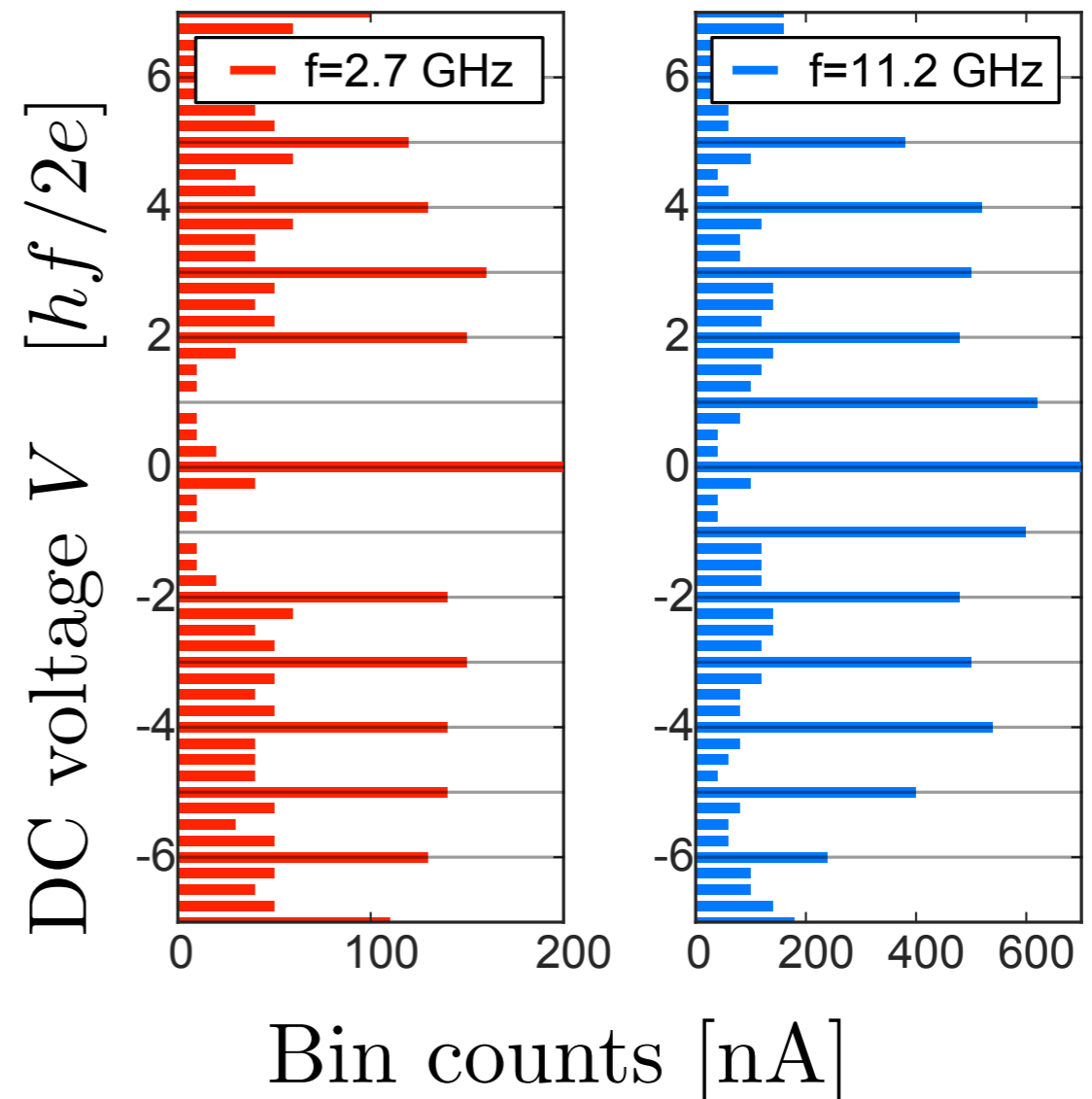
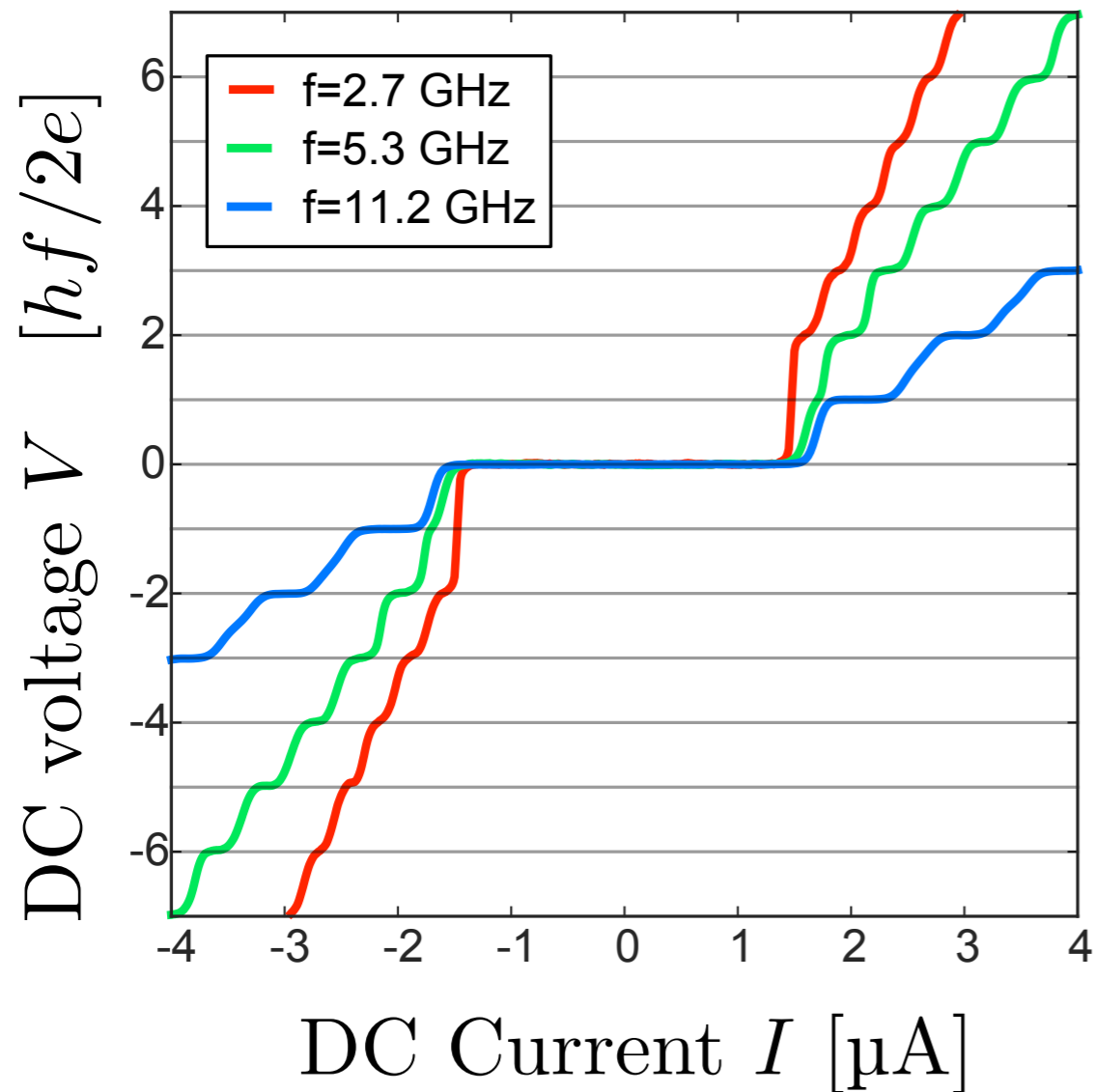
Landau-Zener transitions

- ▷ stronger at high frequency

Dominguez *et al.*, PRB **86**, 146503 (2012)

Wiedenmann *et al.*, Nat. Comms **7**, 10303 (2016)

Shapiro steps



High frequency

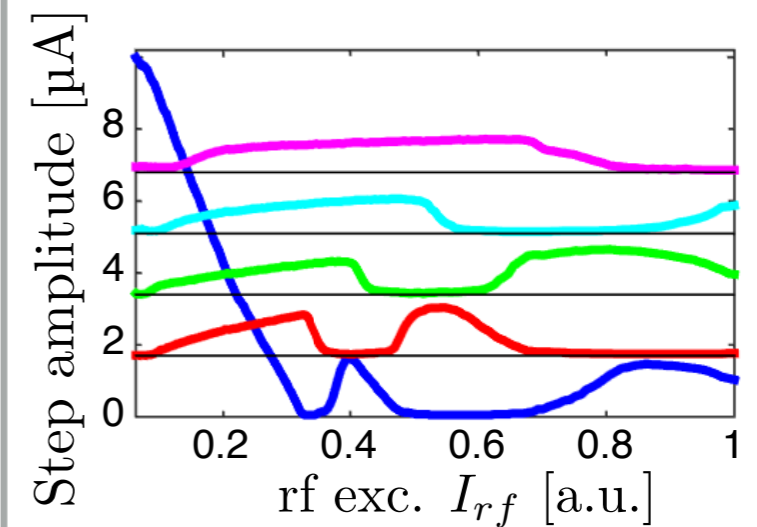
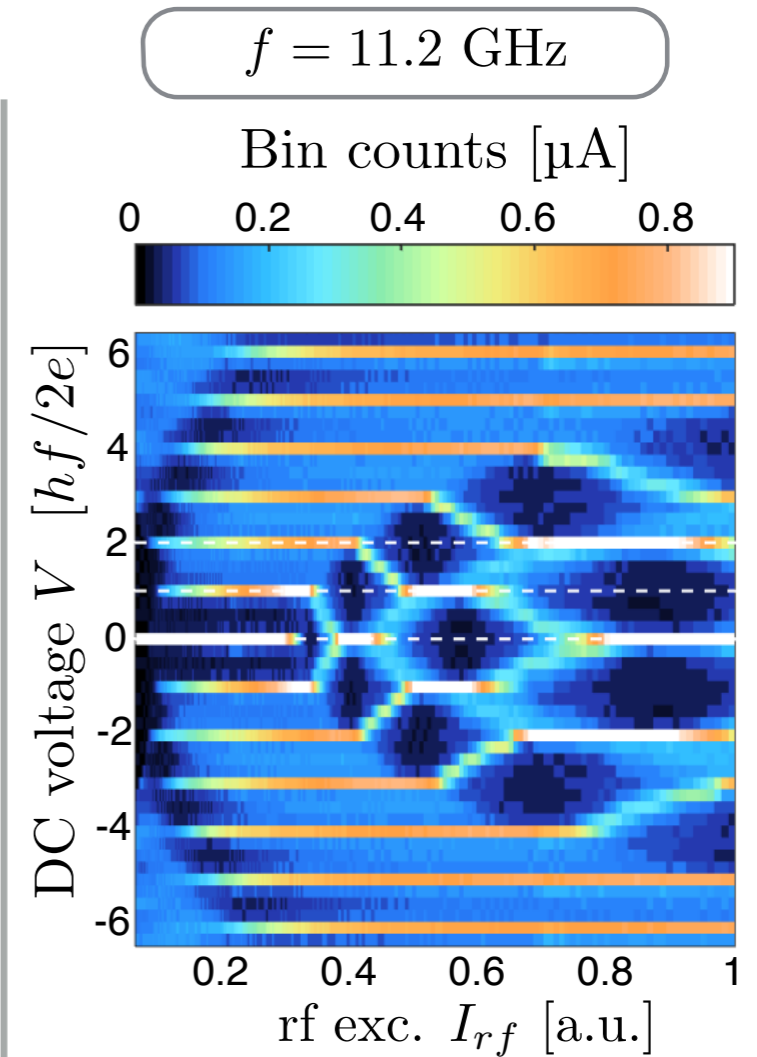
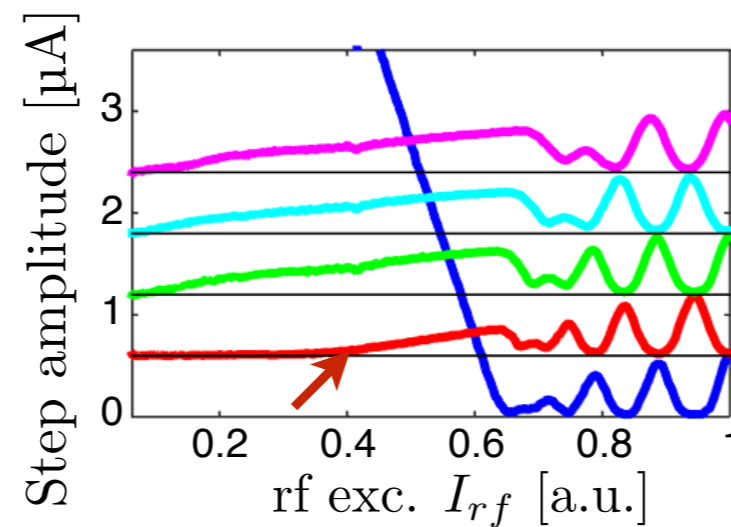
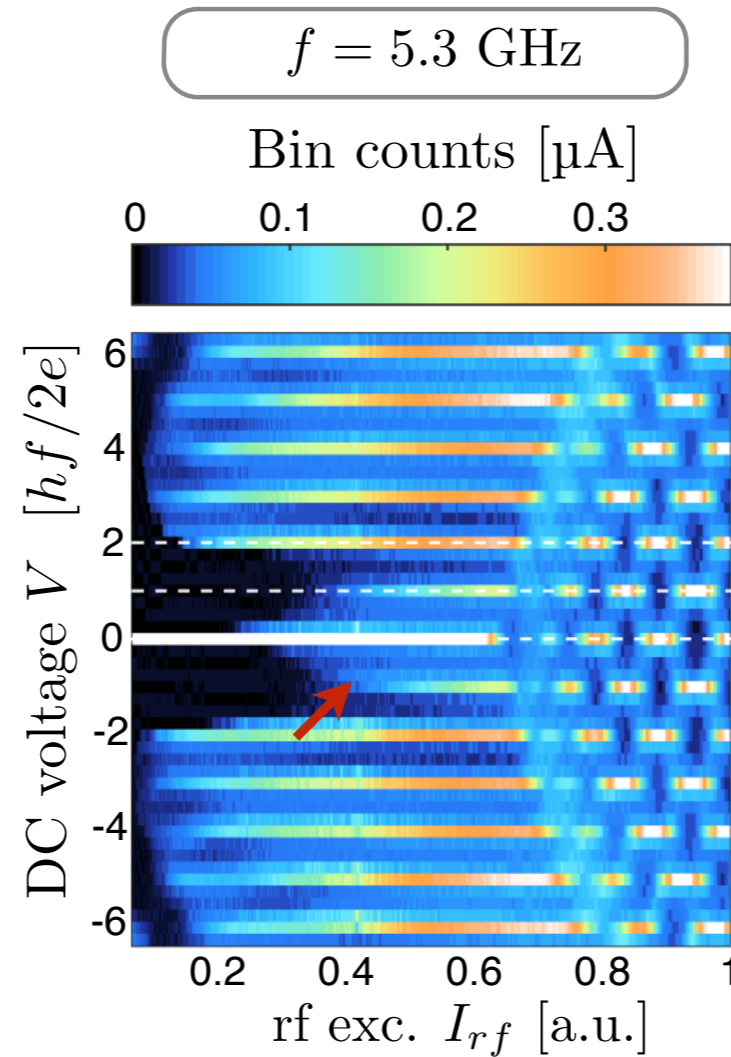
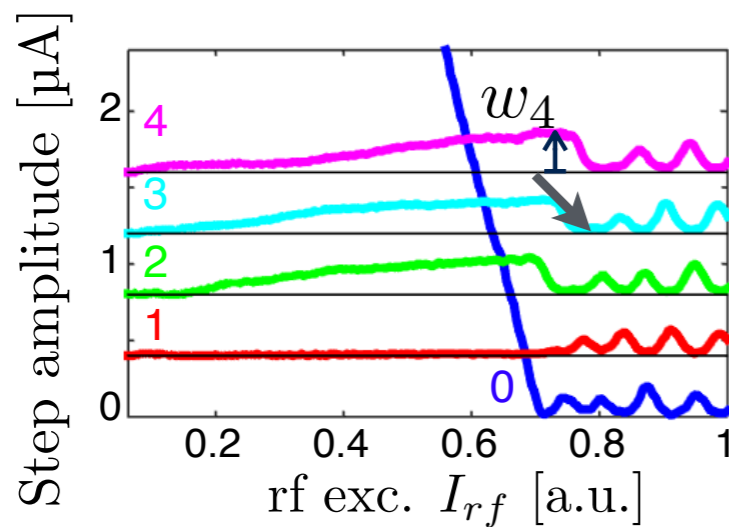
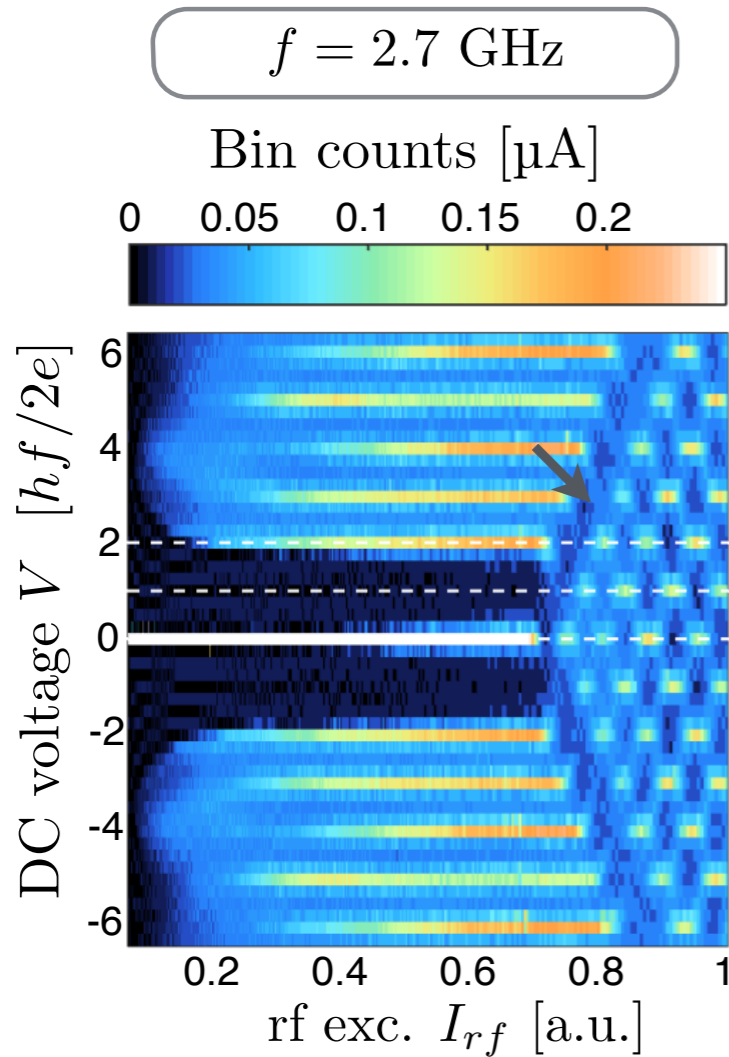
- ▷ >12 steps visible
- ▷ all steps seen: $n = 0, \pm 1, \pm 2, \dots$

Low frequency

- ▷ first step missing: $n = 0, \pm 2, \pm 3, \dots$
- ▷ 4π supercurrent?

Rokhinson *et al.*, Nat. Phys. **86**, 146503 (2012)

Frequency dependence



Mechanism for a missing step?

Additional subharmonic steps

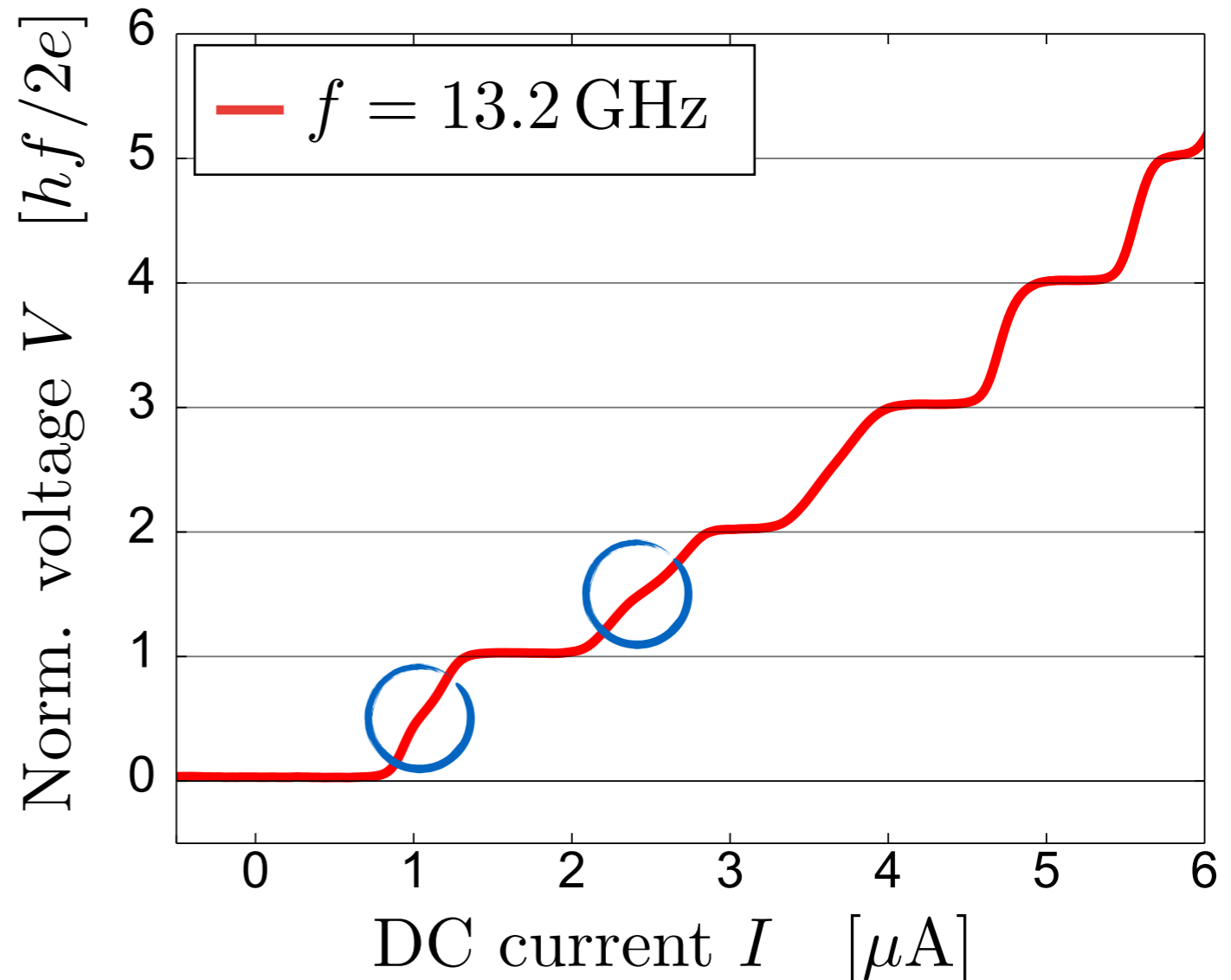
$$n = 1/2, 3/2, 1/3, \dots$$

- ▷ capacitive effects
- ▷ non-linearities
- ▷ higher harmonics in the CPR

Missing steps

- ▷ 4π supercurrent $I_{4\pi} \sin \phi/2$
- ▷ dominates at low frequency

$$f \sim f_{4\pi} = \frac{2eR_n I_{4\pi}}{h}$$
- ▷ confirmed by RSJ simulations



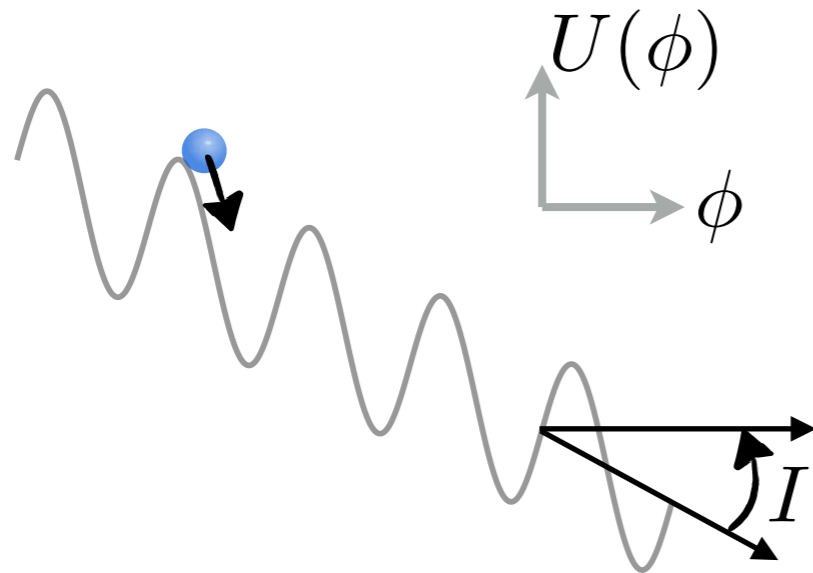
Renne *et al.*, R. Phys. App. **9**, 25 (1974)

Valizadeh *et al.*, JNMP **15**, 407 (2008)

Sochnikov *et al.*, PRL **114** 066801 (2014)

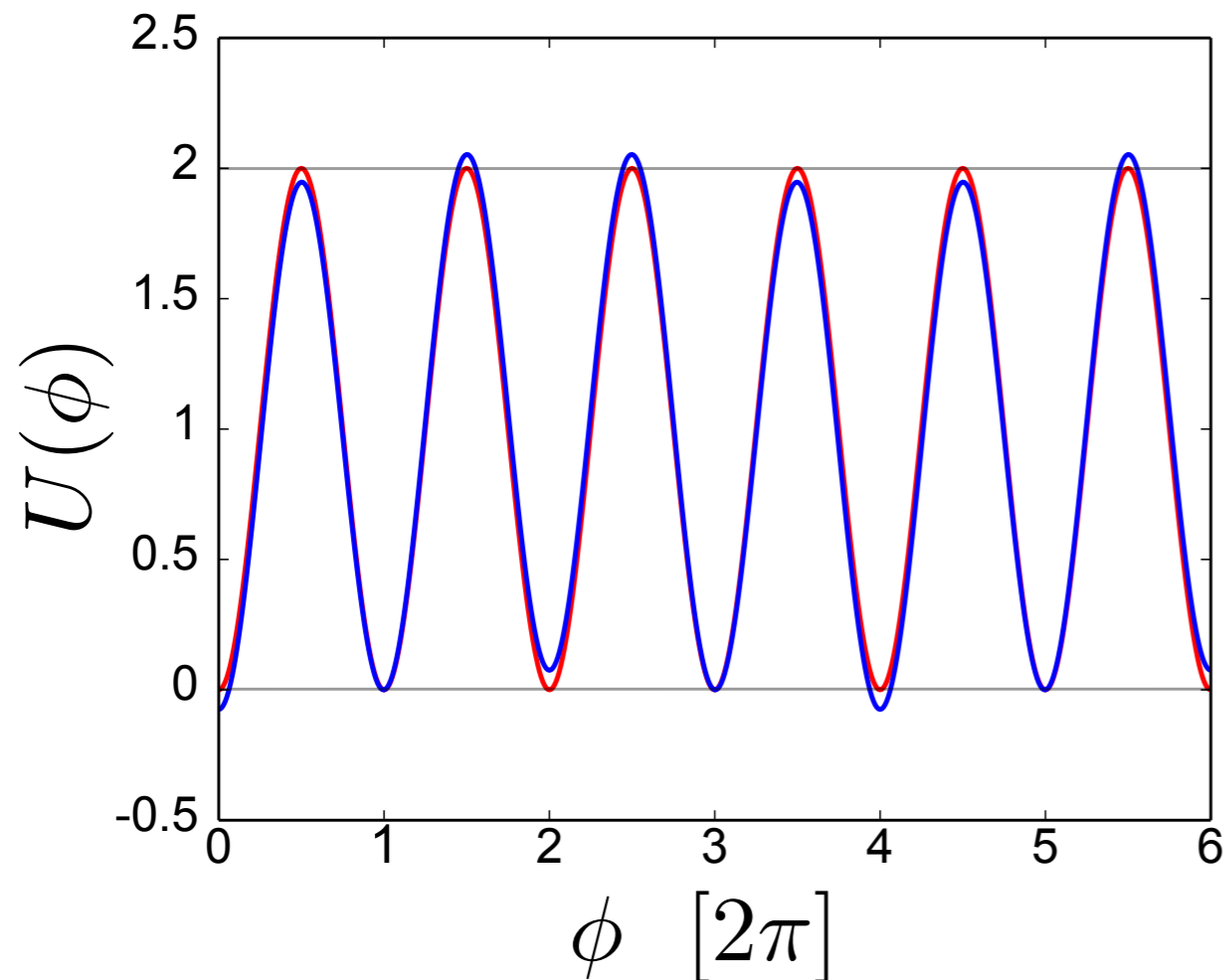
Dominguez *et al.*, PRB **86**, 146503 (2012)

Badiane *et al.*, PRL **107**, 17702 (2011)



Motion of fictitious particle

- ▷ $\partial_\phi U(\phi) = I_S(\phi) - I$
- ▷ $I_S(\phi) = I_{4\pi} \sin \frac{\phi}{2} + I_{2\pi} \sin \phi + \dots$



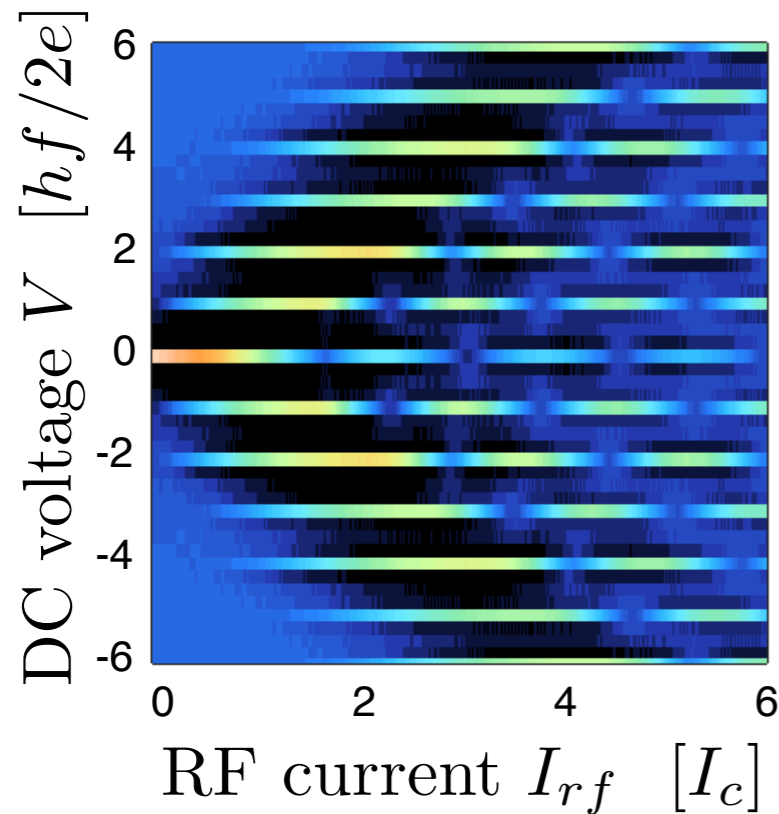
Energy potentials

- $I_{2\pi} = 1, I_{4\pi} = 0$
- $I_{2\pi} = 1, I_{4\pi} = 0.15$

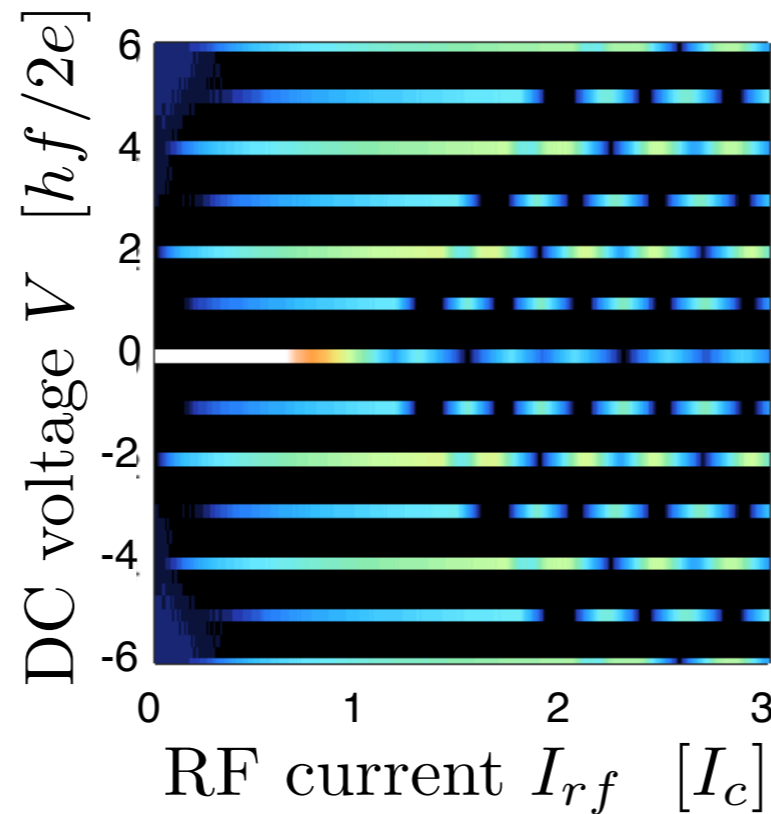
- ▷ 13% of 4π supercurrent
- ▷ 11% of increase of I_c

RSJ simulations

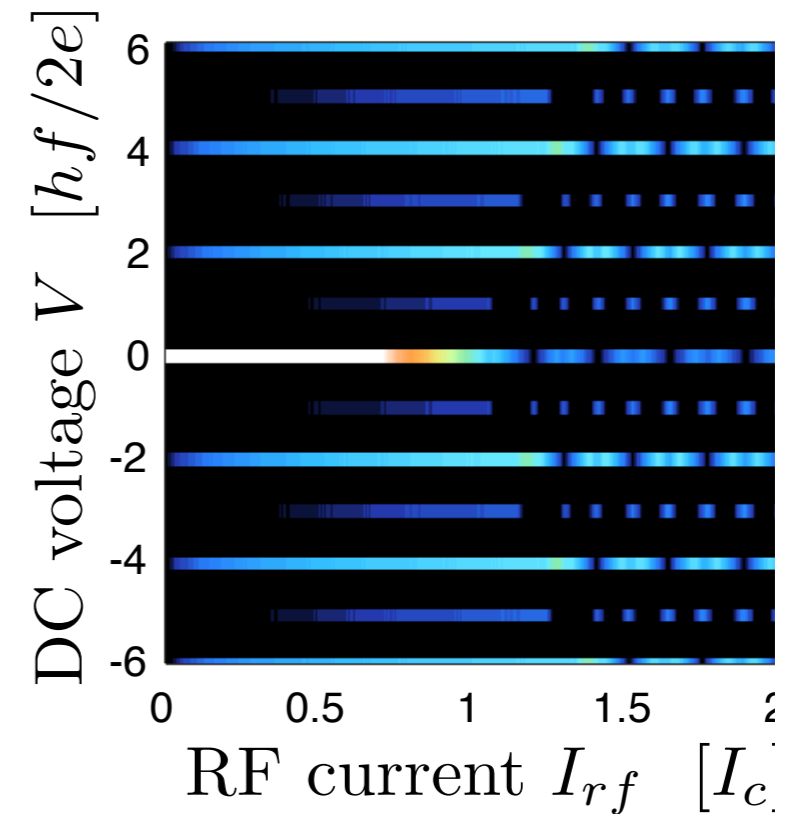
$$f = 0.5 f_J$$



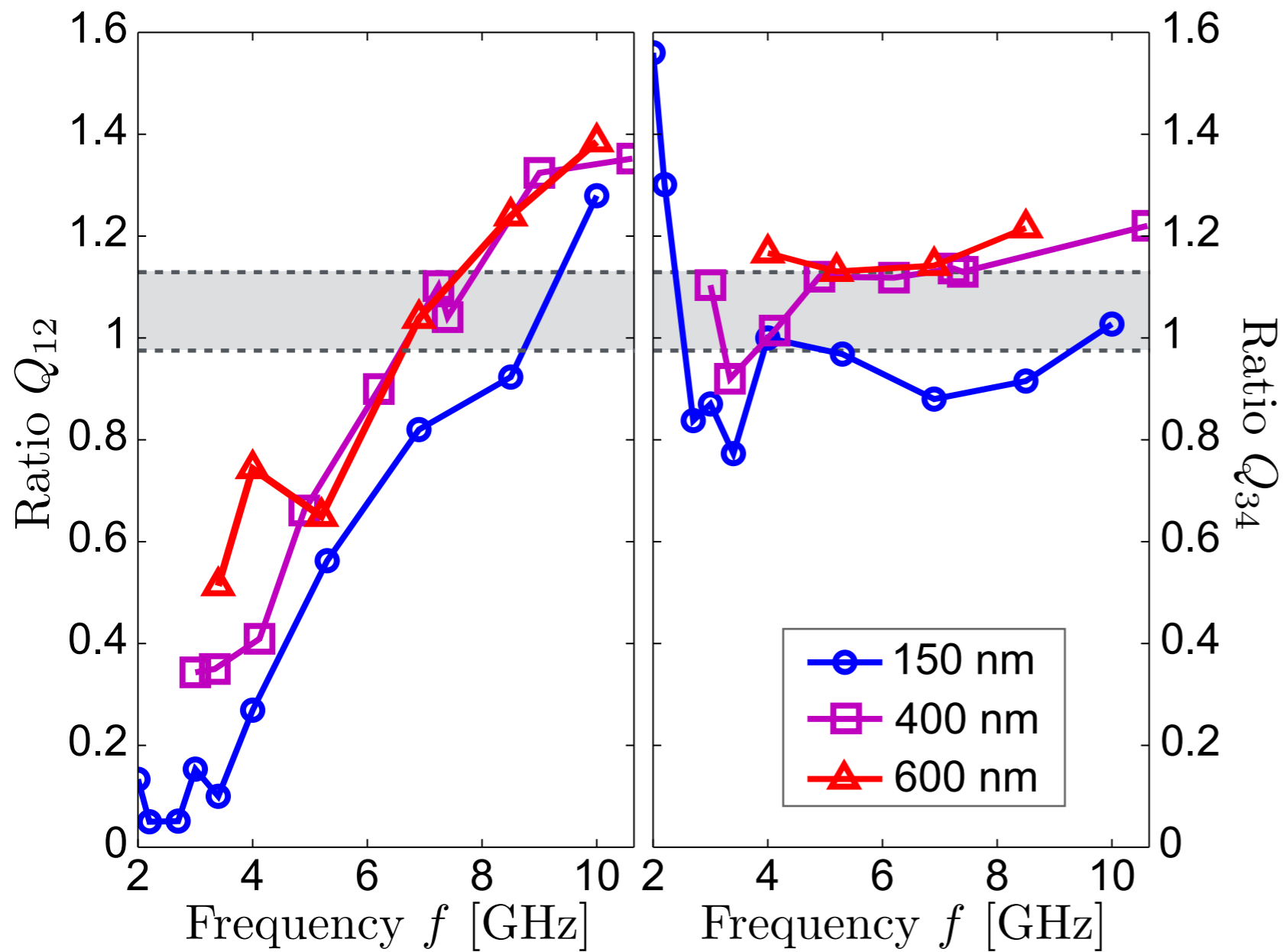
$$f = f_{4\pi} = 0.15 f_J$$



$$f = 0.05 f_J$$



- ▷ progressive disappearance of all odd steps
- ▷ crossover at $f \sim f_{4\pi} = \frac{2eR_n I_{4\pi}}{h}$
- ▷ 4π -periodicity required
- ▷ no effect of 2π -periodic CPR



Ratios

▷ $Q_{12} = \frac{w_1}{w_2} \rightarrow 0$

▷ $Q_{34} = \frac{w_3}{w_4} \simeq 1$

w_n max. amplitude of n^{th} step

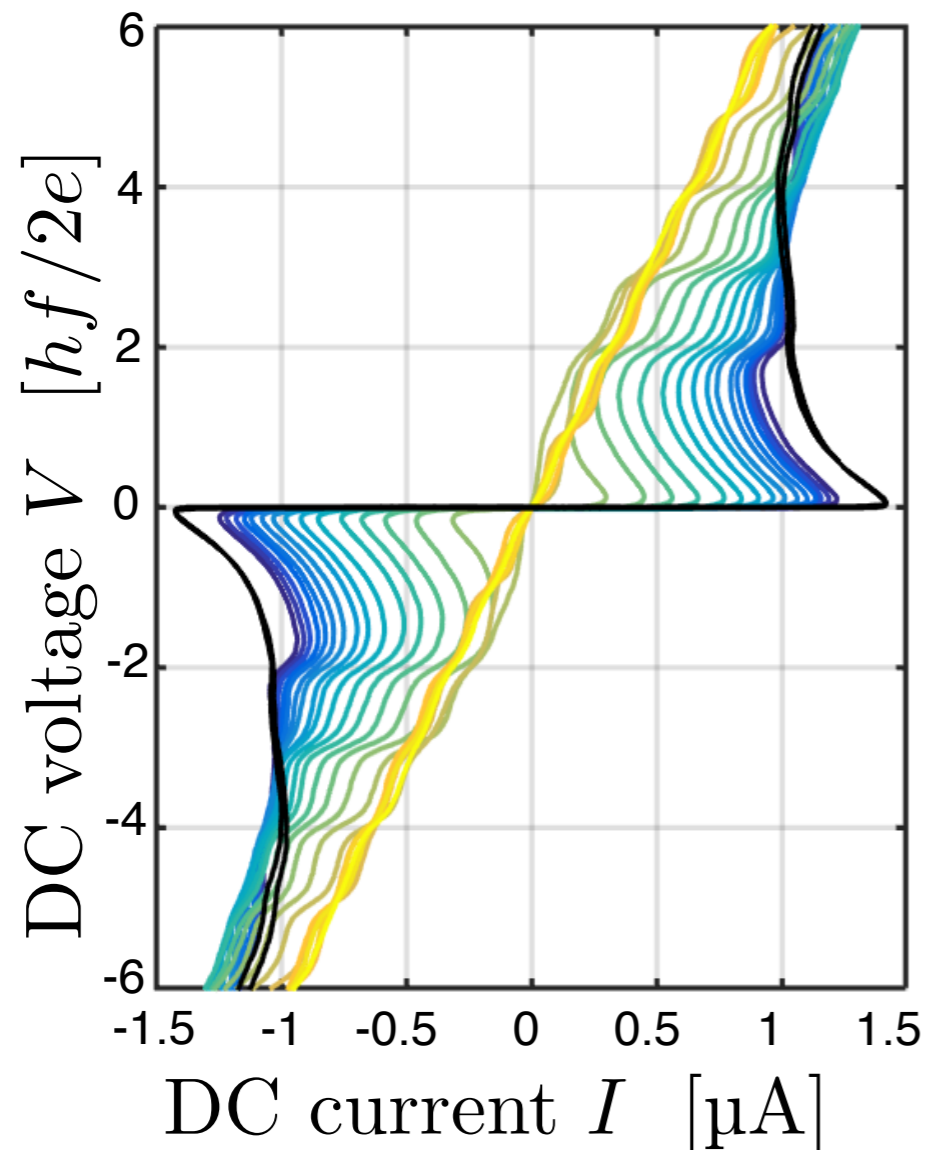
Contribution

▷ crossover $f \sim f_{4\pi}$
 \Rightarrow current $I_{4\pi}$

▷ number of 4π modes
 $I_0 = \frac{e\Delta_i}{\hbar}$
 \Rightarrow 1-3 modes

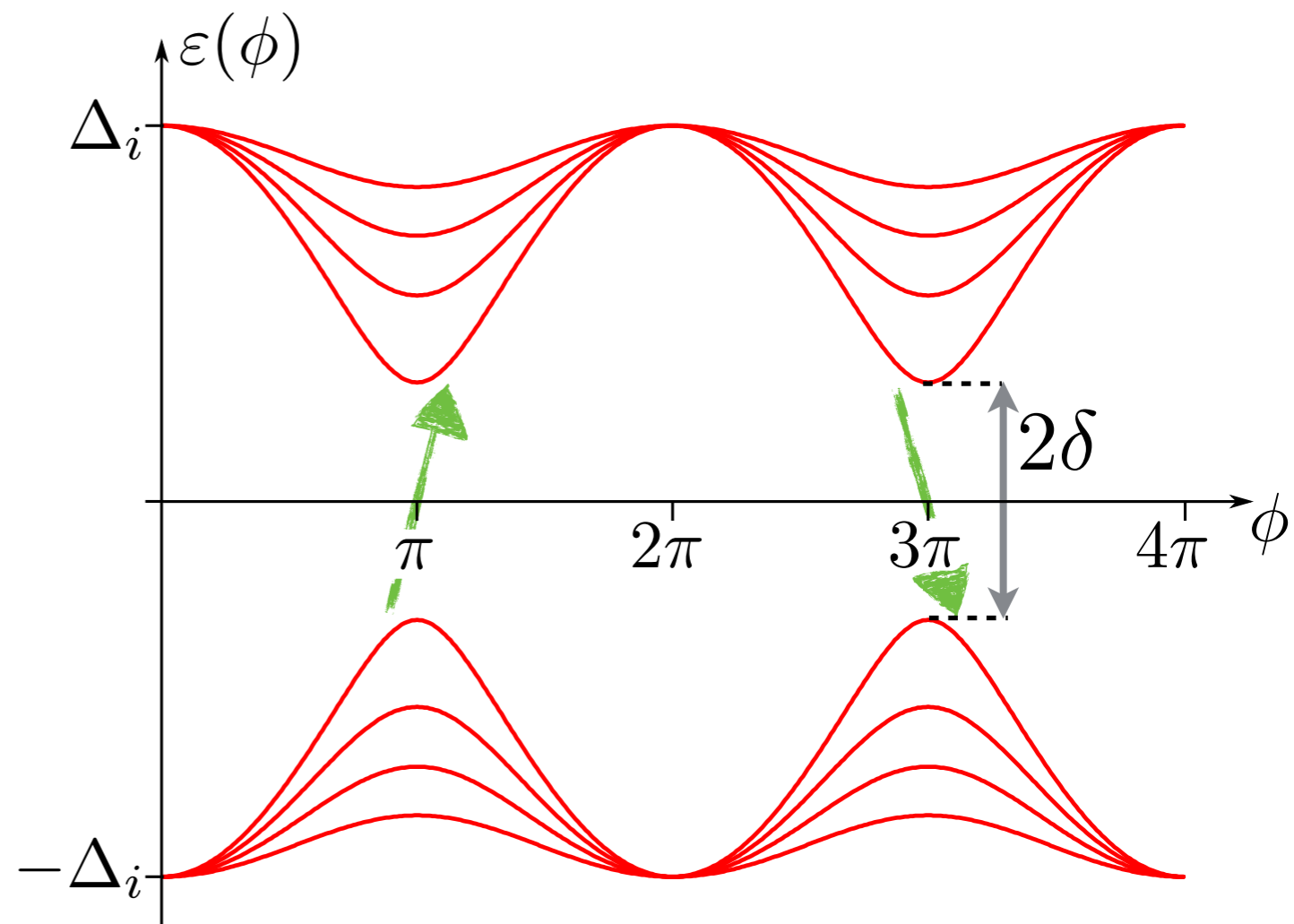
Bias instability/switch

- ▷ instability suppressed by shunt
- ▷ $n=1$ step still missing

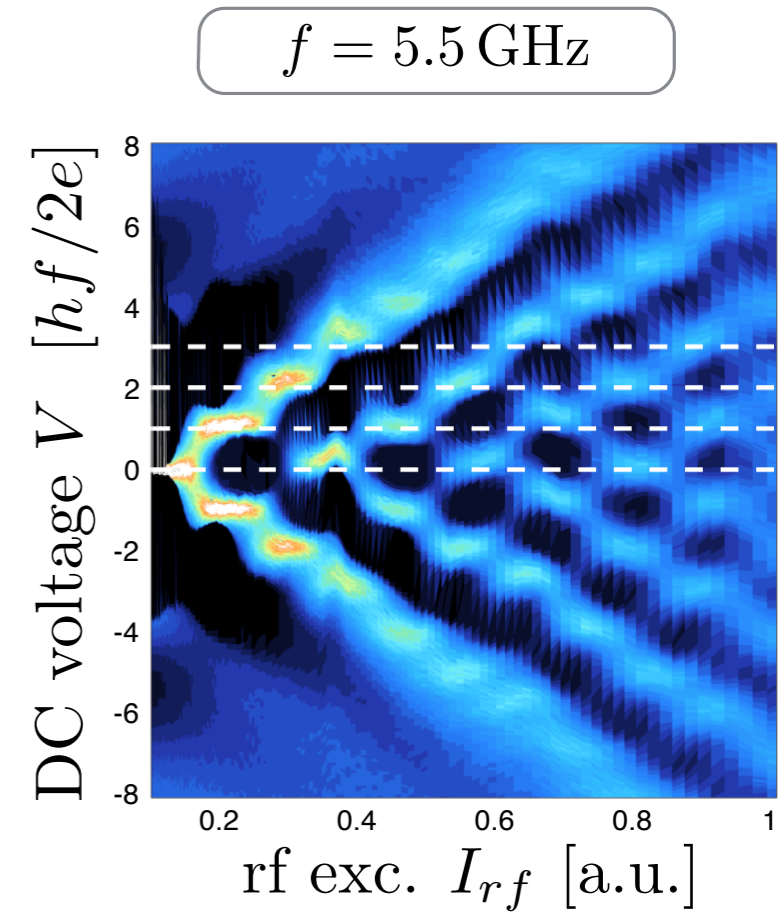
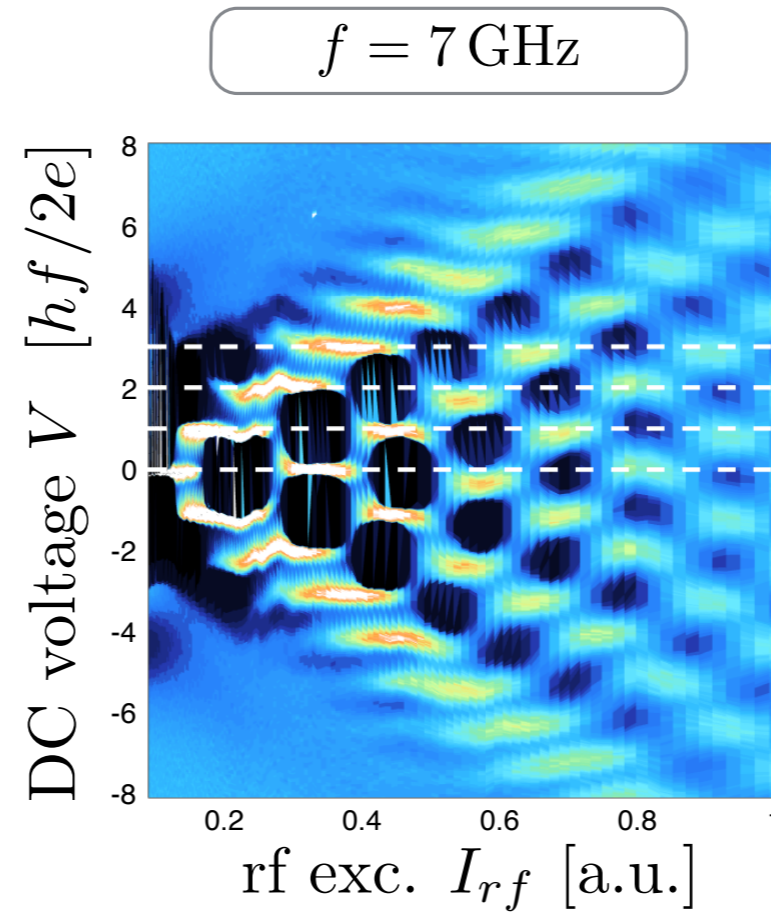
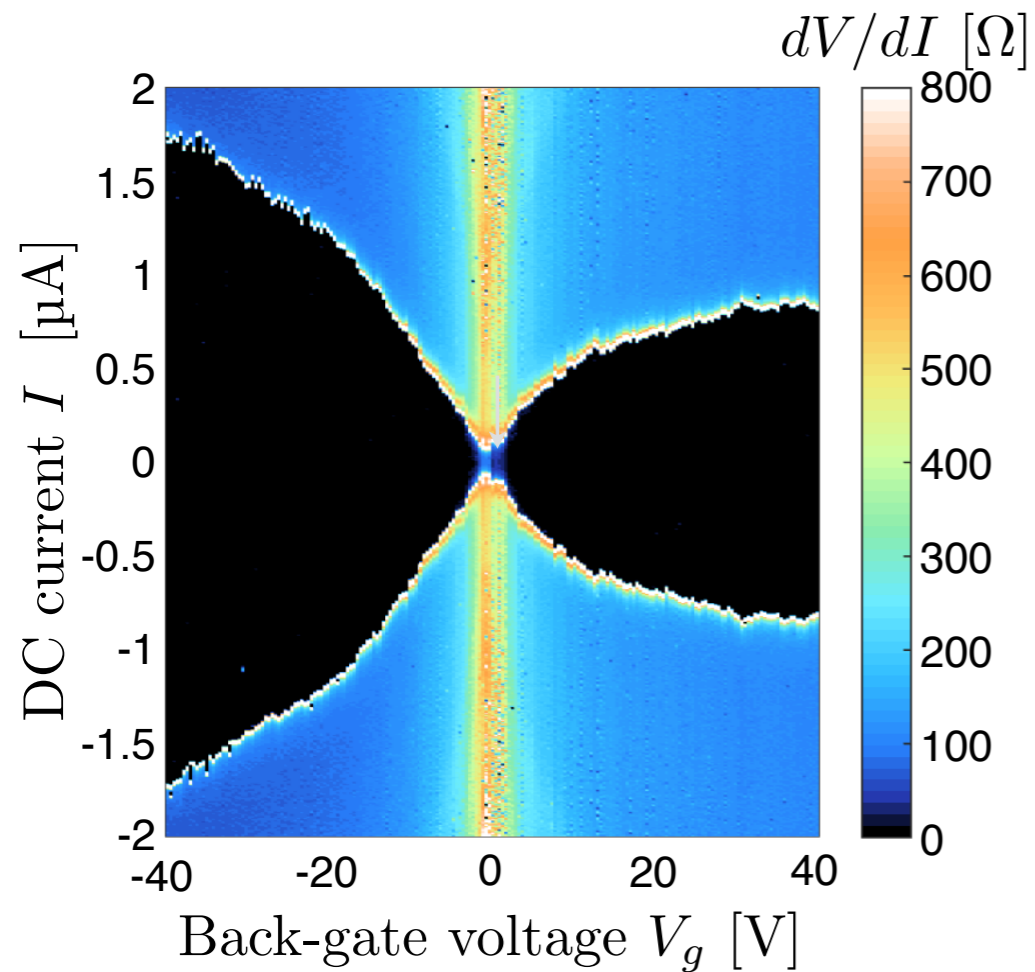


Landau-Zener transitions

- ▷ enhanced at high frequency
- ▷ evaluated minimal transmission $D > 0.994$



Graphene-based junctions



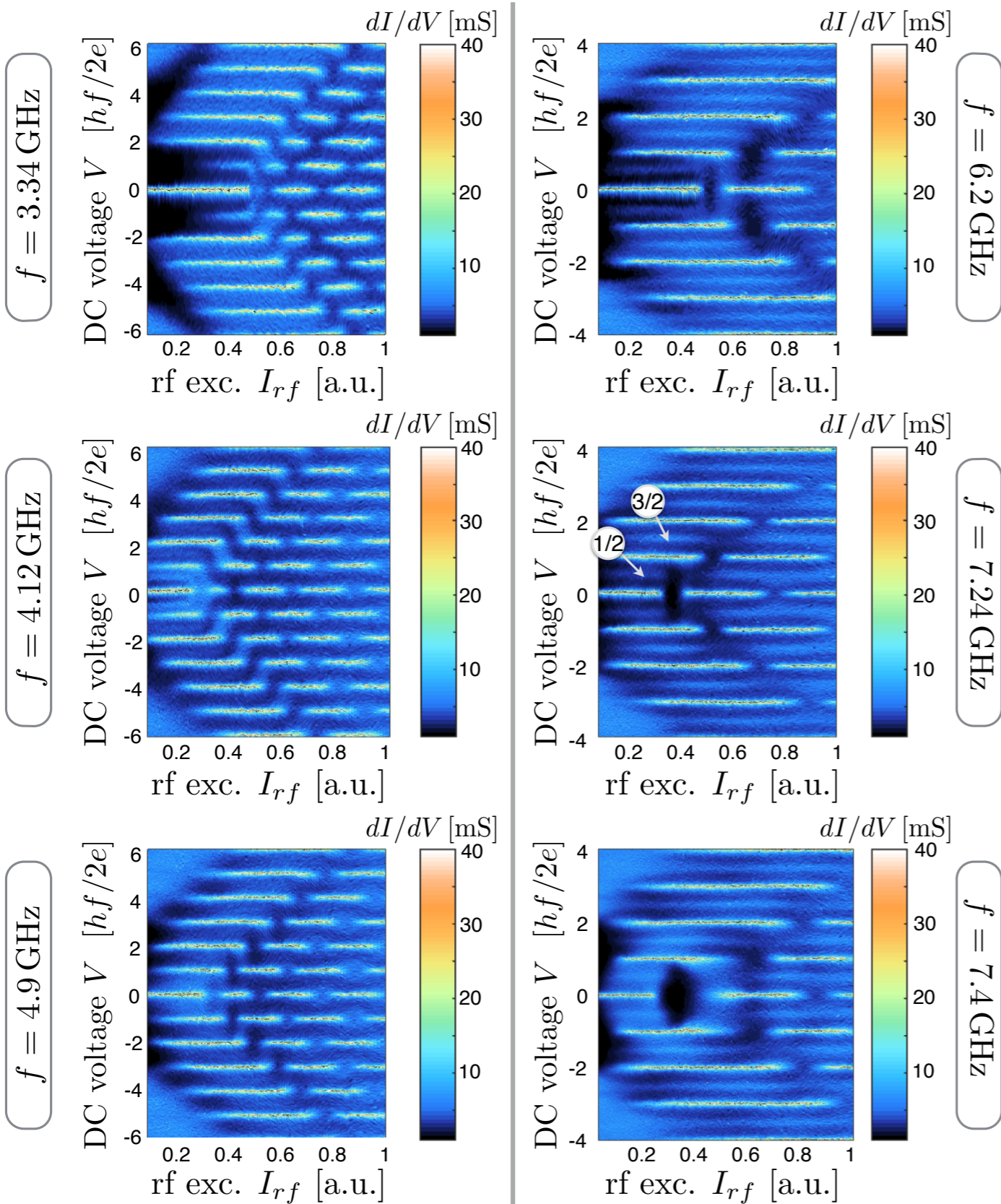
Graphene junctions

- ▷ density : $-2.5 \times 10^{12} \text{cm}^{-2}$ to $2.5 \times 10^{12} \text{cm}^{-2}$
- ▷ mobility : $\sim 5000 \text{cm}^2 \text{V}^{-1} \text{s}^{-1}$

Shapiro steps

- ▷ $n=1$ step always visible
- ▷ different densities (n , p , DP) and frequencies (4-14 GHz)

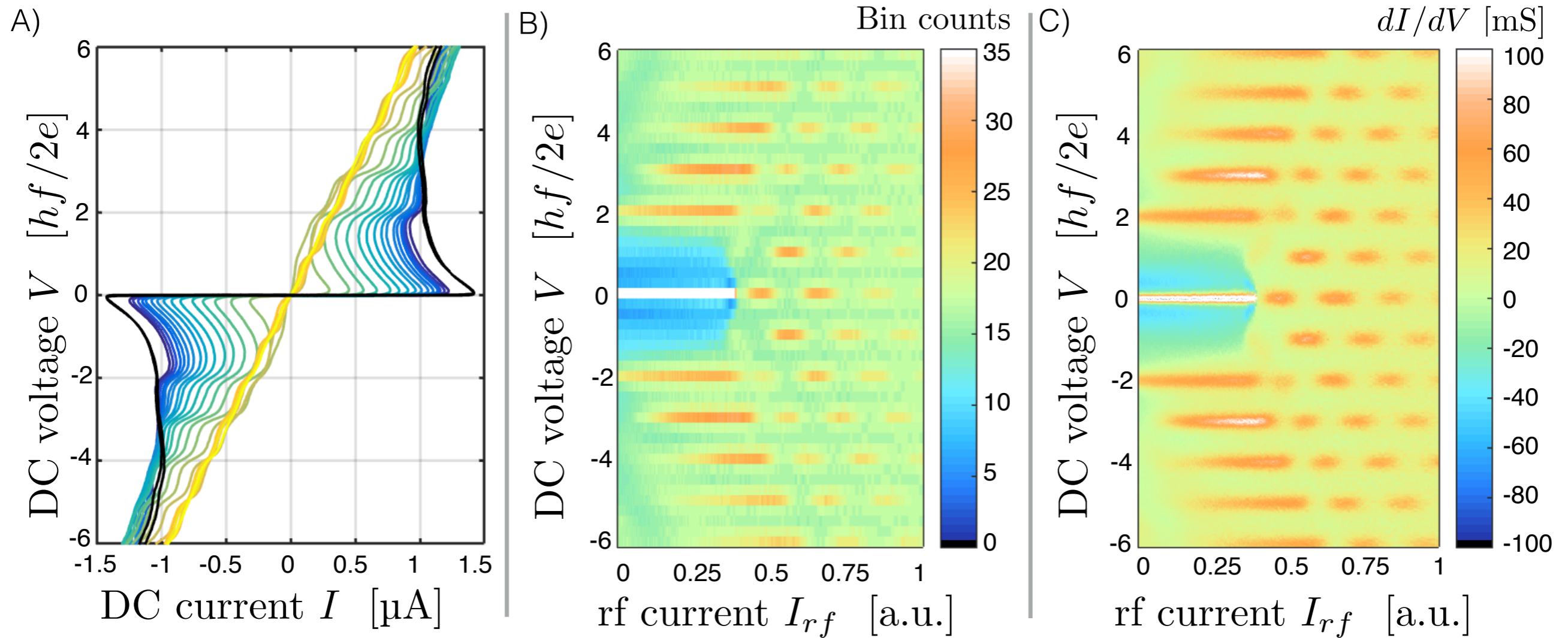
Frequency dependence

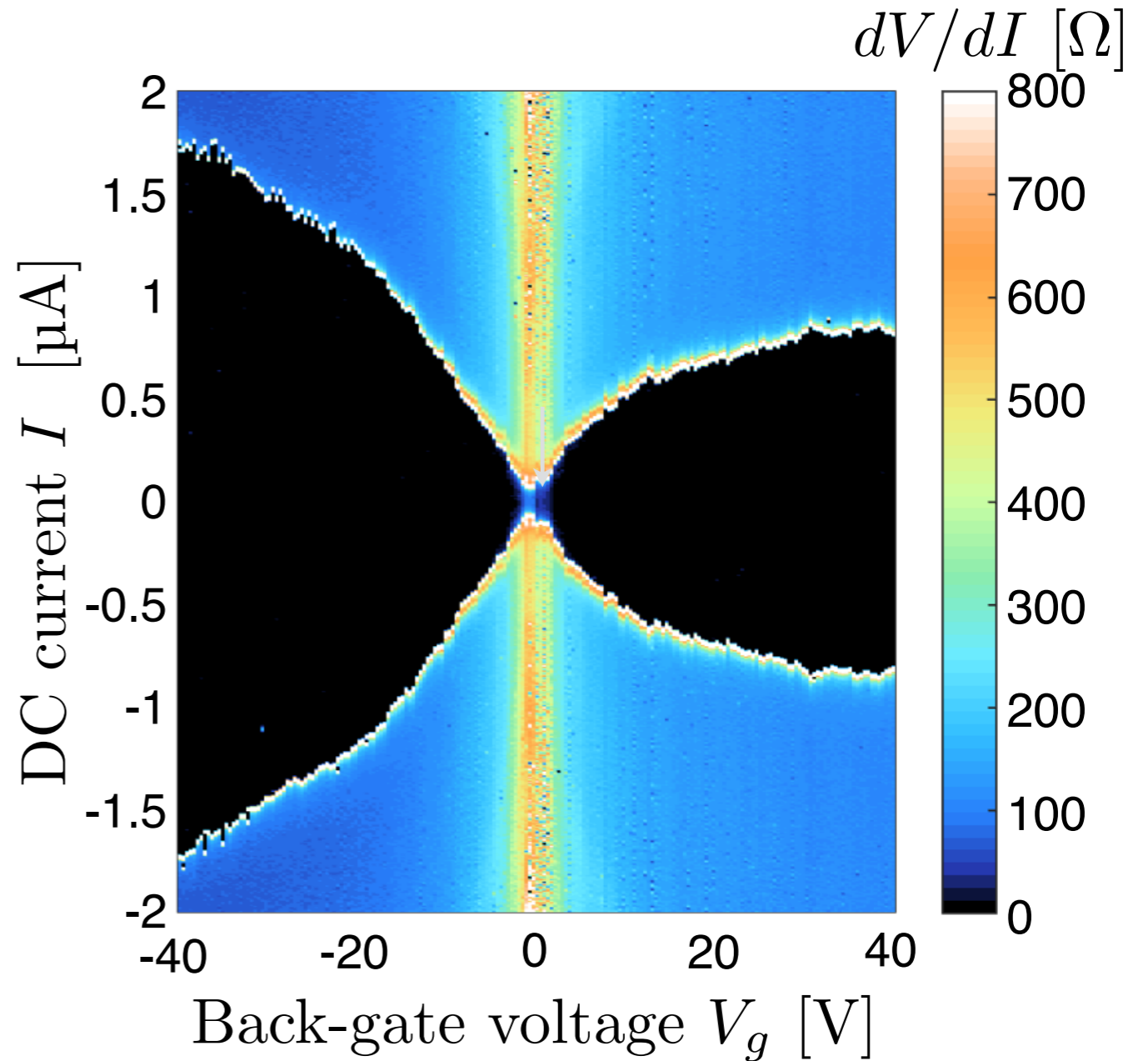
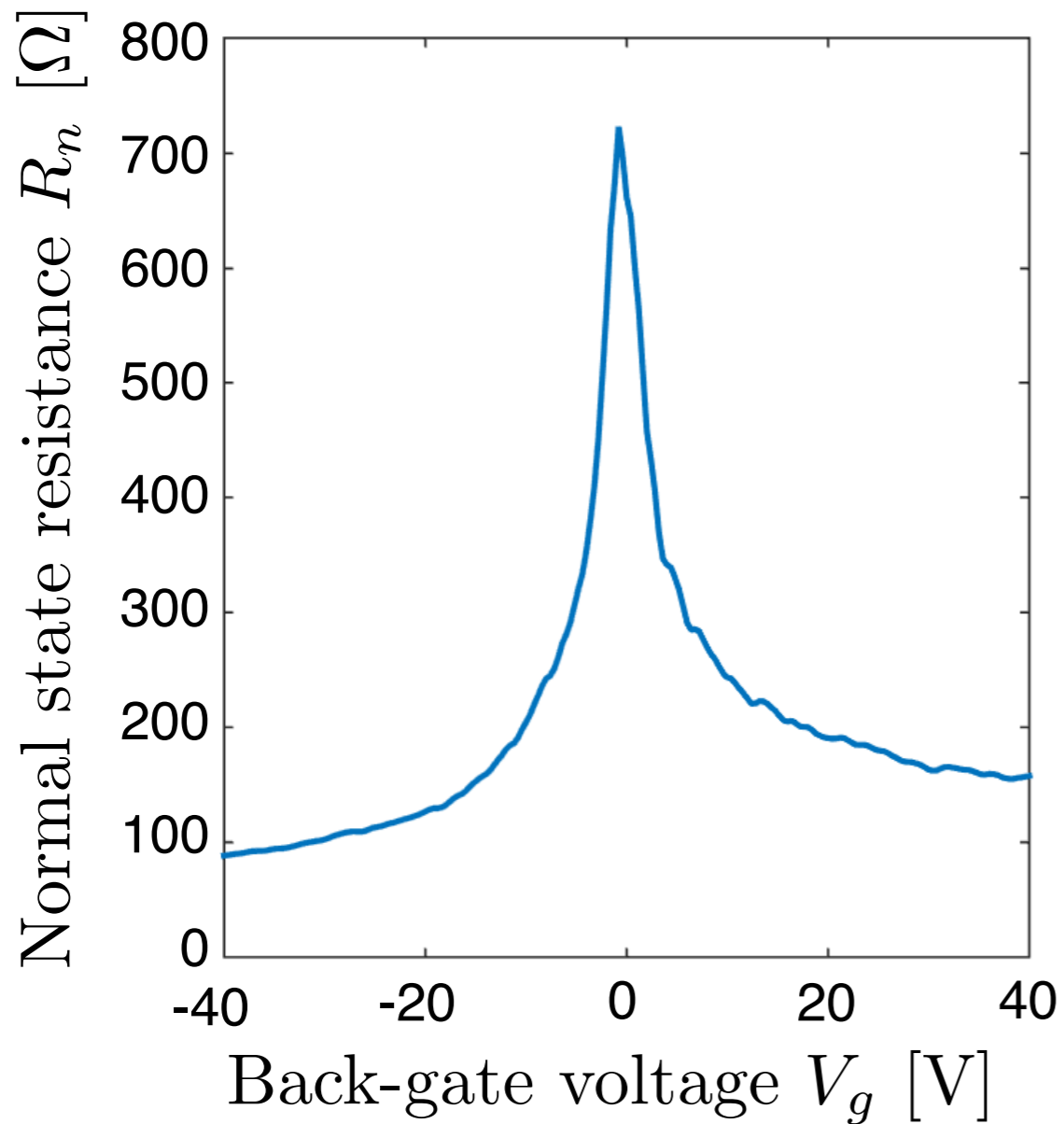


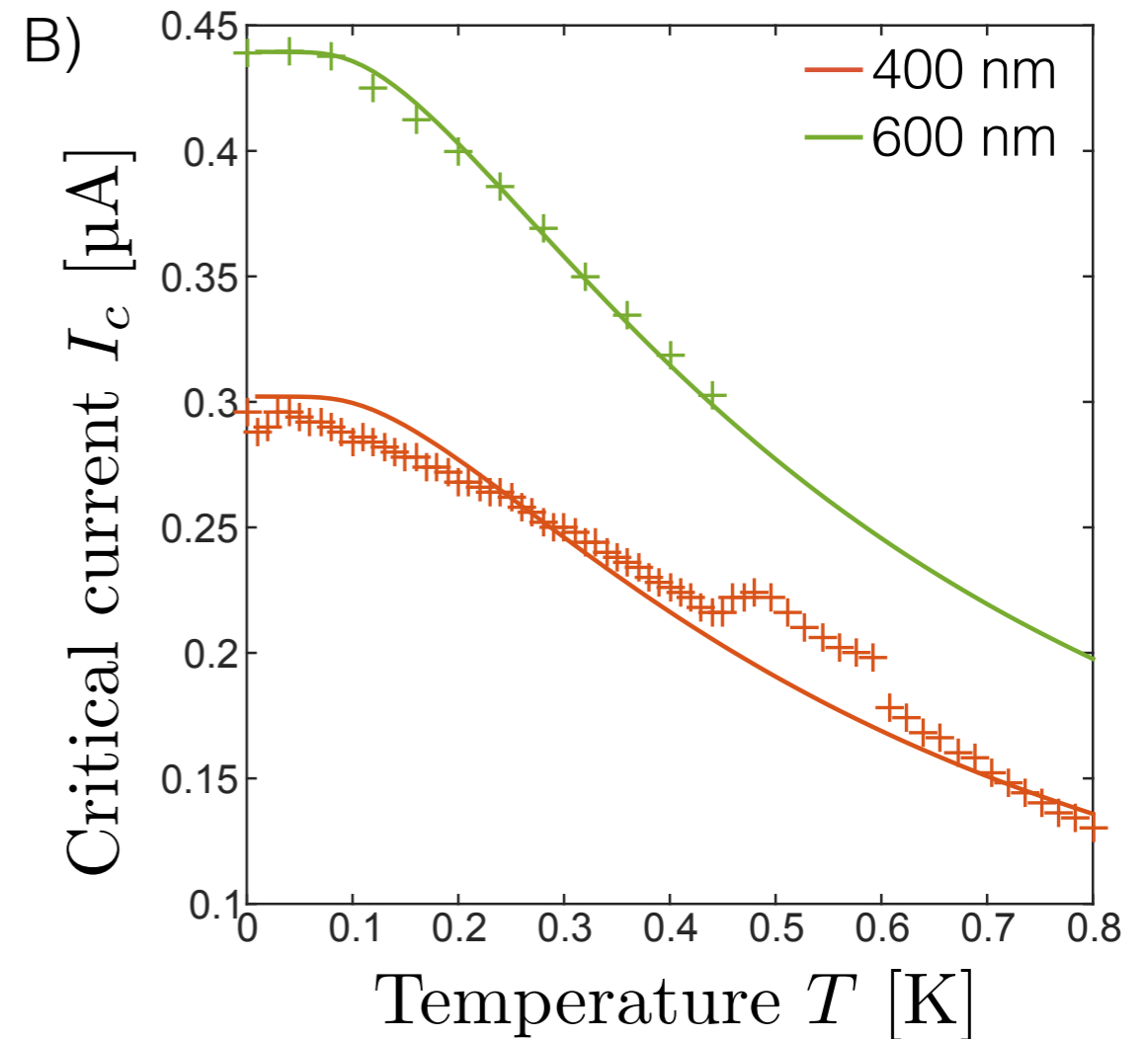
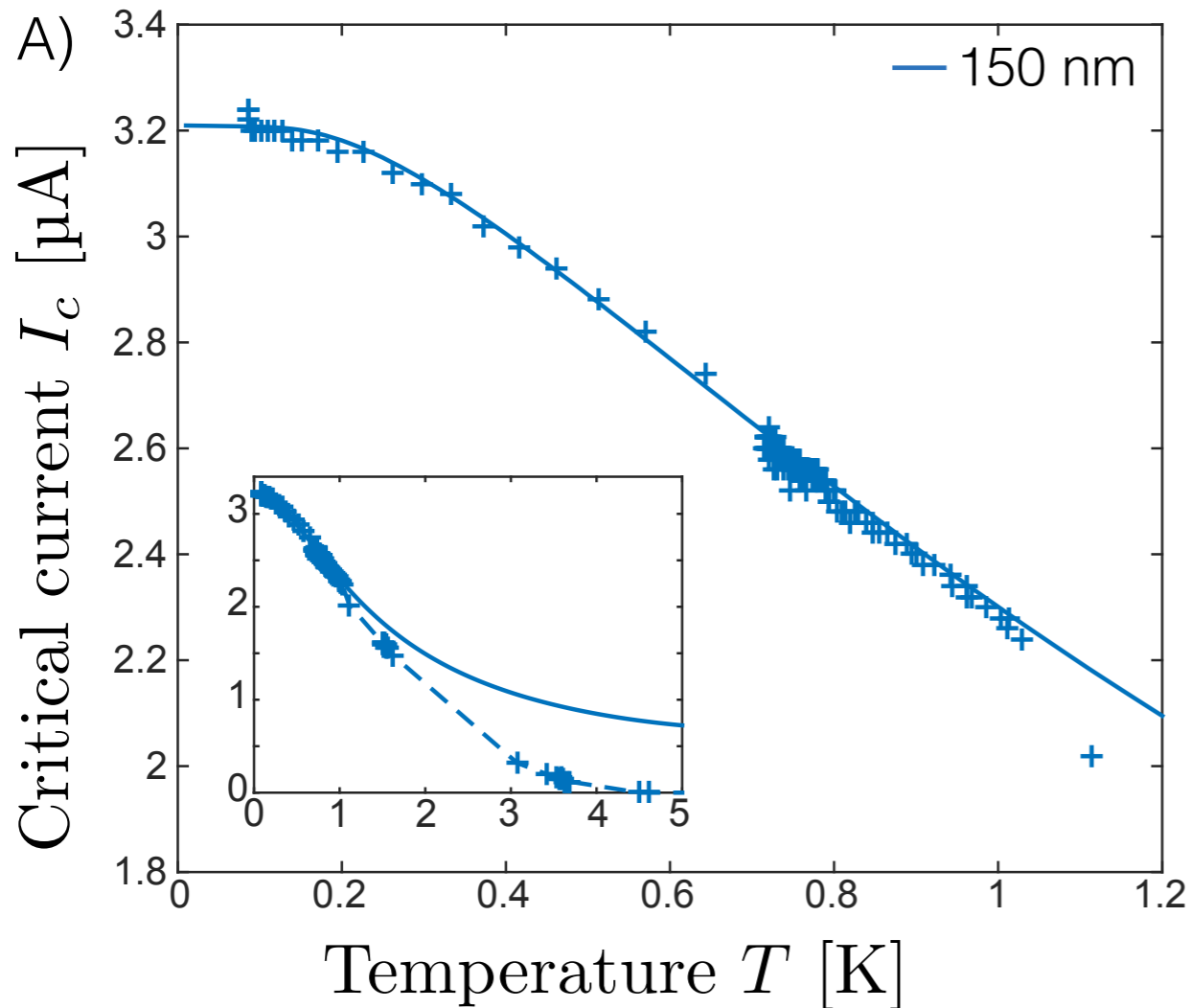
Frequency dependence

- ▶ missing $n=1$ step at low frequency
- ▶ subharmonic steps at high frequency

Shapiro steps : voltage bias



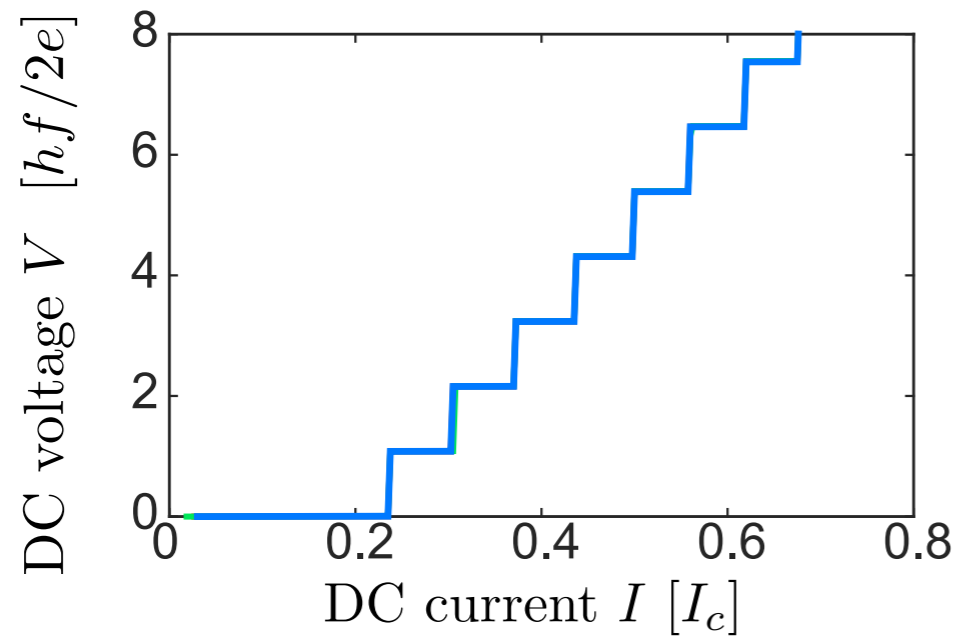
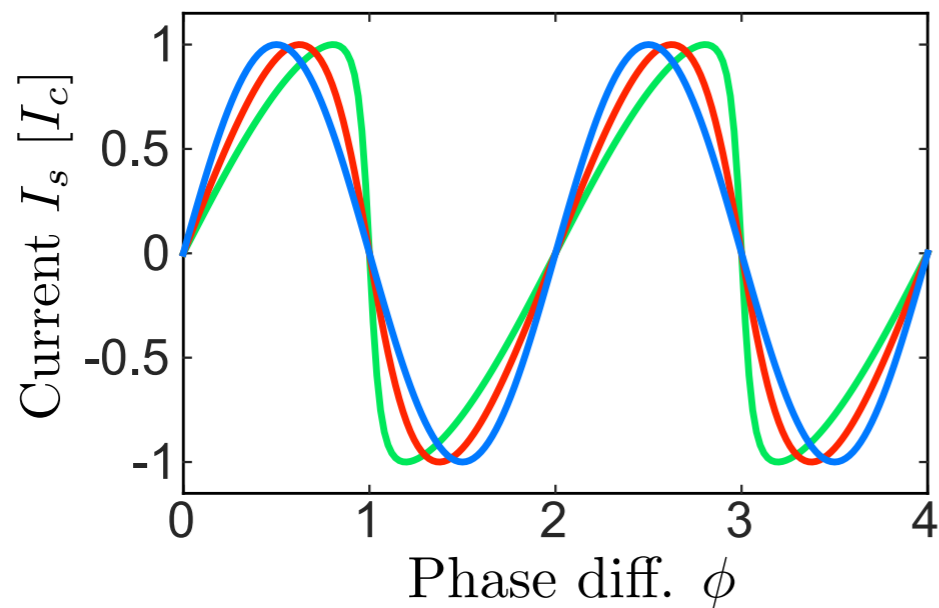
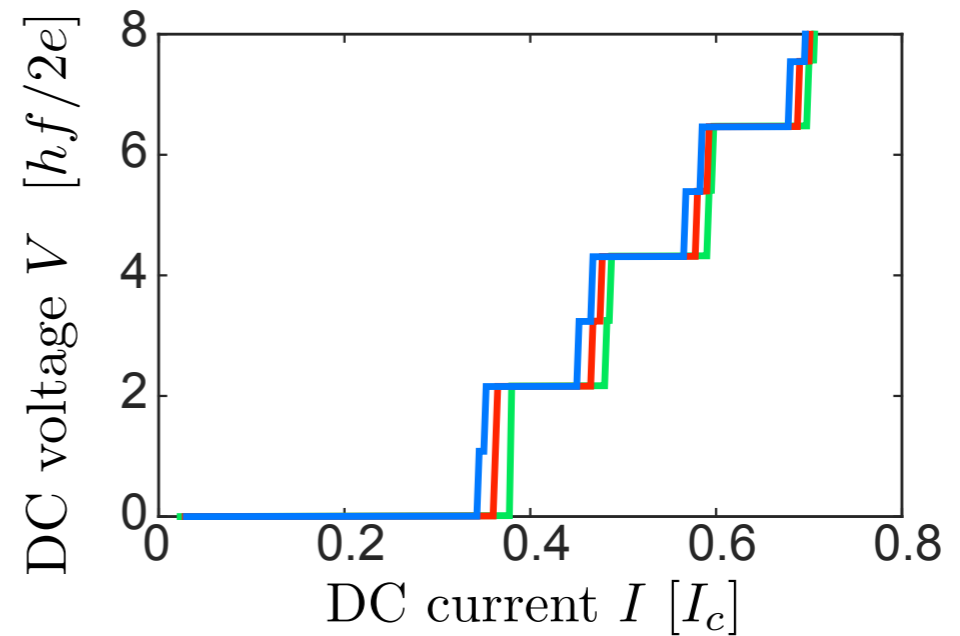
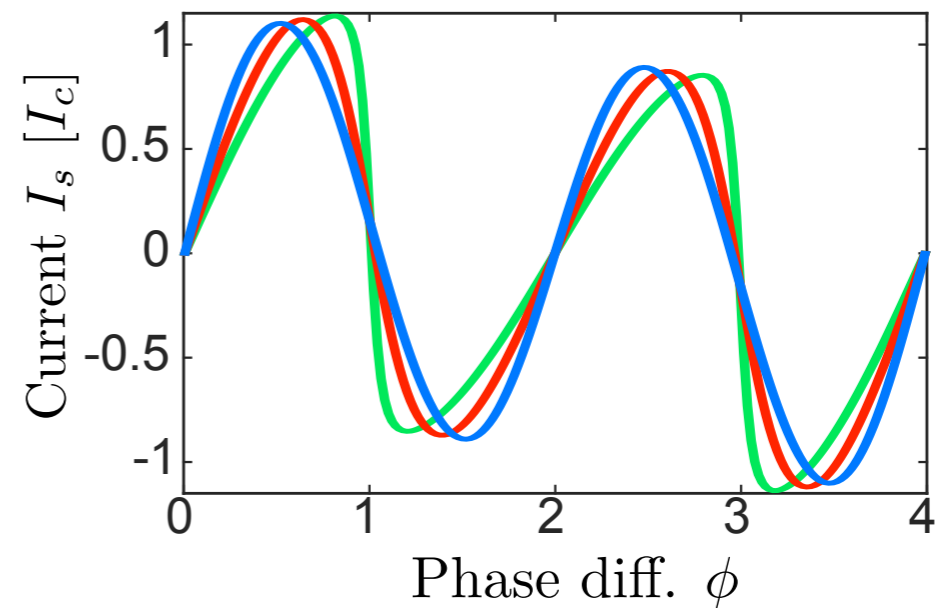


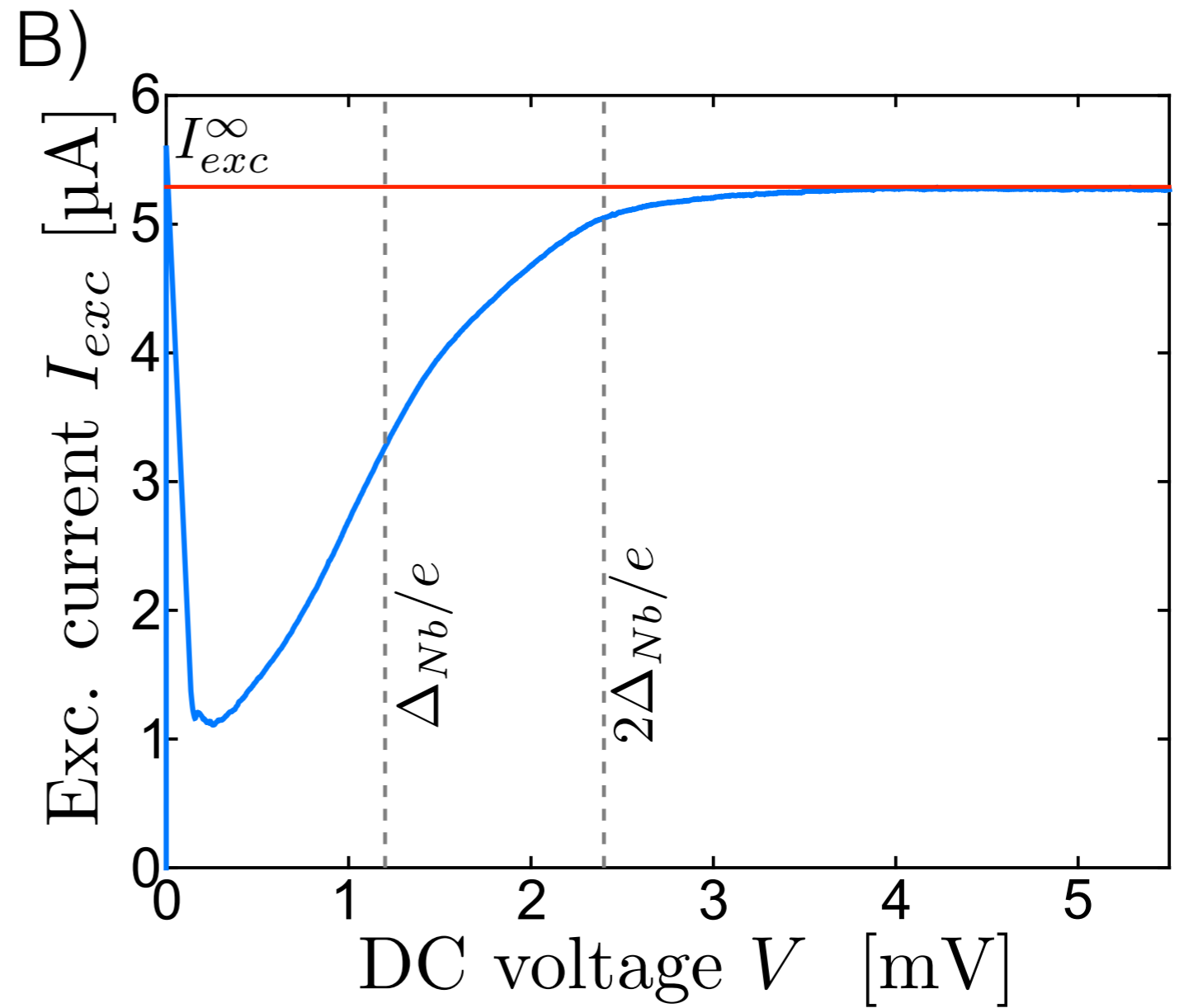
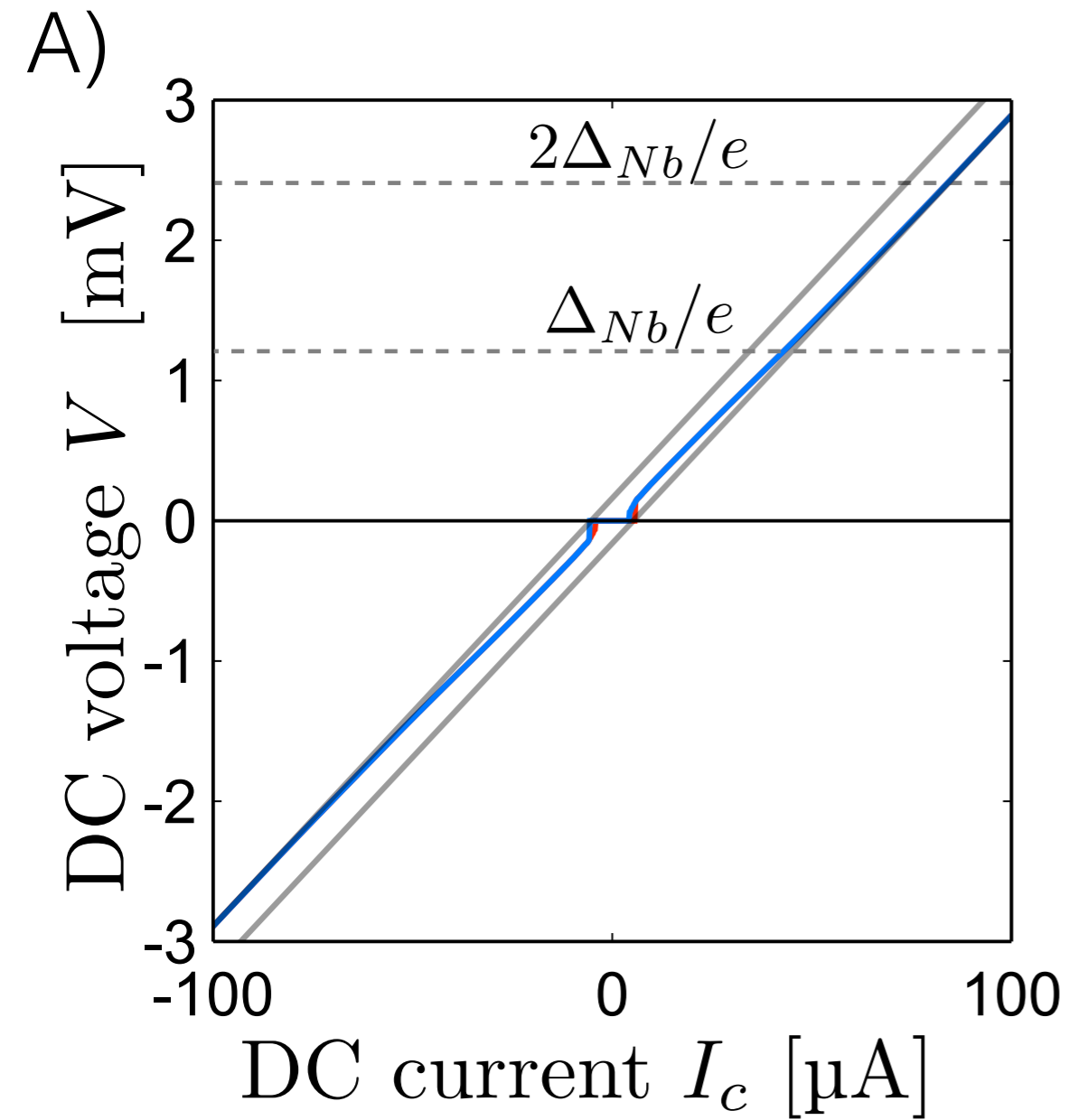
Temperature dependence of I_c 

- ▷ fits for in 0-1K region to access induced gap
- ▷ perturbation theory diverges at high temperature

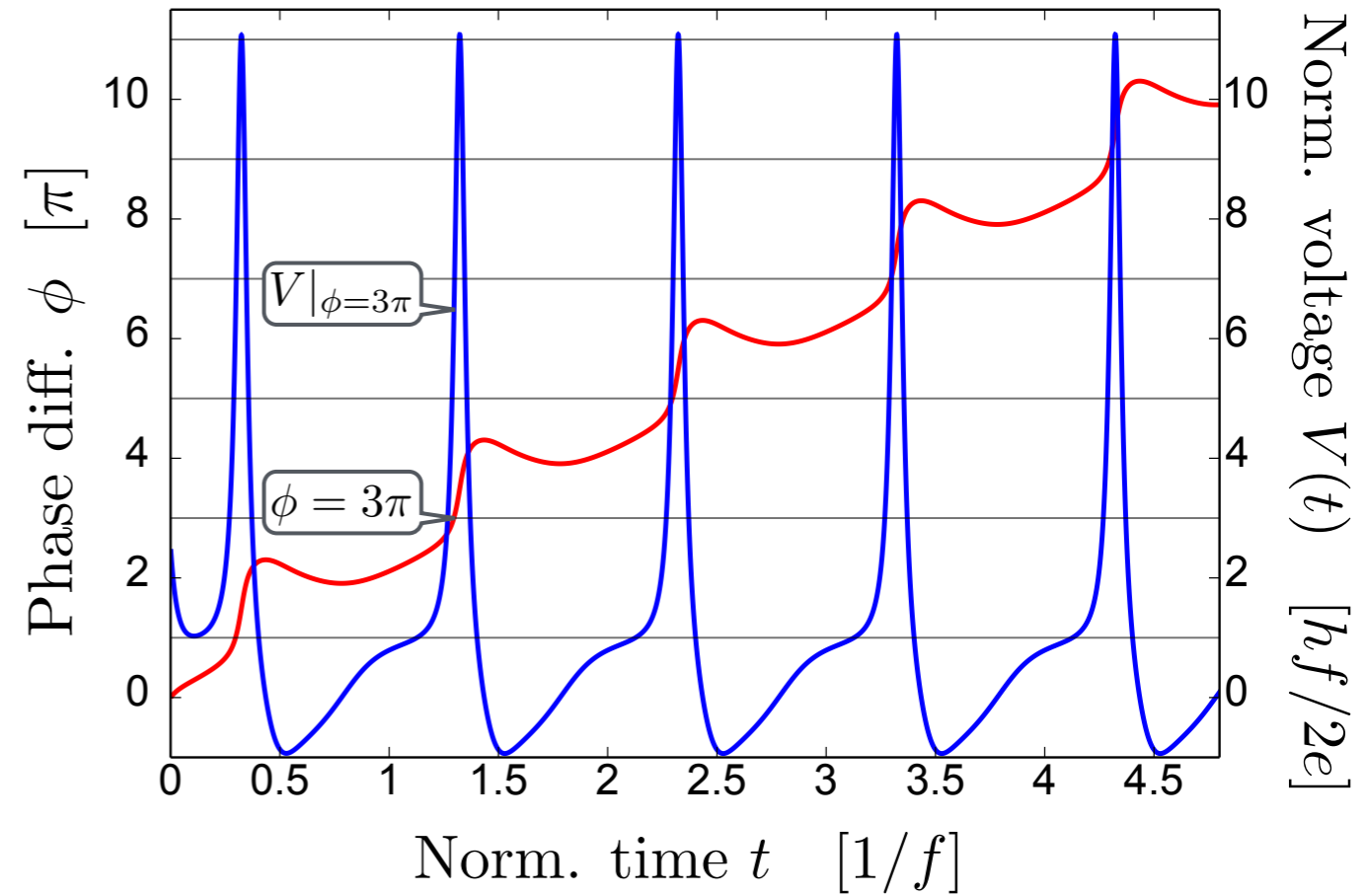
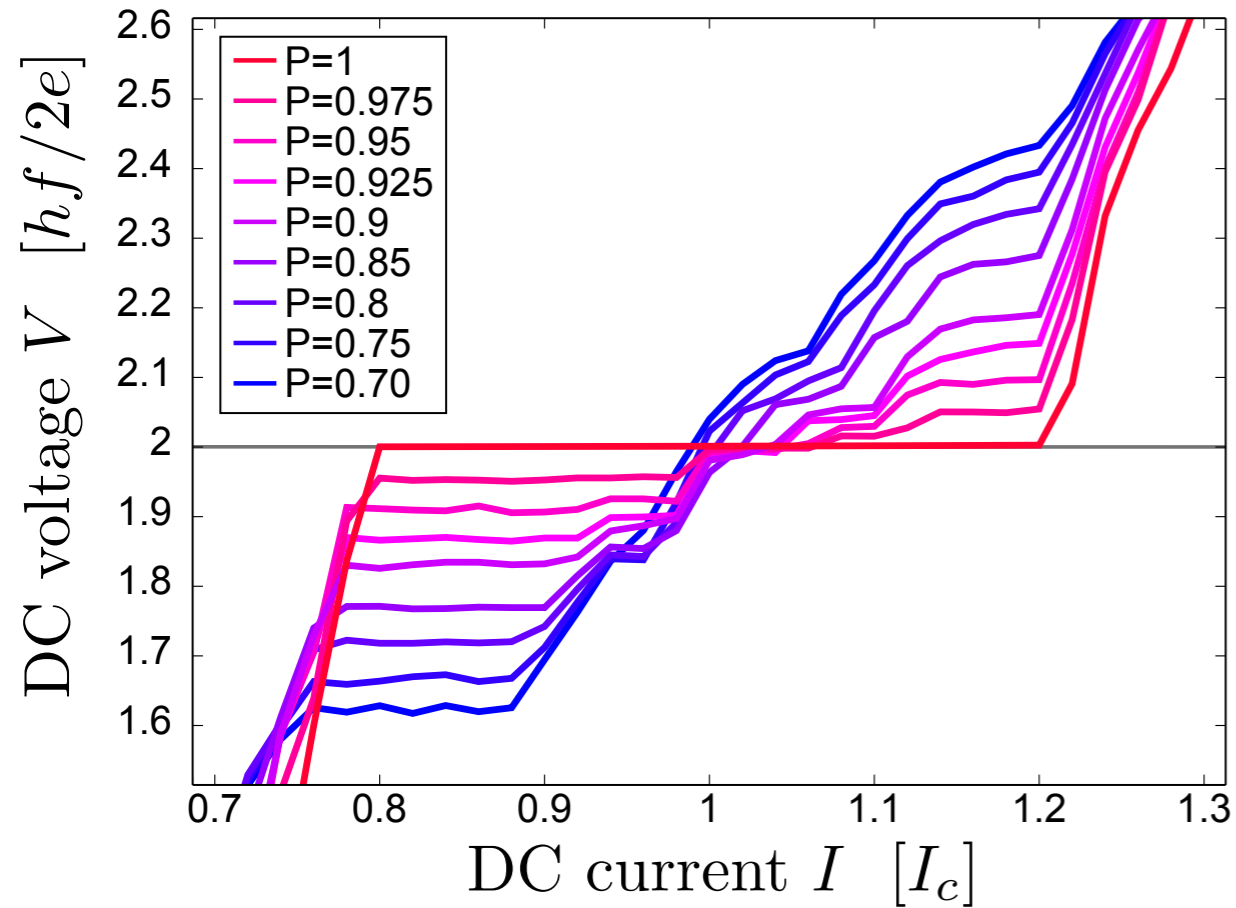
Tkachov *et al.*, PRB **88**, 075401 (2013)

Effect of CPR in RSJ simulations

 2π modes $2\pi + 4\pi$ modes

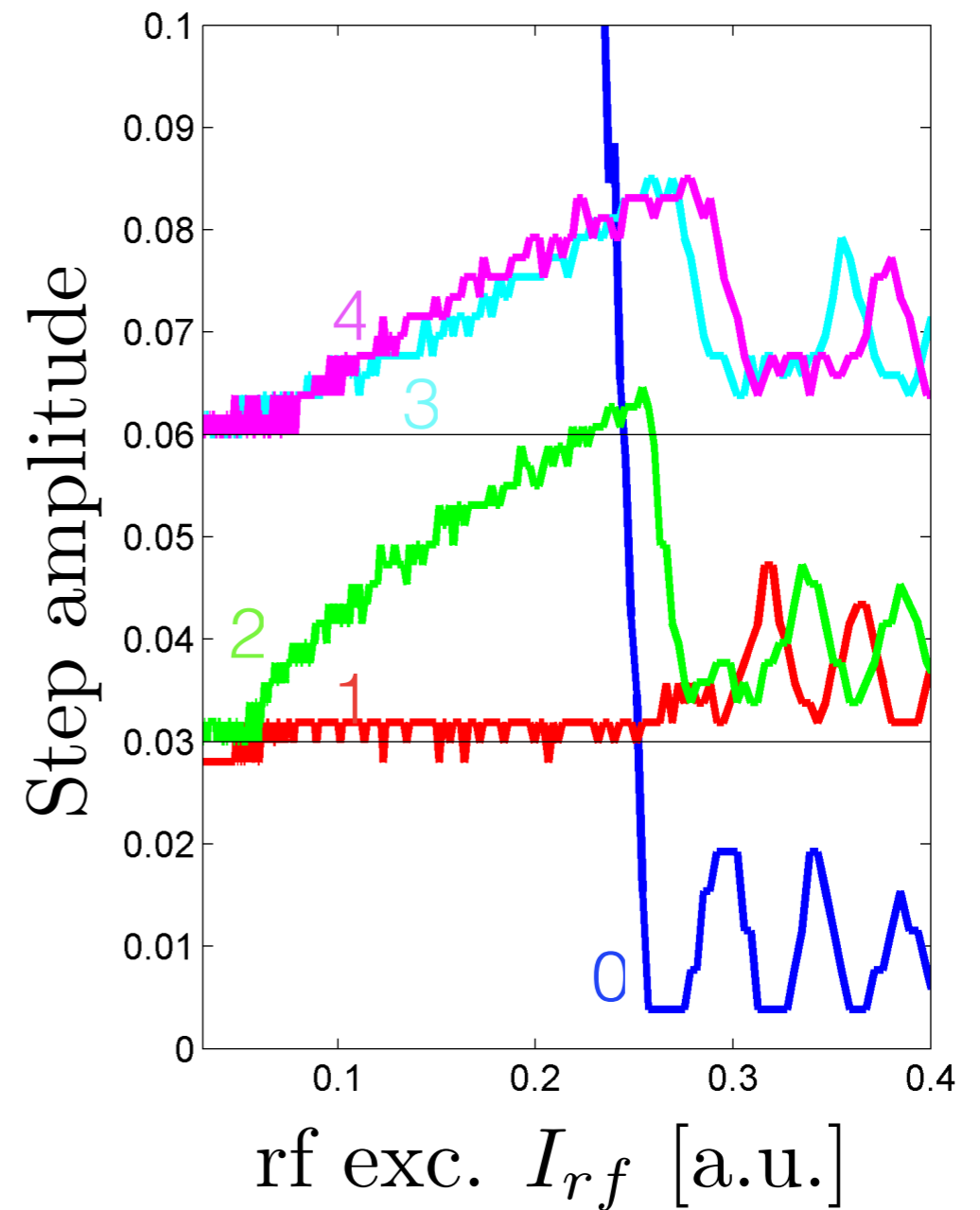
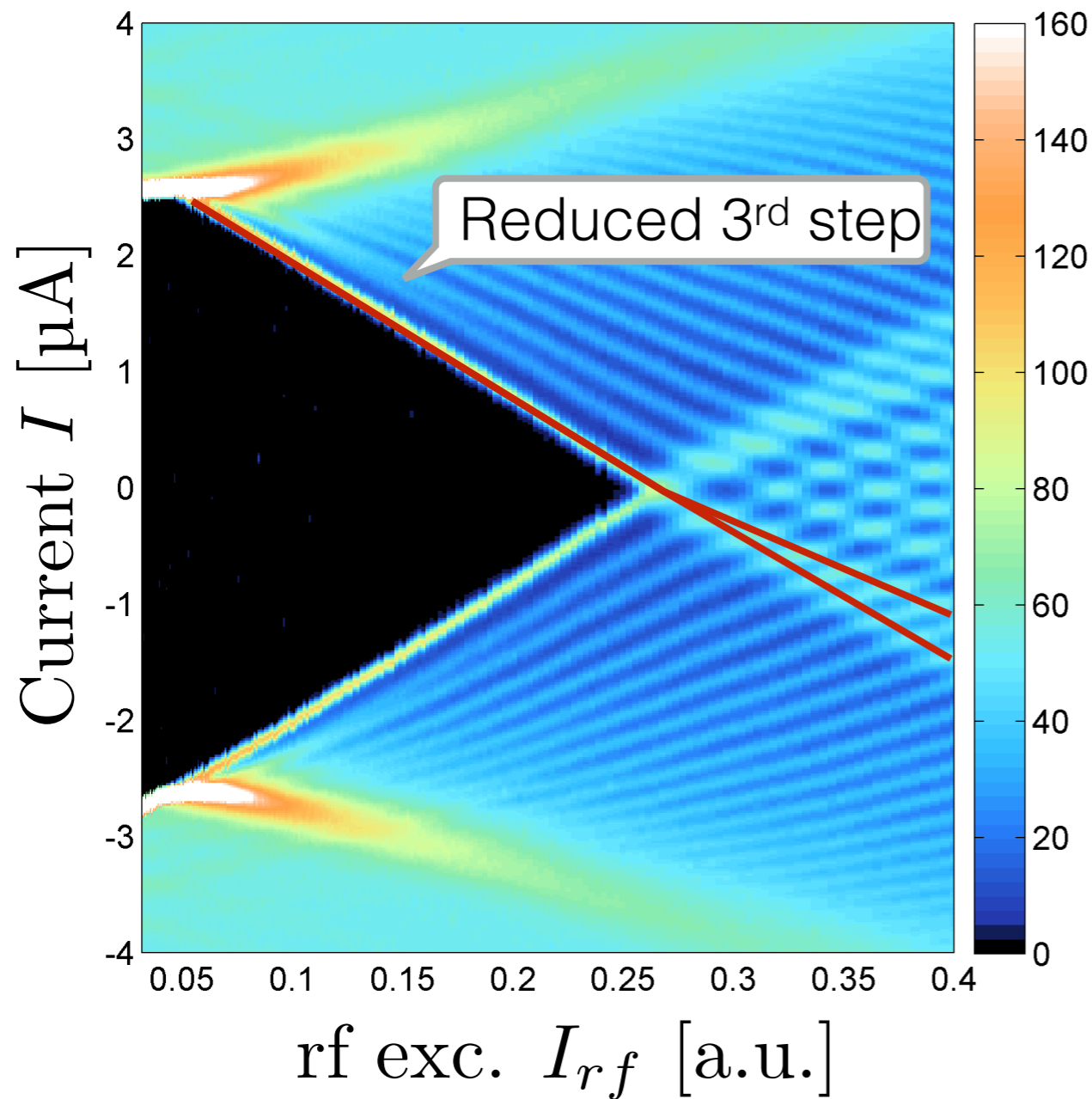


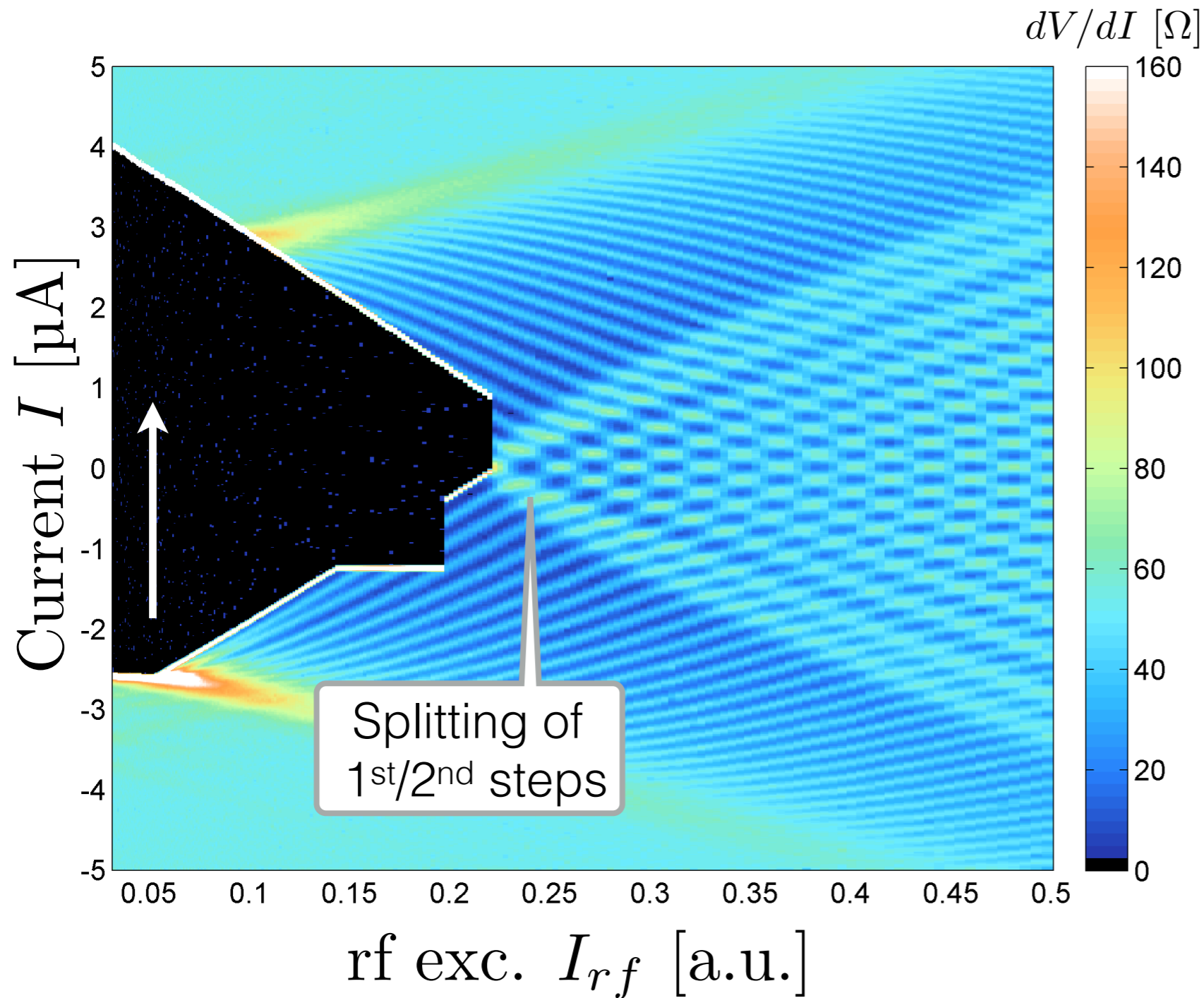
Landau-Zener transitions



- ▶ Andreev bound states $\varepsilon_{\pm}(\phi) = \pm \sqrt{4\delta^2 + \Delta_i^2 \cos^2 \phi/2}$
- ▶ Landau-Zener transition probability $P_{LZ} = \exp\left(-2\pi \frac{4\delta^2}{\Delta_i \hbar \dot{\phi}}\right)$

Weak suppression of third step





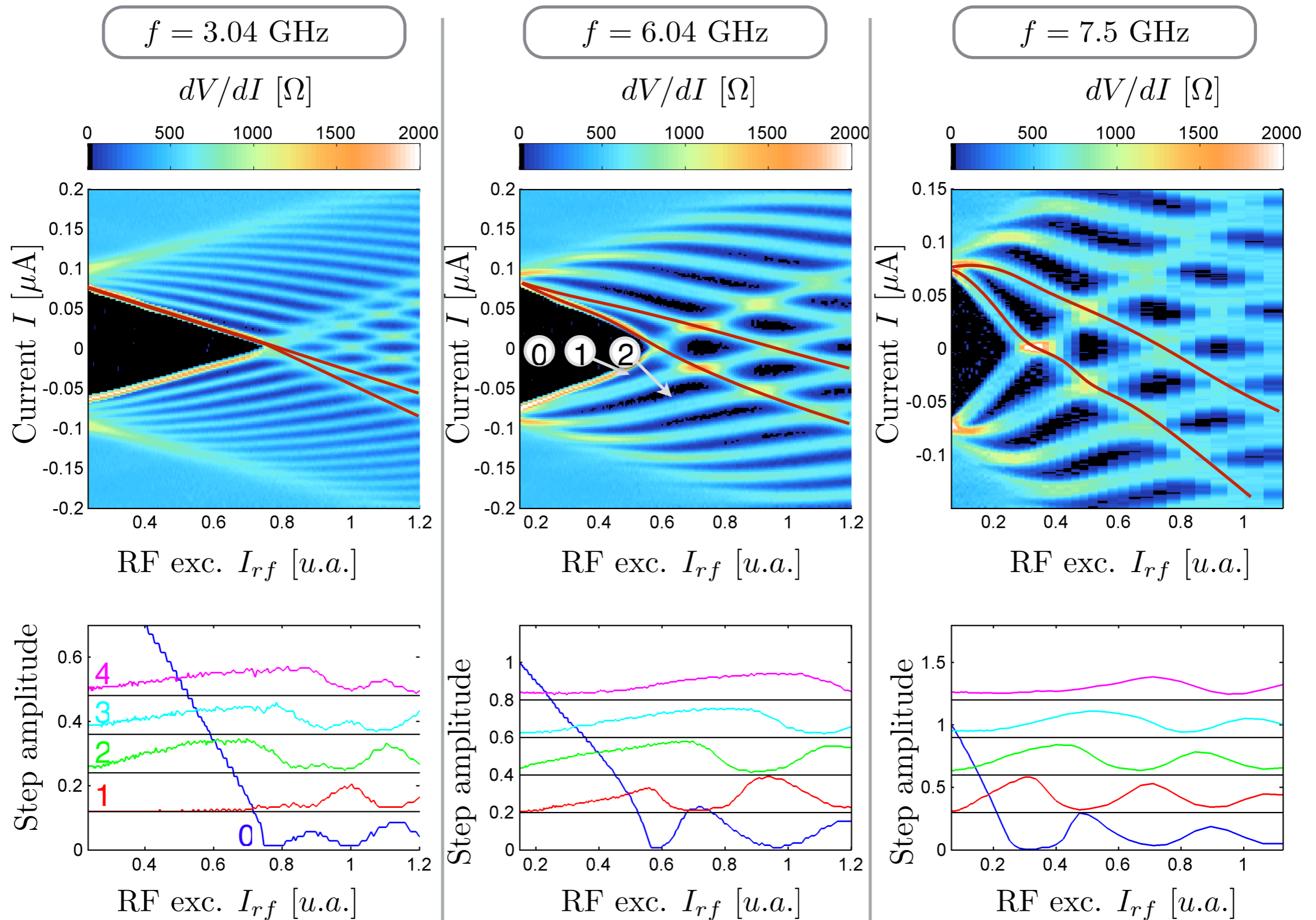
Hysteresis

- ▷ unknown origin
- ▷ common in similar systems (graphene, TIs)

Shapiro steps

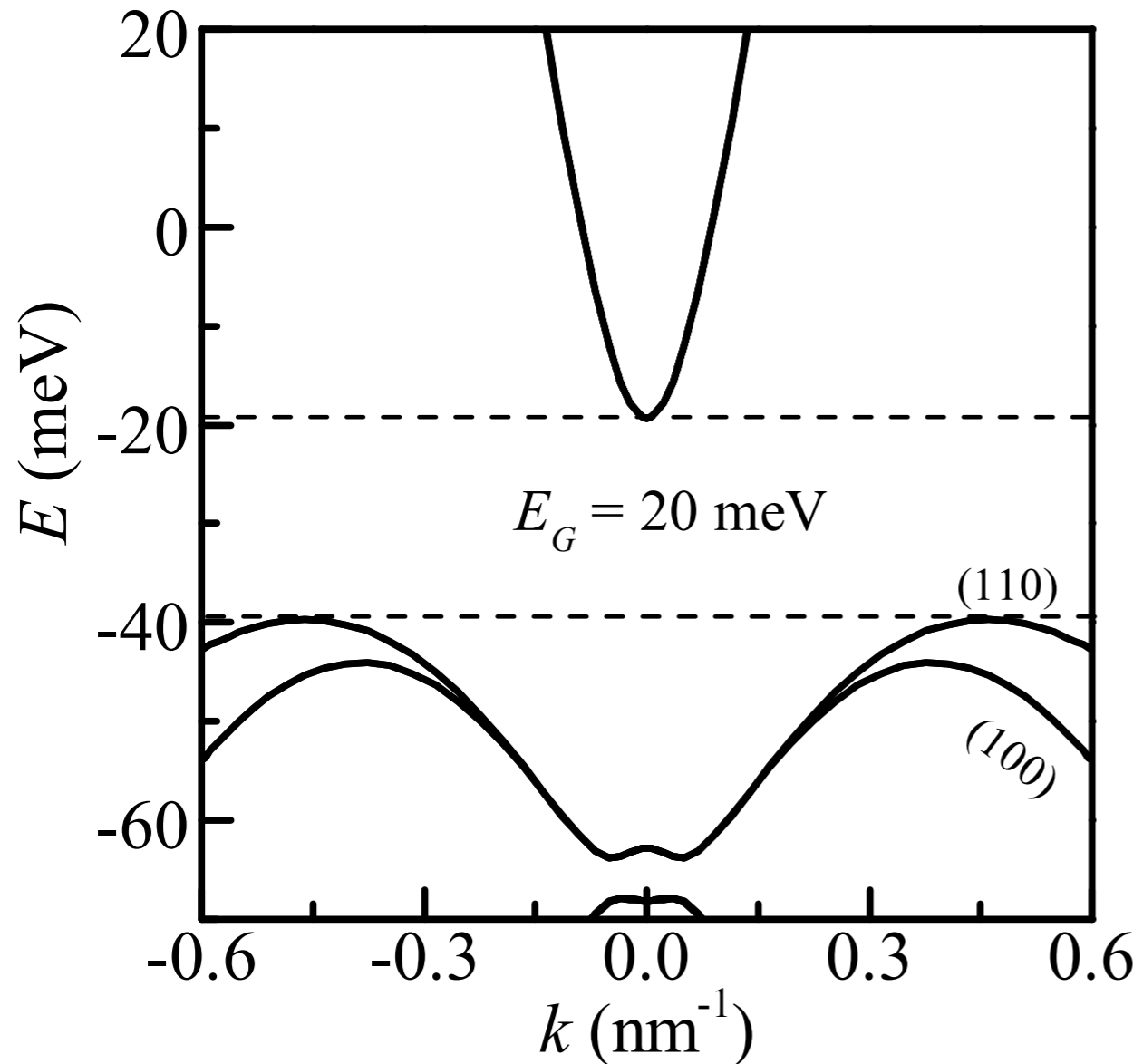
- ▷ bistability
⇒ no phase locking
- ▷ Shapiro steps invisible

Frequency dependence (2)



$$d_{\text{QW}} = 8.0 \text{ nm}$$

$$a_{\text{sub}} = 6.467 \text{ Angstrom (Cd}_{0.96}\text{Zn}_{0.04}\text{Te)}$$



- ▶ « camelback » structure
⇒ low mobility in p-regime
- ▶ gap around 20 meV
- ▶ edge states not calculated

Resonances at high frequencies

